

Machine learning for energy landscapes

Van 't Hoff Institute for Molecular Sciences & Informatics Institute University of Amsterdam

MAX PLANCK INSTITUTE FOR POLYMER RESEARCH

Tristan Bereau





Introduction to kernel-based ML Today

Machine learning for energy landscapes

Van 't Hoff Institute for Molecular Sciences & Informatics Institute University of Amsterdam

X PLANCK INSTITUTE FOR POLYMER RESEARCH

Incorporation of physical symmetries, conservation laws

Tristan Bereau





Introduction to kernel-based ML Today

Machine learning for energy landscapes

Supervised ML in chemistry and materials science

Thursday

X PLANCK INSTITUTE FOR POLYMER RESEARCH

Incorporation of physical symmetries, conservation laws

Tristan Bereau

- Van 't Hoff Institute for Molecular Sciences & Informatics Institute
 - University of Amsterdam





Kevin Jablonka Mohamad Moosavi















$\mathbf{F} = m\mathbf{a}$







$\mathbf{F} = m\mathbf{a}$ Specify interparticle forces: "force field"







Numerically integrate particle positions



$\mathbf{F} = ma^{\prime}$

Specify interparticle forces: "force field"







Numerically integrate particle positions



$\mathbf{F} = ma^{\prime}$

Specify interparticle forces: "force field"













Numerically integrate particle positions



$\mathbf{F} = ma^*$

Specify interparticle forces: "force field"





S

 μ S

MS









Numerically integrate particle positions



$\mathbf{F} = ma^{\prime}$

Specify interparticle forces: "force field"





S

 μ S

MS









Numerically integrate particle positions



$\mathbf{F} = ma^{\prime}$

Specify interparticle forces: "force field"





Timescales of interest











Numerically integrate particle positions



$\mathbf{F} = m\mathbf{a}$

Specify interparticle forces: "force field"











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Data as the 4th pillar of science



A. Agrawal and A. Choudhary APL Mater. 4 053208 (2016); https://www.bigmax.mpg.de





3rd PARADIGM Computational Science, Simulations

4th PARADIGM **Big-Data-Driven Science**

Monte Carlo; molecular dynamics; density-functional theory and beyond

Detection of patterns and anomalies in Big Data; artificial intelligence; etc.



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Databases





Data science

Hardware

Machine learning











Databa







Data science

Machine learning









Links to machine learning

Potential energy surface



Interpolation of a high-dimensional function

Can we build a more accurate PES?

Can we **easily** build an accurate PES?

Can we make the numerical integration faster and/or more efficient?

. . .







Good

Not good



Teaser: let's fit data points









Teaser: let's fit data points



Good

Not good









Teaser: let's fit data points













ruthlessly stolen from A. von Lilienfeld

Teaser: let's fit data points









Multivariate function approximation

Sparse data



infer smooth function



7



Sparse data







7



Sparse data



THEORY GROUP Rupp, International Journal of Quantum Chemistry 115 (2015)







7



Sparse data



THEORY Rupp, International Journal of Quantum Chemistry 115 (2015)



Regression:

prediction of $f: \mathbb{R}^d \to \mathbb{R}$ based on noisy data points $D = (\mathbf{X}, \mathbf{y}) = {\{\mathbf{x}_{n}, y_{n}\}_{n=1}^{N}}$

$$y_n = f(\mathbf{x}_n) + \varepsilon$$

What is



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Kernel methods are vintage





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Kernel

needs a representation linear algebra can be efficient with small data

Deep learning

learns the representation complex mathematical structure data hungry











Learning from experience

Inductive (based on examples)



Input Space

Feature Space





source: xkcd





Extrapolation in machine learning









Extrapolation in machine learning











Prediction



Bayesian inference

Prior beliefs

Sampled data







Prediction





Bayesian inference

Prior beliefs

Sampled data







Prediction



posterior



Bayesian inference



Bayes' formula











THEORY GROUP Rasmussen, Advanced lectures on machine learning. Springer, 63-71 (2004)





Gaussian processes



THEORY GROUP Rasmussen, Advanced lectures on machine learning. Springer, 63-71 (2004)

 $f \sim \mathcal{GP}(m,k)$





 $f \sim \mathcal{GP}(m,k)$ $\mu_i = m(x_i) \qquad \Sigma_{ij} = k(x_i, x_j)$

f(x)

random variable: value of the stochastic function at x

THEORY Rasmussen, Advanced lectures on machine learning. Springer, 63-71 (2004)

Gaussian processes

mean

covariance





Gaussian processes $f \sim \mathcal{GP}(m,k)$ $\mu_i = m(x_i) \qquad \Sigma_{ij} = k(x_i, x_j)$ covariance mean kernel

f(x)

random variable: value of the stochastic function at x



THEORY GROUP Rasmussen, Advanced lectures on machine learning. Springer, 63-71 (2004)





Linear-ridge regression





Linear-ridge vs kernel-ridge regression

- kernel-ridge regression (ML)
 - $\mathbf{K} \alpha = p$








- kernel-ridge regression (ML)
 - $\mathbf{K} \alpha = p$







in general: $m \ll N$

- kernel-ridge regression (ML)
 - $\mathbf{K} \alpha = p$

in general: $m \ll N$

- kernel-ridge regression (ML) $\mathbf{K} \alpha = p$ N ______

in general: $m \ll N$

kernel-ridge regression (ML) $\mathbf{K}\alpha = p$ linear N# training points NN $K_{ij} = K_{ij}(\mathbf{x}_i, \mathbf{x}_j)$ $=K_{ij}(|\mathbf{x}_i - \mathbf{x}_j|)$ $= \exp\left(-\frac{|\mathbf{x}_i - \mathbf{x}_j|}{\sigma}\right)$

Kernel machine learning 101

1) Define representation and kernel

$$K(r, r') = \exp\left(-\frac{(r - r')^2}{2\sigma^2}\right)$$

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Kernel machine learning 101

1) Define representation and kernel

$$K(r, r') = \exp\left(-\frac{(r - r')^2}{2\sigma^2}\right)$$

2) Train your model:

 $(\mathbf{K} + \lambda \mathbf{I}) \alpha = U$

 $\alpha = (K + \lambda \mathbb{I})^{-1} U$

Inverse is ill defined

Regularization: "hyperparameter" scales noise level

Optimize weight coefficients on training set

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Kernel machine learning 101

1) Define representation and kernel

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Conformational space missing from training

Linking conformational and interpolation spaces 📿

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Symmetries and conservation laws

Mechanics 101: Principle of least action

Mechanics 101: Principle of least action $S[x(t)] = \int_{t_1}^{t_2} dt L[x(t), \dot{x}(t), t]$ kinetic energy action potential energy $_$ agrangian L = T

microtrajector

Hamilton's principle: system minimizes action (variational principle)

$\dot{\mathcal{S}}[x^*(t)] = 0$

Mechanics 101: Principle of least action kinetic energy potential energy

microtrajector

Hamilton's principle: system minimizes action (variational principle)

stationarity under small perturbations leads to Euler-Lagrange equations

$$\delta \mathcal{S} = \int_{t_1}^{t_2} \mathrm{d}t \, L(x^* + \varepsilon, \dot{x}^* + \dot{\varepsilon}, t) - L(x^*, \dot{x}^*, t)$$
$$= \int_{t_1}^{t_2} \mathrm{d}t \, \left(\varepsilon \frac{\partial L}{\partial x} + \dot{\varepsilon} \frac{\partial L}{\partial x}\right) = \int_{t_1}^{t_2} \mathrm{d}t \, \left(\varepsilon \frac{\partial L}{\partial x} - \varepsilon \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{x}}\right) = 0$$

$\mathcal{S}[x^*(t)] = 0$

Mechanics 101: Principle of least action $\mathcal{S}[x(t)] = \int_{t_1}^{t_2} dt \, L[x(t), \dot{x}(t), t] \qquad \qquad \text{kinetic end}$ action $Lagrangian \ L = T - Lagrangian$ kinetic energy potential energy

microtrajector

Hamilton's principle: system minimizes action (variational principle)

stationarity under small perturbations leads to Euler-Lagrange equations

 $\varepsilon(t_1) = \varepsilon(t_2) = 0$

$$\delta S = \int_{t_1}^{t_2} dt \, L(x^* + \varepsilon, \dot{x}^* + \dot{\varepsilon}, t) - L(x^*, \dot{x}^*, t)$$
$$= \int_{t_1}^{t_2} dt \, \left(\varepsilon \frac{\partial L}{\partial x} + \dot{\varepsilon} \frac{\partial L}{\partial x}\right) = \int_{t_1}^{t_2} dt \, \left(\varepsilon \frac{\partial L}{\partial x} - \varepsilon \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}\right) = 0$$

Integration by parts &

$\mathcal{S}[x^*(t)] = 0$

From symmetries, to invariants, to conserved quantities

 $\mathcal{S}[x(t), y(t), z(t)] = \left[dt \, \frac{m}{2} \left(\dot{x}^2 + \dot{y}^2 + \dot{z}^2 \right) - mgz \right]$

- $\mathcal{S}[x(t), y(t), z(t)] =$
- Introduce constant translations along x and y:
 - $S[x(t) + x_0, y(t) + y_0, z(t)]$

$$\mathrm{d}t\,\frac{m}{2}\left(\dot{x}^2+\dot{y}^2+\dot{z}^2\right)-mgz$$

$$= \int dt \frac{m}{2} \left(\dot{x}^2 + \dot{y}^2 + \dot{z}^2 \right) - mgz$$
$$= \mathcal{S}[x(t), y(t), z(t)]$$

- $\mathcal{S}[x(t), y(t), z(t)] =$
- Introduce constant translations along *x* and *y*:
 - $S[x(t) + x_0, y(t) + y_0, z(t)]$

(Translational) symmetry leaves the action invariant. It leaves the Euler-Lagrange equation unchanged:

$$\mathrm{d}t\,\frac{m}{2}\left(\dot{x}^2+\dot{y}^2+\dot{z}^2\right)-mgz$$

$$= \int dt \frac{m}{2} \left(\dot{x}^2 + \dot{y}^2 + \dot{z}^2 \right) - mgz$$
$$= \mathcal{S}[x(t), y(t), z(t)]$$

 $\frac{\partial L}{\partial x} - \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{x}} = 0$

$$\mathcal{S}[x(t), y(t), z(t)] = \int dt \, \frac{m}{2} \left(\dot{x}^2 + \dot{y}^2 + \dot{z}^2 \right) - mgz$$

Introduce constant translations along x and y:

$$\mathcal{S}[x(t) + x_0, y(t) + y_0, z(t)] = \int dt \, \frac{m}{2} \left(\dot{x}^2 + \dot{y}^2 + \dot{z}^2 \right) - mgz$$
$$= \mathcal{S}[x(t), y(t), z(t)]$$

(Translational) symmetry leaves the action invariant. It leaves the Euler-Lagrange equation unchanged:

$$\frac{\partial L}{\partial x} = 0 \qquad \qquad \frac{\partial L}{\partial \dot{x}} = m\dot{x} = \text{const.}$$

From symmetries, to invariants, to conserved quantities

 $\frac{\partial L}{\partial L} - \frac{\mathrm{d}}{\partial L} = 0$ $dt \partial \dot{x}$ ∂x

> Translational invariance implies linear momentum conversation

 $\mathcal{S}[\mathbf{r}(t)] =$

Rotational symmetry

Apply transformation $\mathbf{r} \rightarrow \mathbf{r}'$ where $\mathbf{r}'(t) = R\mathbf{r}(t) = \mathbf{r}(t) + \alpha \times \mathbf{r}(t)$

One can show that $\mathcal{S}[\mathbf{r}(t) + \alpha \times \mathbf{r}(t)] = \mathcal{S}[\mathbf{r}(t)]$

Conservation of angular momentum

From symmetries to conserved quantities (cont'd)

$$\int \mathrm{d}t \, \frac{m}{2} \dot{\mathbf{r}}^2 - V(r)$$

Time translation

Apply transformation $\mathbf{r} \rightarrow \mathbf{r}'$ where $\mathbf{r}'(t+\epsilon) = \mathbf{r}(t)$

One can show that $\mathcal{S}[\mathbf{r}'(t+\epsilon)] = \mathcal{S}[\mathbf{r}(t)]$ (up to a boundary term) Conservation of energy

Noether's theorem

To every differentiable symmetry generated by local actions there corresponds a conserved quantity

3 examples:

- Translational symmetry: Linear momentum conservation
- Rotational symmetry: Angular momentum conservation
- Time translation: Energy conservation

mentum conservation mentum conservation ation

Noether's theorem

To every differentiable symmetry generated by local actions there corresponds a conserved quantity

3 examples:

- Translational symmetry: Linear momentum cons
- Rotational symmetry: Angular momentum conset
- Time translation: Energy conservation

mentum conservation

2 ways of encoding symmetries:- Representation- ML model

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Encoding symmetries in the representation

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Translational and rotational symmetries Behler-Parrinello Coulomb matrix

$$G_{i}^{1} = \sum_{j \neq i}^{\text{all}} e^{-\eta(R_{ij} - R_{s})^{2}} f_{c}(R_{ij})$$

$$G_{i}^{2} = 2^{1-\zeta} \sum_{j,k \neq i}^{\text{all}} (1 + \lambda \cos \theta_{ijk})^{\zeta}$$

$$\times e^{-\eta(R_{ij}^{2} + R_{ik}^{2} + R_{jk}^{2})} f_{c}(R_{ij}) f_{c}(R_{ik}) f_{c}(R_{jk})$$

Distances
Angles

Behler & Parrinello, Phys Rev Lett 98 (2007)

Rupp, Tkatchenko, Müller, von Lilienfeld, Phys Rev Lett, 108 (2012)

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Representation: the Coulomb matrix

~ Coulomb's law $E = \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$

Hansen et al., *J Chem Theory Comput*, **9** (2013)

Symmetries of the representation should emulate symmetries of the system

- Translation
- 2. Rotations
- 3. Mirror reflection

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Representation: the Coulomb matrix

Η

0.2

2.9

Η

0.2

2.9

0.5

~ Coulomb's law $E = \frac{q_i q_j}{|\mathbf{r}_i - \mathbf{r}_j|}$

Hansen et al., *J Chem Theory Comput*, **9** (2013)

Symmetries of the representation should emulate symmetries of the system

- 1. Translation
- 2. Rotations
- 3. Mirror reflection

THEORY GROUP Huang and von Lilienfeld, J Chem Phys 145 (2016)

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Optimizing the representation links to the physics

THEORY GROUP Huang and von Lilienfeld, J Chem Phys 145 (2016)

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Encoding symmetries in the ML model

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Action of group G on input sample $x \mapsto T_{g}(x)$

THEORY GROUP Risi Kondor, *Group theoretical methods in machine learning, PhD thesis* (2008)

Encoding symmetries in ML models using group theory



Action of group G on input sample $x \mapsto T_{\rho}(x)$

Can we find a kernel that is invariant to this group action? $f(T_g(x)) = f(x) \forall g \in G$

THEORY GROUP Risi Kondor, Group theoretical methods in machine learning, PhD thesis (2008)

Encoding symmetries in ML models using group theory

 $k(x, x') = k(T_g(x), T_{g'}(x'))$









Action of group G on input sample $x \mapsto T_{\rho}(x)$

Can we find a kernel that is invariant to this group action?

To ensure invariance, symmetrize the kernel

 $k^{G}(x, x') = \frac{1}{|G|} \sum_{k \in V} k(x, T_{g}(x'))$ $g \in G$

Risi Kondor, Group theoretical methods in machine learning, PhD thesis (2008)

Encoding symmetries in ML models using group theory

 $f(T_g(x)) = f(x) \forall g \in G$ $k(x, x') = k(T_g(x), T_{g'}(x'))$









Example of symmetrized kernel

Vanilla/naïve kernel





























Tensorial property (e.g., dipole moment, force) **rotates** with the sample









Tensorial property (e.g., dipole moment, force) **rotates** with the sample











Tensorial property (e.g., dipole moment, force) **rotates** with the sample











Tensorial property (e.g., dipole moment, force) **rotates** with the sample



"Build kernel so as to encode the rotational properties of the target property"











THEORY GROUP Glielmo, Sollich, De Vita, Phys Rev B 95 (2017)

Covariant kernels

Encode rotational properties of the target property in the **kernel**









$\hat{\mathbf{f}}(\mathcal{S}\rho \mid \mathcal{D}) = \hat{\mathbf{S}}\hat{\mathbf{f}}(\rho \mid \mathcal{D})$ Force prediction Transformation (rotation/inversion) Descriptor Training data

THEORY GROUP Glielmo, Sollich, De Vita, Phys Rev B 95 (2017)

Covariant kernels

Encode rotational properties of the target property in the **kernel**









$\mathbf{f}(\mathcal{S}\rho \mid \mathcal{D}) = \mathbf{S}\hat{\mathbf{f}}(\rho \mid \mathcal{D})$ Force prediction Transformation (rotation/inversion) Descripto Training data

THEORY Glielmo, Sollich, De Vita, *Phys Rev B* **95** (2017)

Covariant kernels

Encode rotational properties of the target property in the **kernel**

"Transform the configuration, and the prediction transforms with it"







Covariant kernels $\mathbf{K}(\mathcal{S}\rho,\mathcal{S}'\rho') = \mathbf{S}\mathbf{K}(\rho,\rho')\mathbf{S}'^{\mathrm{T}}$ Kernel Configurations

Transformations (rotation/inversion)

THEORY Glielmo, Sollich, De Vita, Phys Rev B 95 (2017)

 $\mathbf{K}(\boldsymbol{\rho},\boldsymbol{\rho}') = \int d\mathcal{R}\mathbf{R}k_b(\boldsymbol{\rho},\mathcal{R}\boldsymbol{\rho}')$

 $\mathbf{K}^{\mu}(\rho, \rho') = \frac{1}{L} \sum_{ij} \phi(r_i, r_j) \mathbf{r}_i \otimes \mathbf{r}_j'^{\mathrm{T}}$











Conclusions



Extrapolation in ML models of energy landscapes Can lead to catastrophic physics

Take advantage of symmetries Noether: symmetry leads to conservation law

$\mathbf{K}(\mathcal{S}\rho, \mathcal{S}'\rho') = \mathbf{S}\mathbf{K}(\rho, \rho')\mathbf{S}'^{\mathrm{T}}$







Build symmetries in ML model

Work with subset of kernels that a priori satisfy conservation law

