

# *Free Energy and Phase Equilibria*

Thermodynamic Integration (7.1)

Chemical Potentials (7.2)

Overlapping Distributions (7.2)

Umbrella Sampling (7.4)

(Application: Phase Diagram of Carbon)

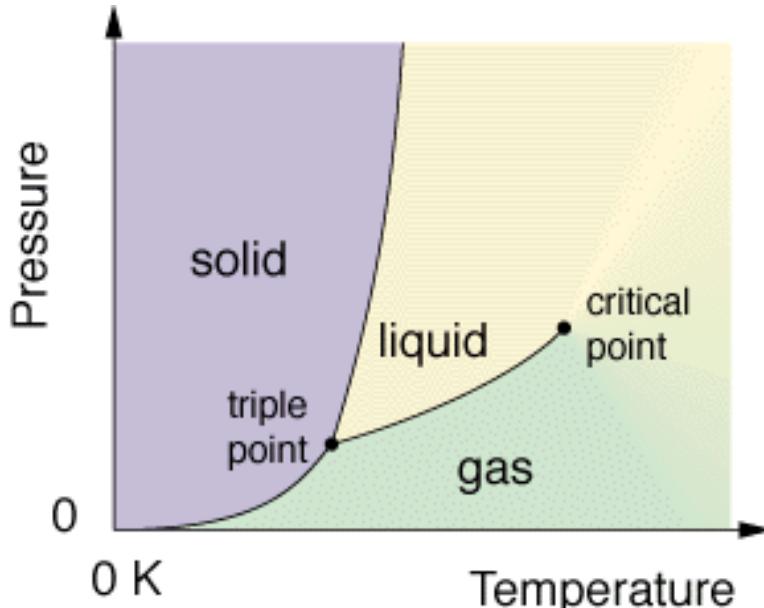
# *Why Free Energies?*

- Reaction equilibrium constants  $A \leftrightarrow B$

$$K = \frac{[B]}{[A]} = \frac{p_B}{p_A} = \exp[-\beta(G_B - G_A)]$$

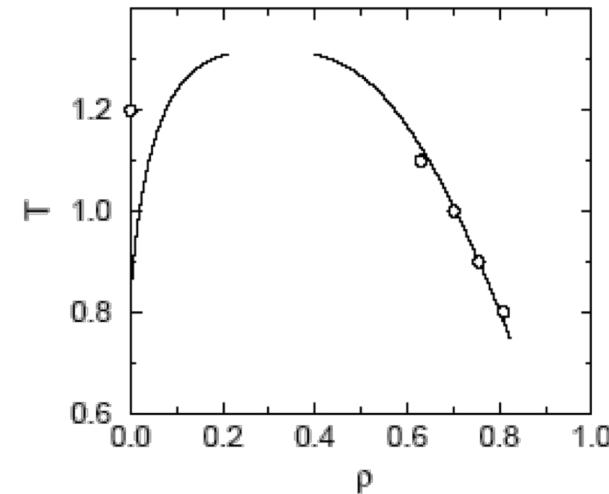
- Examples:
  - Chemical reactions: e.g. catalysis, etc....
  - Protein folding, binding: free energy gives binding constants
- Phase diagrams
  - Prediction of thermodynamic stability of phases
  - Coexistence lines
  - Critical points
  - Triple points
  - First order/second order phase transitions

# Phase diagrams



Critical point: no transition between liquid and vapor

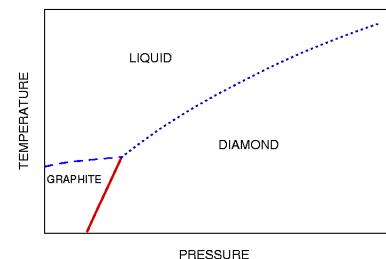
Triple point: liquid, vapor and solid in equilibrium.



Along the liquid-gas coexistence line increasing the pressure and temperature at constant volume the liquid density becomes lower and the vapor density higher.

How do we compute these lines?

Carbon Phase Diagram



# *Phase equilibrium*

Criteria for equilibrium (for single component)

$$T_I = T_{II} \quad P_I = P_{II} \quad \mu_I = \mu_{II}$$

Chemical potential

$$\mu = \left( \frac{\partial F}{\partial N} \right)_{V,T} = \left( \frac{\partial G}{\partial N} \right)_{P,T} = G_m$$

If  $\mu_I > \mu_{II}$  : transport of particles from phase I to phase II.

**Stable phase:**

**Lowest chemical potential (for single phase: lowest Gibbs free energy)**

# *Relation thermodynamic potentials*

Helmholtz free energy:  $F = U - TS$

Gibbs free energy:  $G = F + PV$

Suppose we have  $F(n, V, T)$

Then we can find G from F from:

$$P = -\left(\frac{\partial F}{\partial V}\right)_{n,T}$$

$$G = F - \left(\frac{\partial F}{\partial V}\right)_{n,T} V$$

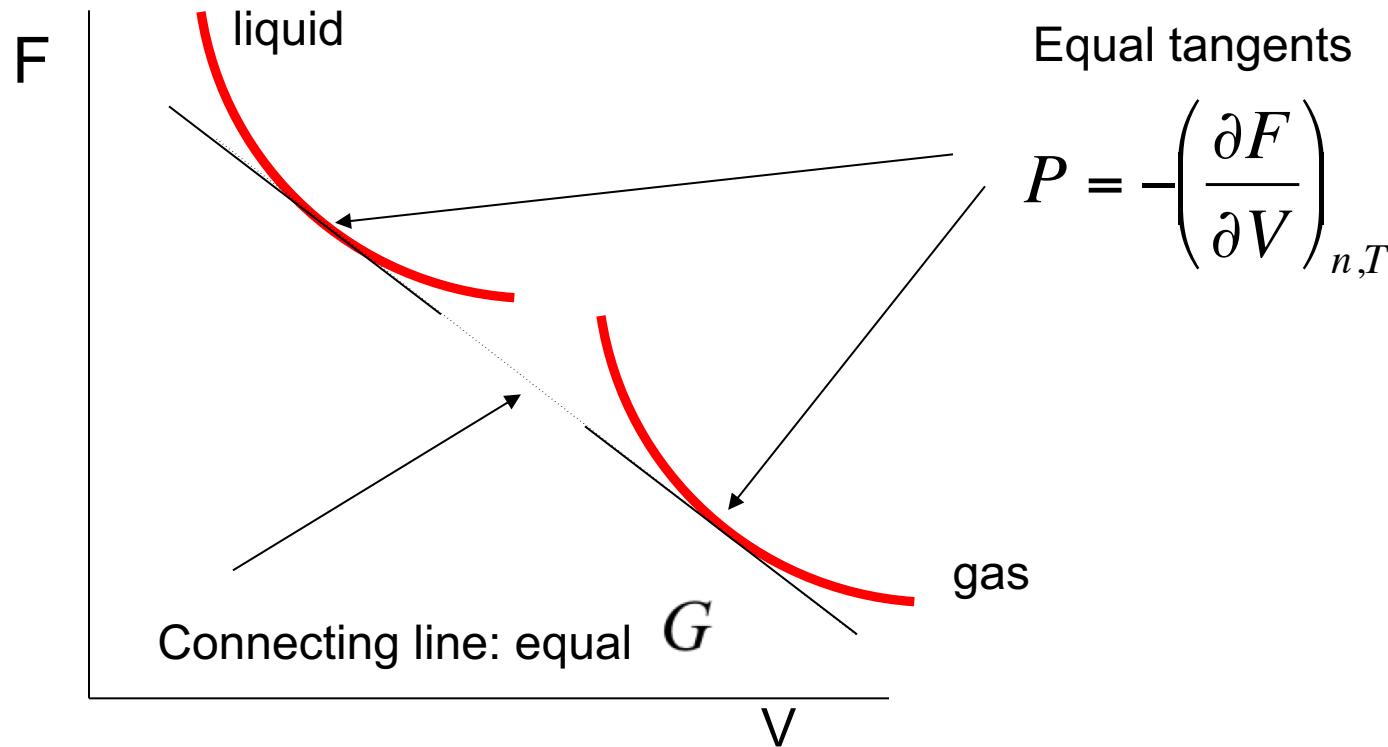
All thermodynamic quantities can be derived from F and its derivatives

# Phase equilibria from $F(V, T)$

Common tangent construction

$$0 = \Delta G = \Delta F + P\Delta V$$

$$\Delta F = -P\Delta V$$



# We need $F$ (or $G$ )

- We can calculate  $F(V)$ , using equation of state  $P(V)$

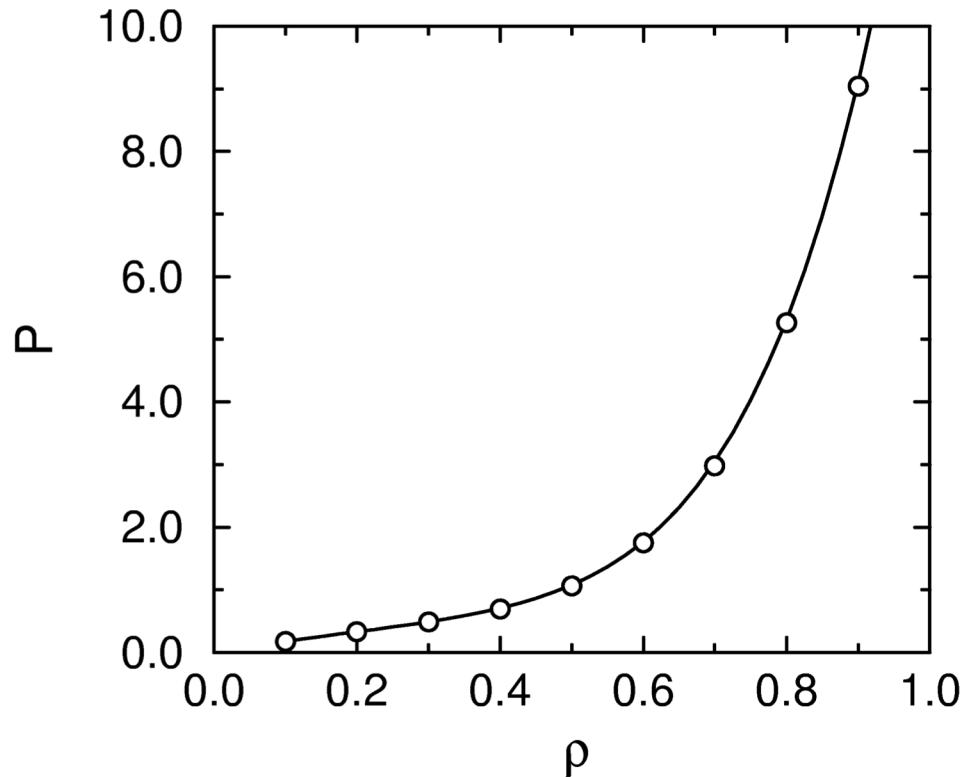
$$F(V) = F(V_0) + \int_{V_0}^V \left( \frac{\partial F}{\partial V} \right)_{N,T} dV = F(V_0) - \int P dV$$

$$F(\rho) = F(\rho_0) + N \int_{\rho_0}^{\rho} \frac{P(\rho')}{\rho'^2} d\rho' \quad (V=N/\rho)$$

- So in fact for only 1 point of the equation of state the  $F$  is needed
- For liquid e.o.s even from ideal gas

$$\beta F(\rho)/N = \beta F^{id}(\rho)/N + \int_0^{\rho} \frac{\beta P(\rho') - \rho'}{\rho'^2} d\rho'$$

# Equation of state



$$P = P(\rho, T)$$

$$\left( \frac{\partial F}{\partial V} \right)_{N,T} = -P$$

$$F(\rho) = F(\rho_0) + N \int_{\rho_0}^{\rho} \frac{P(\rho')}{\rho'^2} d\rho'$$

$$F(\rho) = F_{id}(\rho) + (F(\rho) - F_{id}(\rho))$$

$$\beta F(\rho)/N = \beta F^{id}(\rho)/N + \int_0^{\rho} \frac{\beta P(\rho') - \rho'}{\rho'^2} d\rho'$$

# *Free Energies and Phase Equilibria*

## General Strategies

- Determine free energy of both phases separately, relative to a reference state  
Free-energy difference calculation  
*General applicable: Gas, Liquid, Solid, Inhomogeneous systems, ...*
- Determine free energy difference between two phases  
Gibbs Ensemble  
*Specific applicable: Gas, Liquid*

# **Statistical Thermodynamics**

Probability to find a particular configuration (NVT)

$$P(\mathbf{r}^N) = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}'^N \delta(\mathbf{r}'^N - \mathbf{r}^N) \exp[-\beta U(\mathbf{r}'^N)] \propto \exp[-\beta U(\mathbf{r}^N)]$$

Partition function

$$Q_{NVT} = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]$$

Free energy

$$\beta F = -\ln(Q_{NVT})$$

Ensemble average

$$\langle A \rangle_{NVT} = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]$$

# Ensemble average versus free energy

Generate configuration using MC:  $\{r_1^N, r_2^N, r_3^N, r_4^N \dots, r_M^N\}$

$$\bar{A} = \frac{1}{M} \sum_{i=1}^M A(r_i^N) \approx \frac{\int dr^N A(r^N) \exp[-\beta U(r^N)]}{\int dr^N \exp[-\beta U(r^N)]} = \langle A \rangle_{NVT}$$

Generate configuration using MD:  $\{r_1^N, r_2^N, r_3^N, r_4^N \dots, r_M^N\}$

$$\bar{A} = \frac{1}{M} \sum_{i=1}^M A(r_i^N) \approx \frac{1}{T} \int_0^T dt A(t) \int \approx \langle A \rangle_{NVT}$$

ergodicity

$$\beta F = -\ln Q_{NVT} = -\ln \frac{1}{\Lambda^{3N} N!} \int dr^N \exp[-\beta U(r^N)]$$

F is difficult, because requires measuring the phase space volume

# I - Thermodynamic integration

- Known reference state       $\lambda=0$
- Unknown target state       $\lambda=1$

The diagram illustrates the coupling parameter  $\lambda$  as a bridge between two systems. A light blue rounded rectangle contains the text "Coupling parameter". Two lines extend from the bottom-left corner of this box to two separate labels: "Reference System" at the top right and "Target System" below it. Below the box, the equation for the potential energy  $U(\lambda)$  is given as a weighted sum of the reference and target potentials:  $U(\lambda) = (1 - \lambda)U_I + \lambda U_{II}$ . Further down, the expression for the partition function  $Q_{NVT}(\lambda)$  is shown as:  $Q_{NVT}(\lambda) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\lambda)]$ .

$$U(\lambda) = (1 - \lambda)U_I + \lambda U_{II}$$
$$Q_{NVT}(\lambda) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\lambda)]$$

$$F(\lambda = 1) - F(\lambda = 0) = \int_{\lambda=0}^{\lambda=1} d\lambda \left( \frac{\partial F(\lambda)}{\partial \lambda} \right)_{N,V,T}$$

# Thermodynamic integration

$$\left( \frac{\partial F(\lambda)}{\partial \lambda} \right)_{N,T} = -\frac{1}{\beta} \frac{\partial}{\partial \lambda} \ln(Q) = -\frac{1}{\beta} \frac{1}{Q} \frac{\partial Q}{\partial \lambda}$$

$$= \frac{\int d\mathbf{r}^N (\partial U(\lambda)/\partial \lambda) \exp[-\beta U(\lambda)]}{\int d\mathbf{r}^N \exp[-\beta U(\lambda)]}$$

$$= \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_\lambda$$

Free-energy difference  
as ensemble average!

$$F(\lambda = 1) - F(\lambda = 0) = \int d\lambda \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_\lambda$$

$$Q_{NVT}(\lambda) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\lambda)]$$

$$F(\lambda = 1) - F(\lambda = 0) = \int_{\lambda=0}^{\lambda=1} d\lambda \left( \frac{\partial F(\lambda)}{\partial \lambda} \right)_{N,V,T}$$

# *Example*

- In general

$$U(\lambda) = (1 - \lambda)U_I + \lambda U_{II}$$
$$\left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_\lambda = \langle U_{II} - U_I \rangle_\lambda$$

- Specific example

$$U(\lambda) = U^{LJ} + \lambda U^{\text{dipole-dipole}}$$

$$U(0) = U^{LJ} \qquad \qquad \qquad \text{Lennard-Jones}$$

$$U(1) = U^{\text{Stockm}} \qquad \qquad \qquad \text{Stockmayer}$$

$$\left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_\lambda = \langle U^{\text{dip-dip}} \rangle_\lambda$$

# Free energy of solid

More difficult. What is reference?

Not the ideal gas.

One (natural) choice is an Einstein crystal: harmonic oscillators around  $r_0$

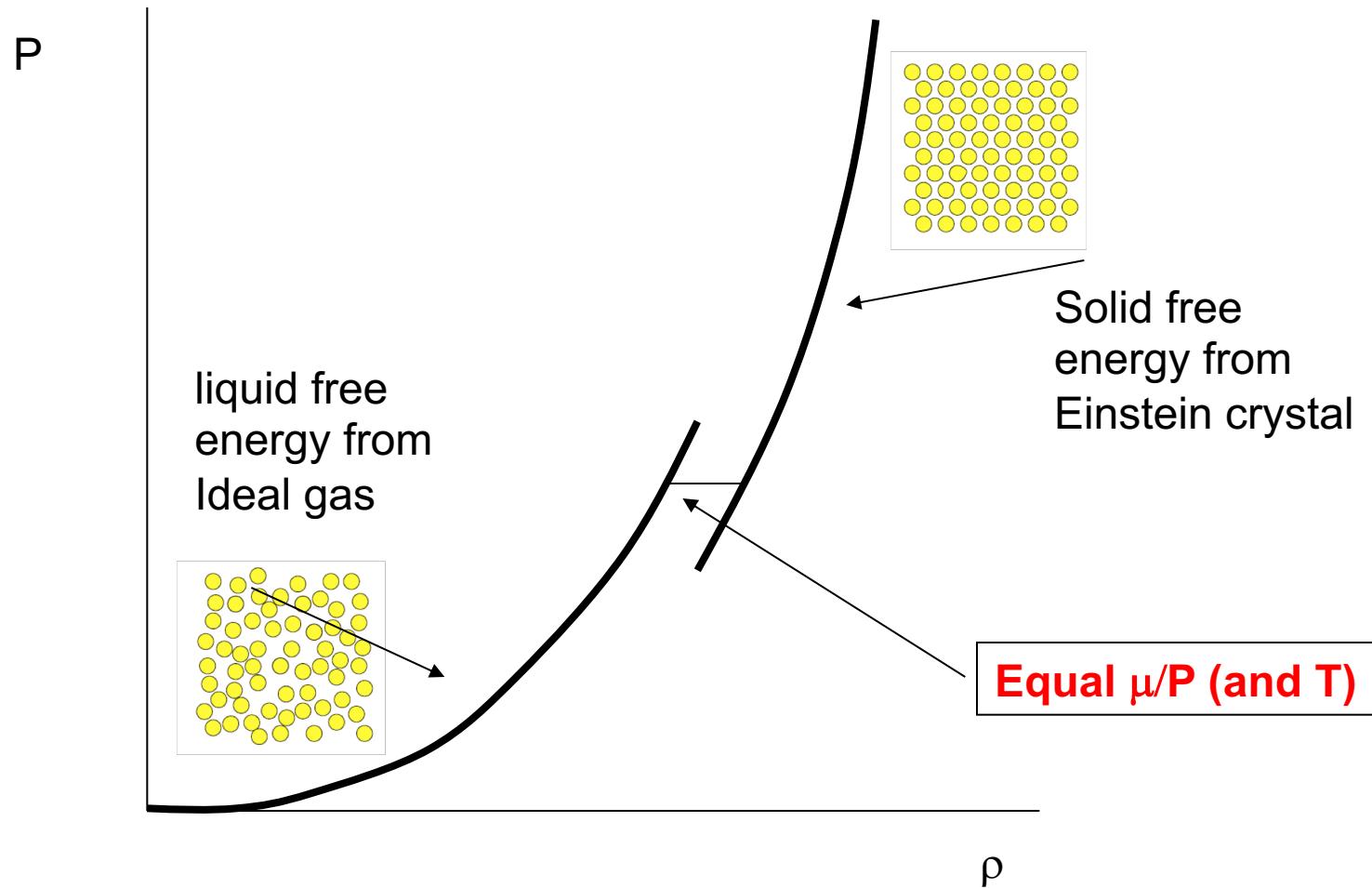
$$U(\lambda; r^N) = (1 - \lambda)U(r^N) + \boxed{\lambda U(r_0^N) + \lambda \sum_{i=1}^N \alpha(r_i - r_{0,i})^2}$$

$$F = F_{ein} - \int_{\lambda=0}^{\lambda=1} d\lambda \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_\lambda$$

Note, here:  
 $\lambda = 1$  Reference System  
 $\lambda = 0$  Target System

$$F = F_{ein} - \int_{\lambda=0}^{\lambda=1} d\lambda \left\langle -U(r^N) + U(r_0^N) + \sum_{i=1}^N \alpha(r_i - r_{0,i})^2 \right\rangle_\lambda$$

# *Hard sphere freezing*



## *II - Thermodynamic perturbation*

Two systems:

System 0:  $N, V, T, U_0$

System 1:  $N, V, T, U_1$

$$Q_0 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_0)$$

$$Q_1 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_1)$$

$$\Delta\beta F = \beta F_1 - \beta F_0 = -\ln(Q_1/Q_0)$$

$$= -\ln \frac{\int d\mathbf{s}^N \exp[-\beta U_1]}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$= -\ln \frac{\int d\mathbf{s}^N \exp[-\beta(U_1 - U_0)] \exp[-\beta U_0]}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$\Delta\beta F = -\ln \langle \exp[-\beta(U_1 - U_0)] \rangle_0$$

# **Particle Insertion Method**

$$Q_{NVT} = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]$$

$$\beta F = -\ln(Q_{NVT})$$

$$= -\ln\left(\frac{V^N}{\Lambda^{3N} N!}\right) - \ln\left(\int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]\right)$$

$$= -N \ln\left(\frac{1}{\Lambda^3 \rho}\right) + N - \ln\left(\int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]\right)$$

$$\beta F = \beta F^{IG} + \beta F^{ex}$$

$$\mu = \left( \frac{\partial F}{\partial N} \right)_{V,T}$$

$$\left. \begin{aligned} \beta \mu &= \beta \mu^{IG} + \beta \mu^{ex} \\ \beta \mu^{IG} &\equiv \left( \frac{\partial \beta F^{IG}}{\partial N} \right)_{V,T} \\ \beta \mu^{ex} &\equiv \left( \frac{\partial \beta F^{ex}}{\partial N} \right)_{V,T} \end{aligned} \right\}$$

# *Widom test particle insertion*

$$\beta\mu \equiv \left( \frac{\partial \beta F}{\partial N} \right)_{V,T}$$

$$\begin{aligned}\beta\mu &= \frac{\beta F(N+1) - \beta F(N)}{N+1 - N} \\ &= -\ln \frac{Q(N+1)}{Q(N)} \\ &= -\ln \left( \frac{\frac{V^{N+1}}{\Lambda^{3N+3}(N+1)!}}{\frac{V^N}{\Lambda^{3N}N!}} \right) - \ln \left( \frac{\int d\mathbf{s}^{N+1} \exp[-\beta U(\mathbf{s}^{N+1}; L)]}{\int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]} \right) \\ &= -\ln \left( \frac{V}{\Lambda^3(N+1)} \right) - \ln \left( \frac{\int d\mathbf{s}^{N+1} \exp[-\beta U(\mathbf{s}^{N+1}; L)]}{\int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]} \right)\end{aligned}$$

$$\beta\mu = \beta\mu^{IG} + \beta\mu^{ex}$$

$$\beta\mu^{ex} = -\ln \left( \frac{\int d\mathbf{s}^{N+1} \exp[-\beta U(\mathbf{s}^{N+1}; L)]}{\int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)]} \right)$$

# *Widom test particle insertion*

$$\beta\mu^{ex} = -\ln \left( \frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$U(s^{N+1}; L) = \Delta U^+ + U(s^N; L)$$

$$\beta\mu^{ex} = -\ln \left( \frac{\int ds^N \int ds_{N+1} \exp[-\beta(\Delta U^+ + U(s^N; L))] }{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$= -\ln \left( \frac{\int ds_{N+1} \int ds^N \{ \exp[-\beta \Delta U^+] \} \exp[-\beta U(s^N; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$= -\ln \left( \int ds_{N+1} \langle \exp[-\beta \Delta U^+] \rangle_{NVT} \right)$$

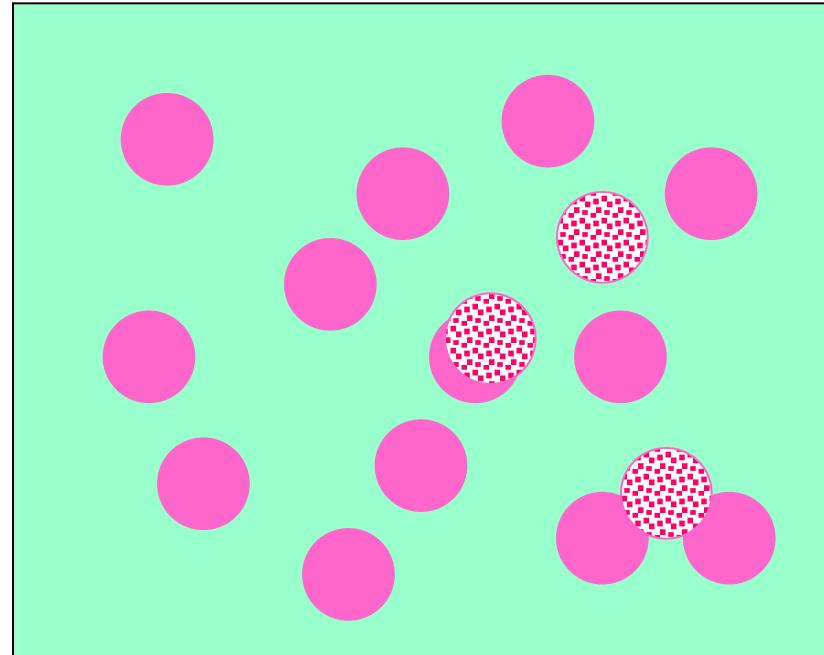
Ghost particle!

# *Hard spheres*

$$\beta\mu^{ex} = -\ln \left( \int d\mathbf{s}_{N+1} \left\langle \exp[-\beta\Delta U^+] \right\rangle_{NVT} \right)$$

$$U(r) = \begin{cases} \infty & r \leq \sigma \\ 0 & r > \sigma \end{cases}$$

$$\exp[-\beta\Delta U^+] = \begin{cases} 0 & \text{if overlap} \\ 1 & \text{no overlap} \end{cases}$$



$\left\langle \exp[-\beta\Delta U^+] \right\rangle$  probability to insert a test particle!

But, ... may fail at high density

# *Thermodynamic perturbation – Umbrella Sampling*

Two systems:

System 0:  $N, V, T, U_0$

System 1:  $N, V, T, U_1$

$$Q_0 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_0)$$

$$Q_1 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_1)$$

$$\Delta\beta F = \beta F_1 - \beta F_0 = -\ln(Q_1/Q_0)$$

$$= -\ln \frac{\int d\mathbf{s}^N \exp[-\beta U_1]}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$= -\ln \frac{\int d\mathbf{s}^N \exp[-\beta(U_1 - U_0)] \exp[-\beta U_0]}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$\Delta\beta F = -\ln \langle \exp[-\beta(U_1 - U_0)] \rangle_0$$

# *Umbrella sampling*

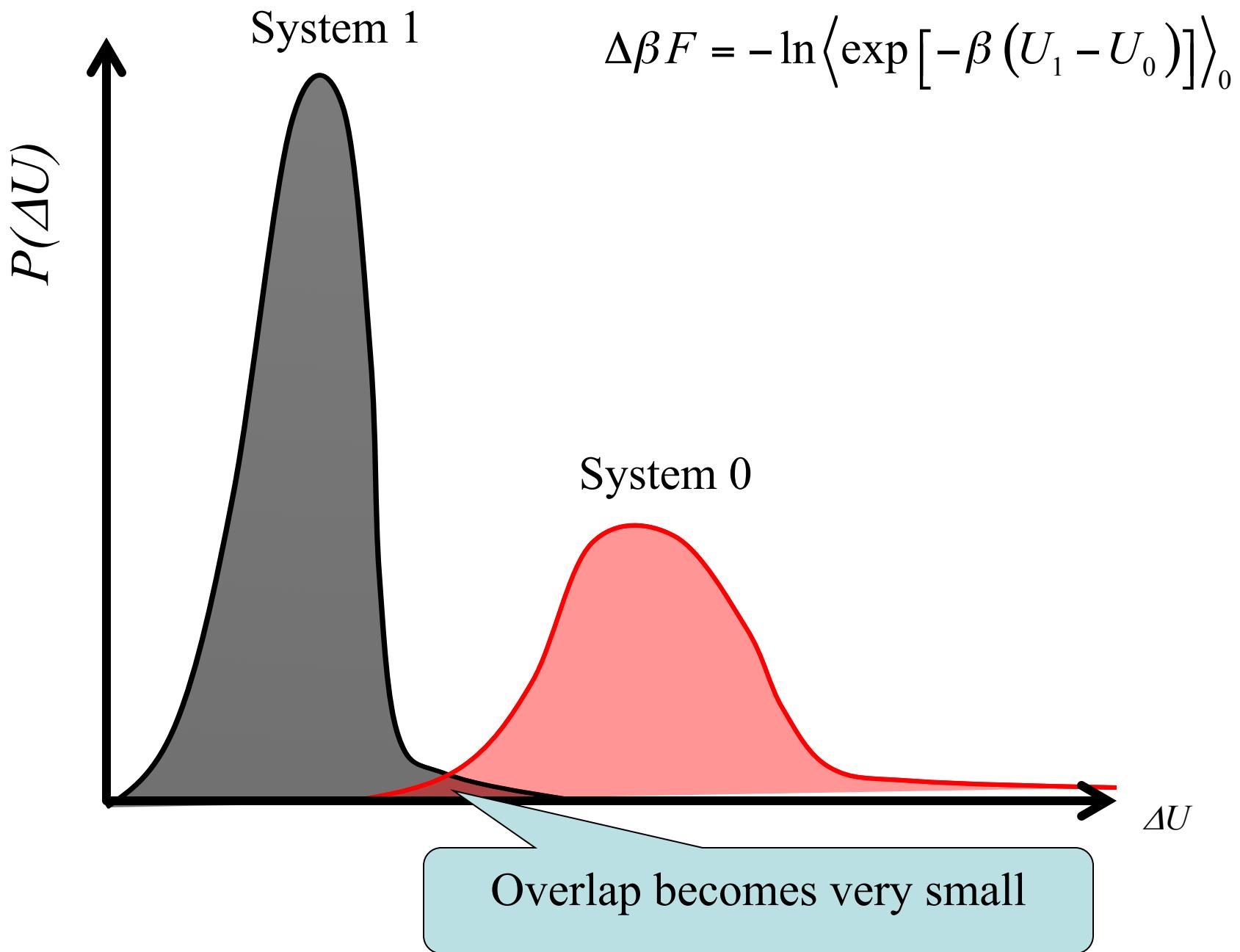
- Start with thermodynamic perturbation

$$\Delta\beta F = -\ln(Q_1/Q_0) = -\ln\left(\frac{\int d\mathbf{s}^N \exp(-\beta U_1)}{\int d\mathbf{s}^N \exp(-\beta U_0)}\right)$$

$$\exp(-\beta\Delta F) = \left( \frac{\int d\mathbf{s}^N \exp(-\beta U_0) \exp(-\beta\Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_0)} \right)$$

$$\exp(-\beta\Delta F) = \langle \exp(-\beta\Delta U) \rangle_0$$

Can we use this for free energy difference between arbitrary systems?



# *Bridging function*

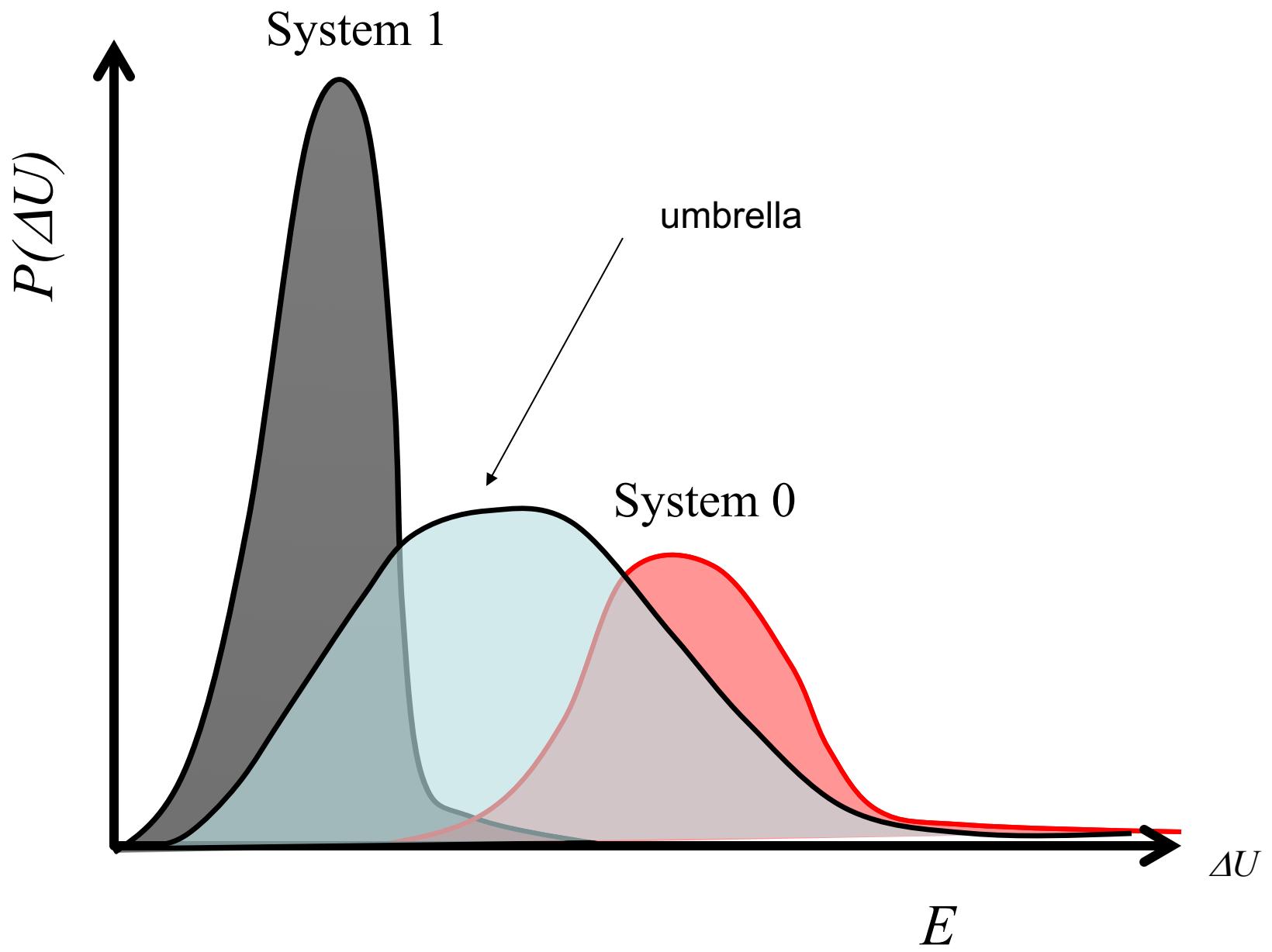
- Introduce function  $\pi(\mathbf{s}^N)$  altering distribution.

$$\exp(-\beta\Delta F) = \left( \frac{\int d\mathbf{s}^N \pi(\mathbf{s}^N) \exp(-\beta U_1) / \pi(\mathbf{s}^N)}{\int d\mathbf{s}^N \pi(\mathbf{s}^N) \exp(-\beta U_0) / \pi(\mathbf{s}^N)} \right)$$

$$\exp(-\beta\Delta F) = \langle \exp(-\beta\Delta U) \rangle_0$$

$$= \frac{\langle \exp(-\beta U_1) / \pi \rangle_\pi}{\langle \exp(-\beta U_0) / \pi \rangle_\pi}$$

- This approach is called umbrella sampling



### III - Overlapping Distribution Method

Two systems:

System 0:  $N, VT, U_0$

System 1:  $N, VT, U_1$

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$$Q_0 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_0)$$

$$Q_1 = \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp(-\beta U_1)$$

$$\Delta\beta F = \beta F_1 - \beta F_0 = -\ln(Q_1/Q_0) = -\ln\left(\frac{\int d\mathbf{s}^N \exp(-\beta U_1)}{\int d\mathbf{s}^N \exp(-\beta U_0)}\right) = -\ln\left(\frac{Q_1}{Q_0}\right)$$

$= \Delta U$  ( $\delta$  function)

$$p_0(\Delta U) = \frac{\int d\mathbf{s}^N \exp(-\beta U_0) \delta(U_1 - U_0 - \Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$p_1(\Delta U) = \frac{\int d\mathbf{s}^N \exp(-\beta U_1) \delta(U_1 - U_0 - \Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_1)}$$

$$p_1(\Delta U) = \frac{\int d\mathbf{s}^N \exp[-\beta(U_1 - U_0)] \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_1)}$$

$$\frac{Q_0}{Q_1} = \exp(\beta \Delta F) = \frac{Q_0}{Q_1} \exp(-\beta \Delta U) \frac{\int d\mathbf{s}^N \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{Q_0}$$

$$= \frac{1}{Q_1} = \frac{Q_0}{Q_1} \frac{1}{Q_0}$$

$$p_1(\Delta U) = \frac{Q_0}{Q_1} \exp(-\beta \Delta U) p_0(\Delta U)$$

$$\ln p_1(\Delta U) = \beta(\Delta F - \Delta U) + \ln p_0(\Delta U)$$

# Overlapping Distribution Method

$$\ln p_1(\Delta U) = \beta(\Delta F - \Delta U) + \ln p_0(\Delta U)$$

$$f_0(\Delta U) \equiv \ln p_0(\Delta U) - 0.5\beta\Delta U$$

$$f_1(\Delta U) \equiv \ln p_1(\Delta U) + 0.5\beta\Delta U$$

Simulate system 0: compute  $f_0$   
Simulate system 1: compute  $f_1$

$$\beta\Delta F = f_1(\Delta U) - f_0(\Delta U)$$

# **Chemical potential (LJ fluid)**

System 0:  $N-1, V, T, U + 1$  ideal gas

$$\Delta\beta F = \beta F_1 - \beta F_0 \equiv \beta\mu^{ex}$$

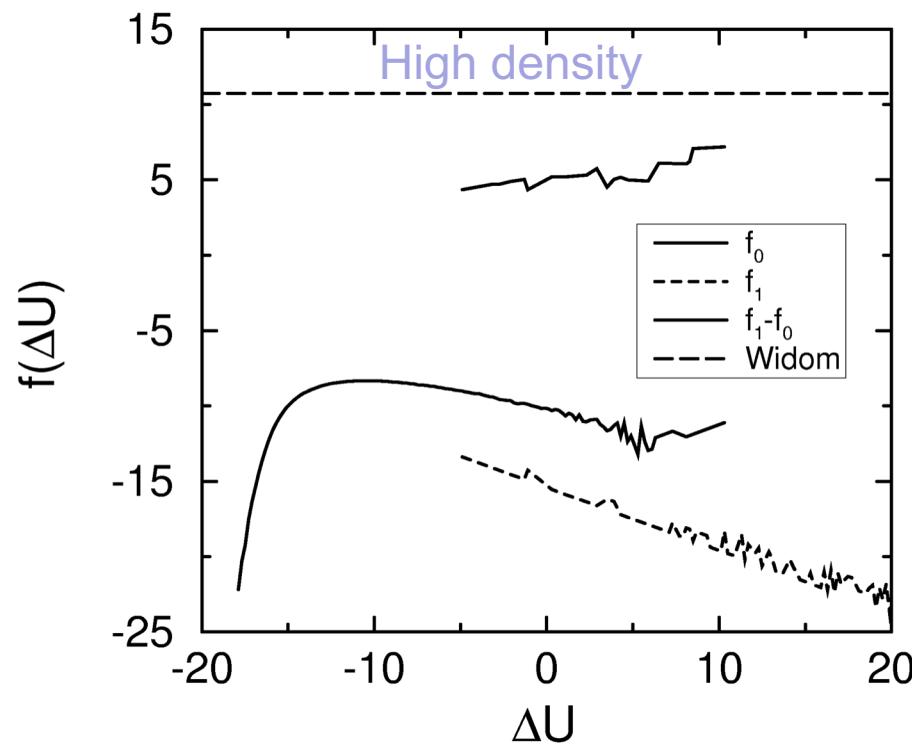
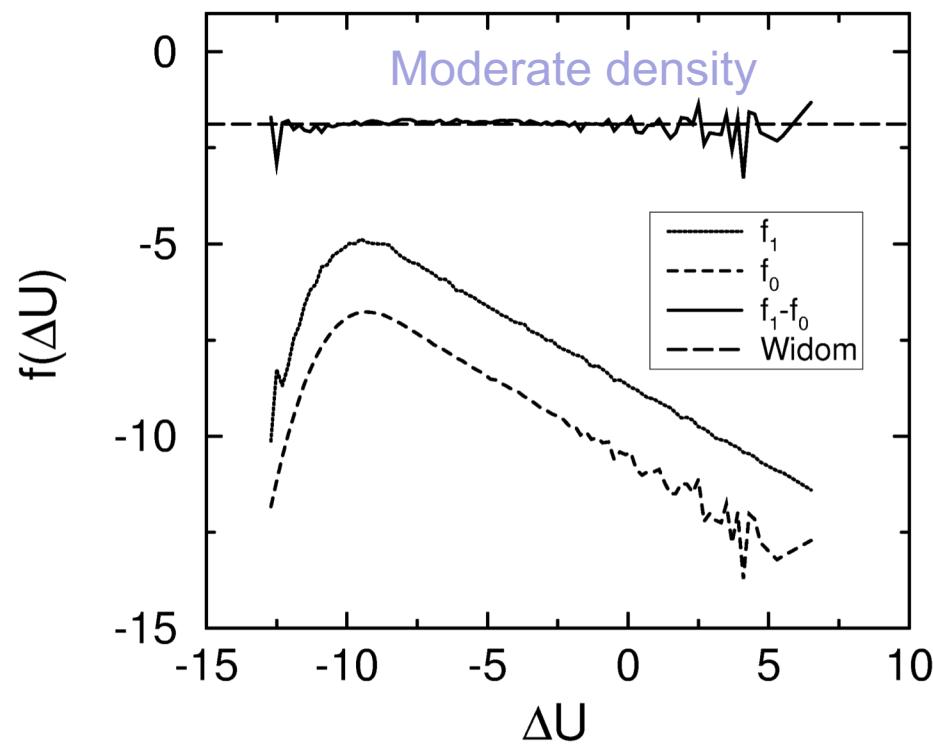
System 0: test particle energy

$$\beta\mu^{ex} = f_1(\Delta U) - f_0(\Delta U)$$

System 1:  $N, V, T, U$

$$\Delta U = U_1 - U_0$$

System 1: real particle energy



## IV - Non-Boltzmann sampling

$T_1$  is arbitrary!

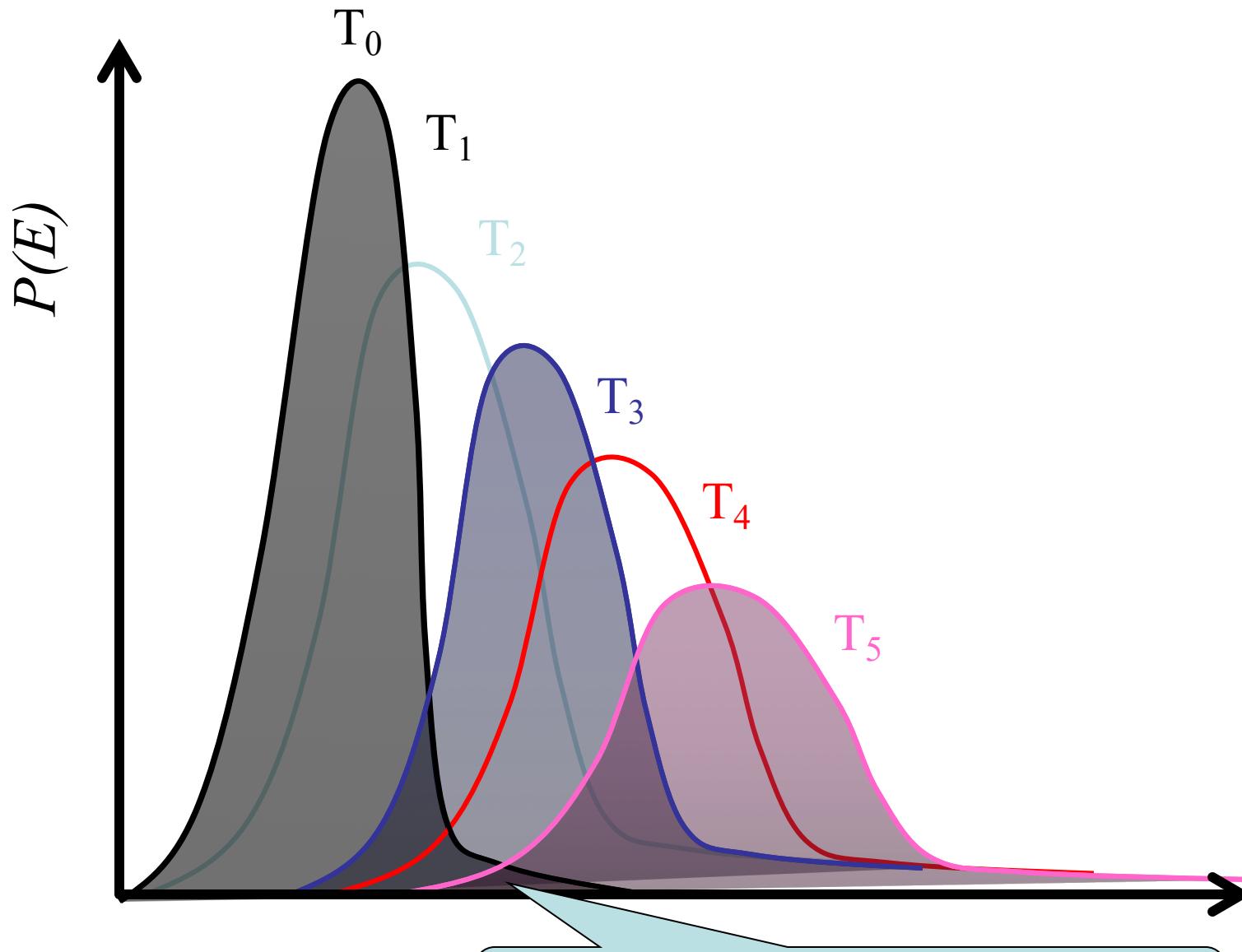
$$\langle A \rangle_{NVT_1} = \frac{1}{Q_{NVT_1}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta_1 U(\mathbf{r}^N)]$$

$$\begin{aligned}
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta_1 U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta_1 U(\mathbf{r}^N)]} \\
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta_1 U(\mathbf{r}^N)] \exp[\beta_2 U(\mathbf{r}^N) - \beta_2 U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta_1 U(\mathbf{r}^N)] \exp[\beta_2 U(\mathbf{r}^N) - \beta_2 U(\mathbf{r}^N)]} \\
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[\beta_2 U(\mathbf{r}^N) - \beta_1 U(\mathbf{r}^N)] \exp[-\beta_2 U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[\beta_2 U(\mathbf{r}^N) - \beta_1 U(\mathbf{r}^N)] \exp[-\beta_2 U(\mathbf{r}^N)]}
 \end{aligned}$$

We only  
need a  
*single*  
simulation!

$$= \frac{\langle A \exp[(\beta_2 - \beta_1)U] \rangle_{NVT_2}}{\langle \exp[(\beta_2 - \beta_1)U] \rangle_{NVT_2}}$$

We perform a simulation at  $T=T_2$  and  
we determine  $A$  at  $T=T_1$



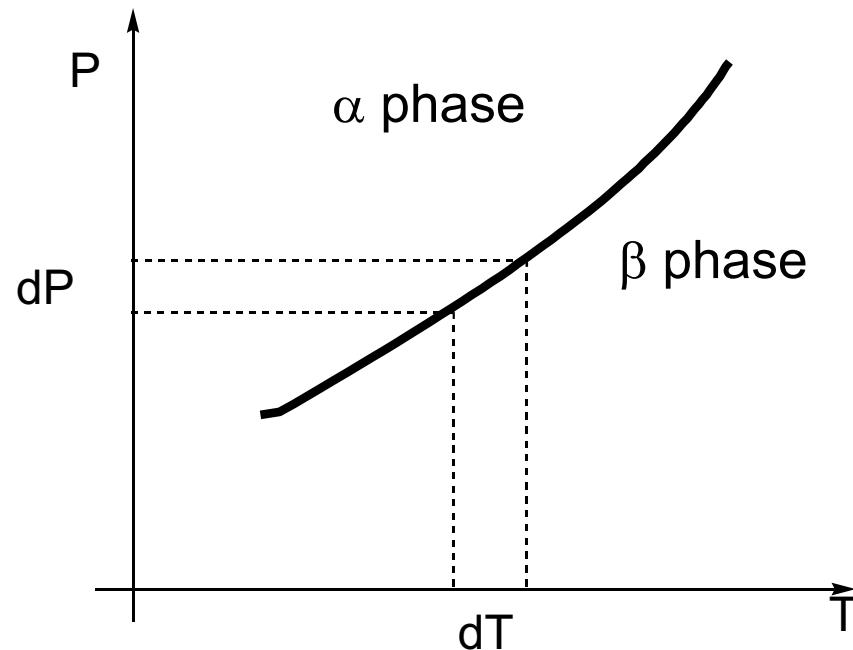
Overlap becomes very small

# Tracing coexistence curves

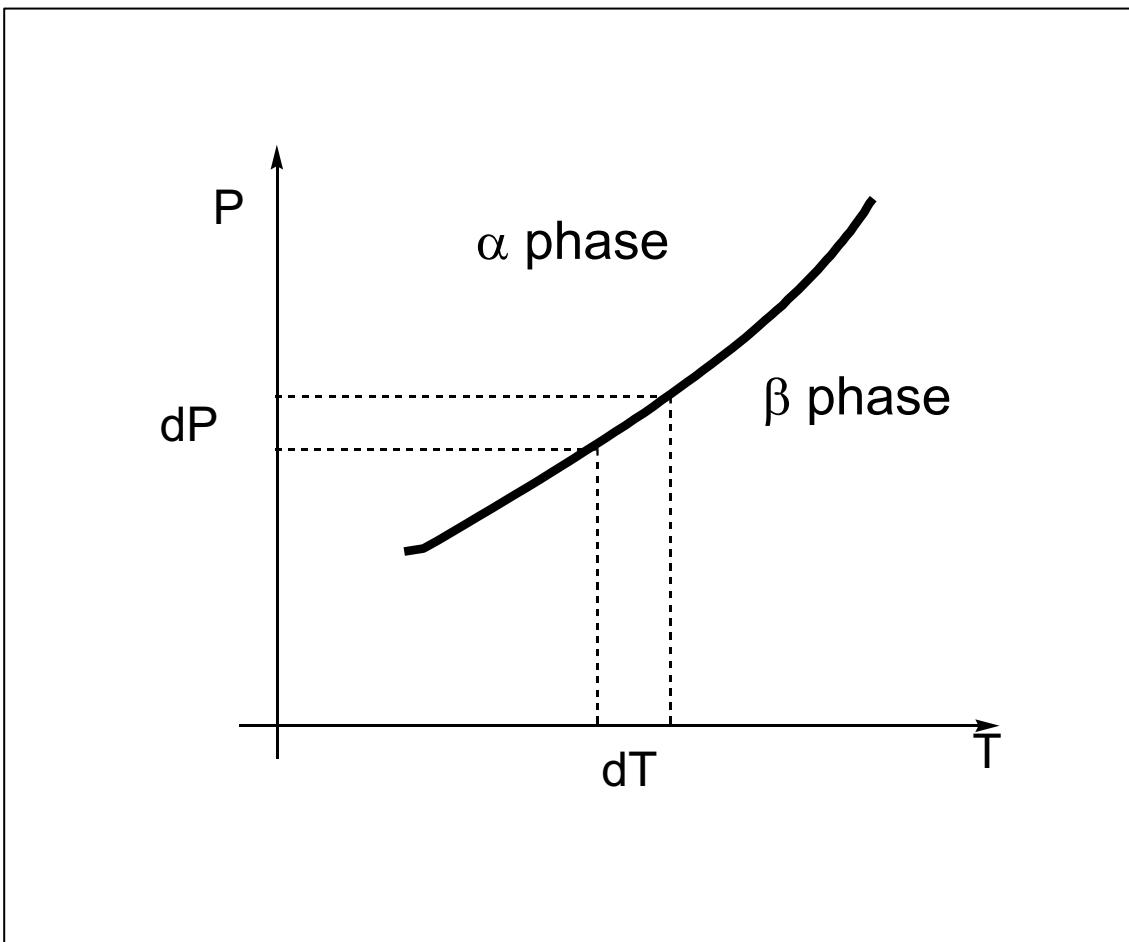
- If we have a coexistence point on the phase diagram we can integrate allong the line while maintaining coexistence.

P en T are equal along  
coexistence line

$$d\mu_\alpha = d\mu_\beta$$



# Tracing coexistence curves



Clapeyron equation

$$\frac{dP}{dT} = \frac{\Delta(U + PV)}{T\Delta V}$$

$$dP = \frac{\Delta(U + PV)}{T\Delta V} dT$$