A blue spiral-bound notebook with silver rings at the top. The title 'Free Energy Calculation' is written in white on a light blue rectangular background.

Free Energy Calculation

Outline

- ❑ Free energies: classical thermodynamics
- ❑ Free energies: statistical thermodynamics
- ❑ Monte Carlo simulations: what went wrong?
- ❑ Experiments: how to make a chemical potential-meter?
- ❑ Free energy techniques:
 - ❑ Thermodynamic integration
 - ❑ Krikwood coupling parameter
 - ❑ Widom test particle insertion
 - ❑ Overlapping distribution
 - ❑ Histogram method
 - ❑ Umbrella sampling

Chemical potential

Until now we have assumed that most of our systems are closed: NVE

Let us now assume that our system can exchange matter

First law

$$dU = TdS - pdV + \sum_i \left(\frac{\partial U}{\partial n_i} \right)_{S, V, n_j} dn_i$$

Change in U if we add one mole of component i at constant S, V, n_j

Chemical potential

$$\mu_i \equiv \left(\frac{\partial U}{\partial n_i} \right)_{S, V, n_j}$$

How is the chemical potential is defined in term of A,G,H

$$A \equiv U - TS$$

$$dA = dU - TdS - SdT$$

Using

$$dU = TdS - pdV + \sum_i \left(\frac{\partial U}{\partial n_i} \right)_{S,V,n_j} dn_i$$

$$dA = -SdT - pdV + \sum_i \left(\frac{\partial U}{\partial n_i} \right)_{S,V,n_j} dn_i$$

$$dA = -SdT - pdV + \sum_i \left(\frac{\partial A}{\partial n_i} \right)_{T,V,n_j} dn_i$$

Chemical potential

$$\mu_i \equiv \left(\frac{\partial A}{\partial n_i} \right)_{T,V,n_j} = \left(\frac{\partial U}{\partial n_i} \right)_{S,V,n_j}$$

Energy

$$dU = TdS - pdV + \sum_i \mu_i dn_i$$

Enthalpy

$$dH = TdS + Vdp + \sum_i \mu_i dn_i$$

Helmholtz free energy

$$dA = -SdT - pdV + \sum_i \mu_i dn_i$$

Gibbs free energy

$$dG = -SdT + Vdp + \sum_i \mu_i dn_i$$

Chemical potential

$$\mu_i \equiv \left(\frac{\partial U}{\partial n_i} \right)_{S, V, n_j} = \left(\frac{\partial H}{\partial n_i} \right)_{S, p, n_j} = \left(\frac{\partial A}{\partial n_i} \right)_{T, V, n_j} = \left(\frac{\partial G}{\partial n_i} \right)_{T, p, n_j}$$

G is extensive

How does the Gibbs free energy change if the system is increased by a factor k?

$$\Delta G = kG - G = (k - 1)G$$

$$dG = -SdT + Vdp + \sum_i \mu_i dn_i$$

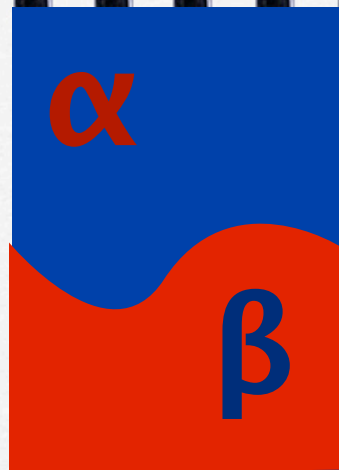
Integration gives

$$\Delta G = 0 + 0 + \sum_i \mu_i \Delta n_i$$

$$\Delta G = 0 + 0 + \sum_i \mu_i (k - 1)n_i$$

Which gives:

$$G = \sum_i \mu_i n_i$$



Equilibrium

Let us consider a system with a constant number of particles, volume, and energy.

This system consists of two phases α and β

Question: When are these two systems in equilibrium?

$$dS = dS^\alpha + dS^\beta = 0 \quad dS^\alpha = \frac{dU^\alpha}{T^\alpha} + \frac{p^\alpha}{T^\alpha} dV^\alpha - \sum_i \frac{\mu_i^\alpha}{T^\alpha} dn_i^\alpha$$

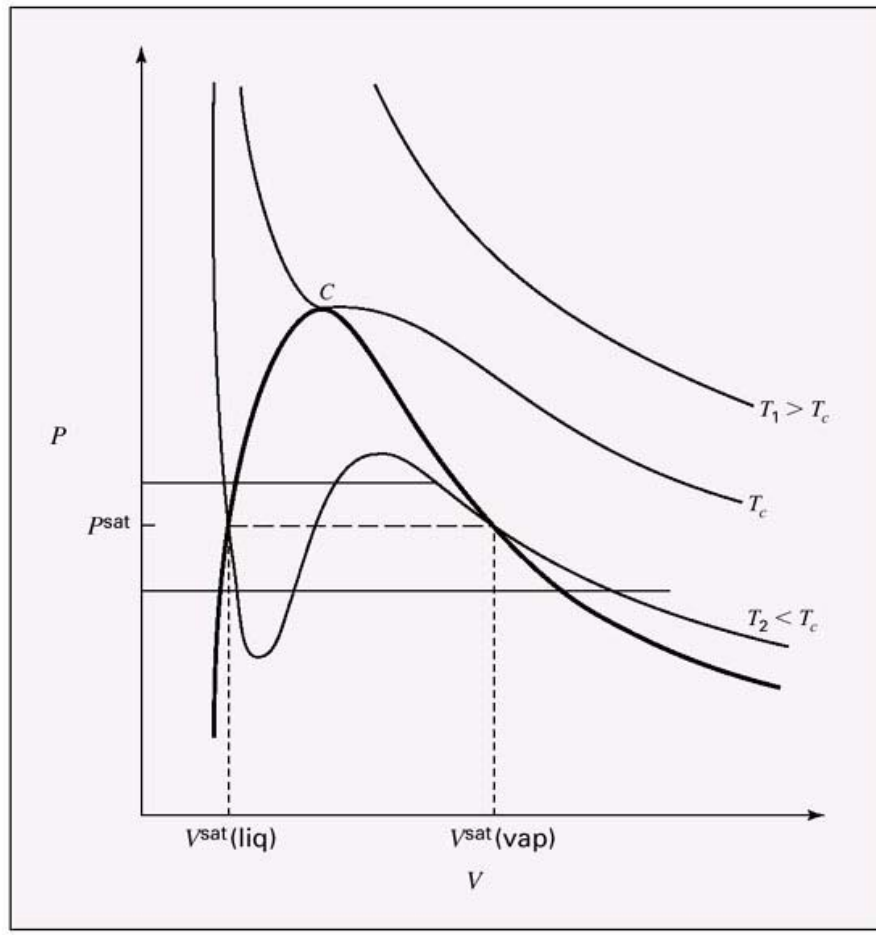
$$dS^\alpha + dS^\beta = \left(\frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) dU^\alpha + \left(\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) dV^\alpha - \sum_i \left(\frac{\mu_i^\alpha}{T^\alpha} - \frac{\mu_i^\beta}{T^\beta} \right) dn_i^\alpha$$

Equilibrium:

$$T^\alpha = T^\beta \wedge p^\alpha = p^\beta \wedge \mu_i^\alpha = \mu_i^\beta$$

Vapor-liquid equilibria

Equilibrium: $T^\alpha = T^\beta \wedge p^\alpha = p^\beta \wedge \mu^\alpha = \mu^\beta$



Van der Waals equation of

$$P = \frac{k_B T}{v - b} - \frac{a}{v^2}$$

We need to compute the chemical potential

Chemical potential

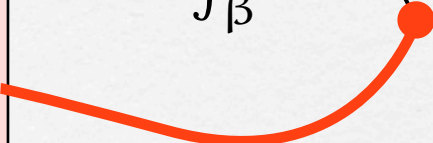
In equilibrium we have: $\mu^\alpha(T, V^\alpha) = \mu^\beta(T, V^\beta)$

$$\Delta\mu^{\alpha\beta} = \mu^\alpha(T, V^\alpha) - \mu^\beta(T, V^\beta)$$

Equilibrium if:

$$\Delta\mu^{\alpha\beta} = \int_{\beta}^{\alpha} dV \left(\frac{\partial \mu}{\partial V} \right)_{T,n} = 0$$

How to relate this to the equation of state?



Recall that we derived the relation between changes in the chemical potential and temperature and pressure (Gibbs-Duhem)

$$G = \sum_i \mu_i n_i \quad \text{or, for a pure component} \quad G = n\mu$$

$$\sum_i n_i d\mu_i = -SdT + Vdp \quad \text{or} \quad \begin{aligned} nd\mu &= -SdT + Vdp \\ d\mu &= -sdT + vdp \end{aligned}$$

$$\Delta\mu^{\alpha\beta} = \int_{\beta}^{\alpha} dv \left(\frac{\partial\mu}{\partial v} \right)_{T,n} = 0$$

$$d\mu = -sdT + vdp$$

$$\mu = \mu(T, P)$$

We need to find

$$\left(\frac{\partial\mu}{\partial v} \right)_T$$

$$\left(\frac{\partial\mu}{\partial v} \right)_T = \left(\frac{\partial\mu}{\partial T} \right)_p \underbrace{\left(\frac{\partial T}{\partial v} \right)_T}_{=0} + \left(\frac{\partial\mu}{\partial p} \right)_T \underbrace{\left(\frac{\partial p}{\partial v} \right)_T}_{=v}$$

$$\left(\frac{\partial\mu}{\partial v} \right)_T = 0 + v \left(\frac{\partial p}{\partial v} \right)_T$$

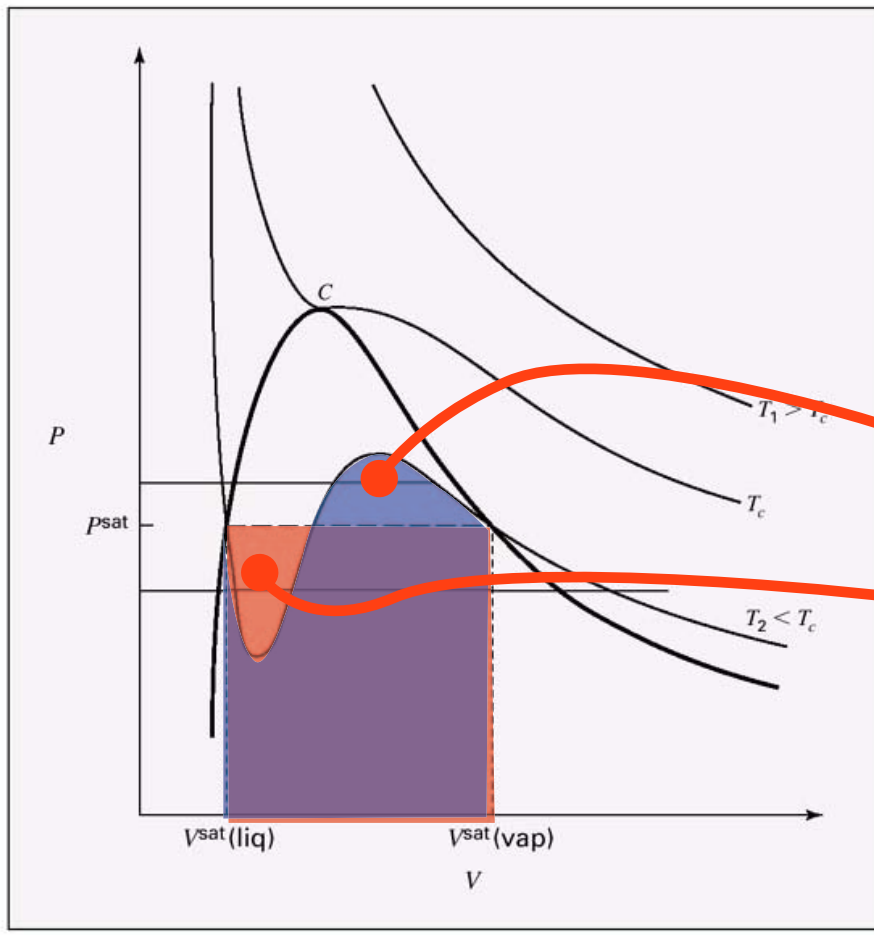
Equation of state gives us the chemical potential

Which gives:

$$\begin{aligned} \Delta\mu^{\alpha\beta} &= \int_{\beta}^{\alpha} dv v \left(\frac{\partial p}{\partial v} \right)_T = vp \Big|_{\alpha}^{\beta} - \int_{\beta}^{\alpha} p dv \\ &= - \int_{\beta}^{\alpha} (p - p_{\text{coex}}) dv \end{aligned}$$

Vapor-liquid equilibria

Equilibrium: $T^\alpha = T^\beta \wedge p^\alpha = p^\beta \wedge \mu^\alpha = \mu^\beta$



Equal chemical potential if

$$\int_{\beta}^{\alpha} (p - p_{\text{coex}}) dv = 0$$

Equal chemical potential if these areas are equal

NVT ensemble

Define de Broglie wave length:

$$\Lambda \equiv \left(\frac{h^2 \beta}{2\pi m} \right)^{1/2}$$

Partition function:

$$Q(N, V, T) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp \left[-\beta U(\mathbf{r}^N) \right]$$

Free energy:

$$F = -\frac{1}{\beta} \ln Q(N, V, T)$$

Example: ideal gas

$$\begin{aligned} Q(N, V, T) &= \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)] \\ &= \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N 1 = \frac{V^N}{\Lambda^{3N} N!} \end{aligned}$$

Free energy:

$$\begin{aligned} \beta F &= -\ln \left(\frac{V^N}{\Lambda^{3N} N!} \right) \\ &\approx N \ln \Lambda^3 + N \ln \left(\frac{N}{V} \right) = N \ln \Lambda^3 + N \ln \rho \end{aligned}$$

Thermo recall (3)

Helmholtz Free energy:

$$dF = -SdT - pdV$$

Pressure

$$\left(\frac{\partial F}{\partial V} \right)_T = -P$$

Energy:

$$\left(\frac{\partial F/T}{\partial 1/T} \right) = \left(\frac{\partial \beta F}{\partial \beta} \right) = E$$

Free energy:

$$\beta F = N \ln \Lambda^3 + N \ln \left(\frac{N}{V} \right)$$

Pressure:

$$P = - \left(\frac{\partial F}{\partial V} \right)_T = \frac{N}{\beta V}$$

Energy:

$$E = \left(\frac{\partial \beta F}{\partial \beta} \right) = \frac{3N}{\Lambda} \frac{\partial \Lambda}{\partial \beta} = \frac{3}{2} N k_B T$$

Chemical potential:

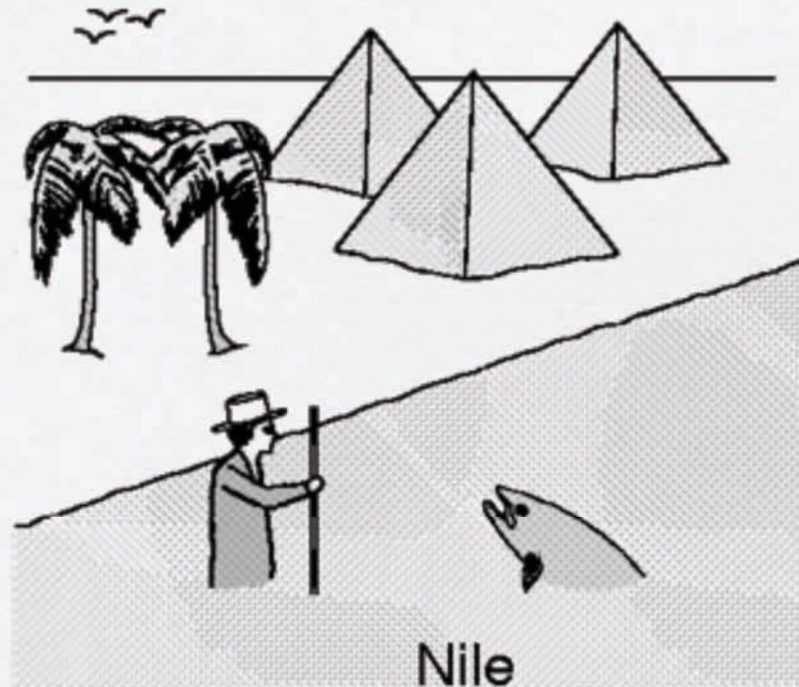
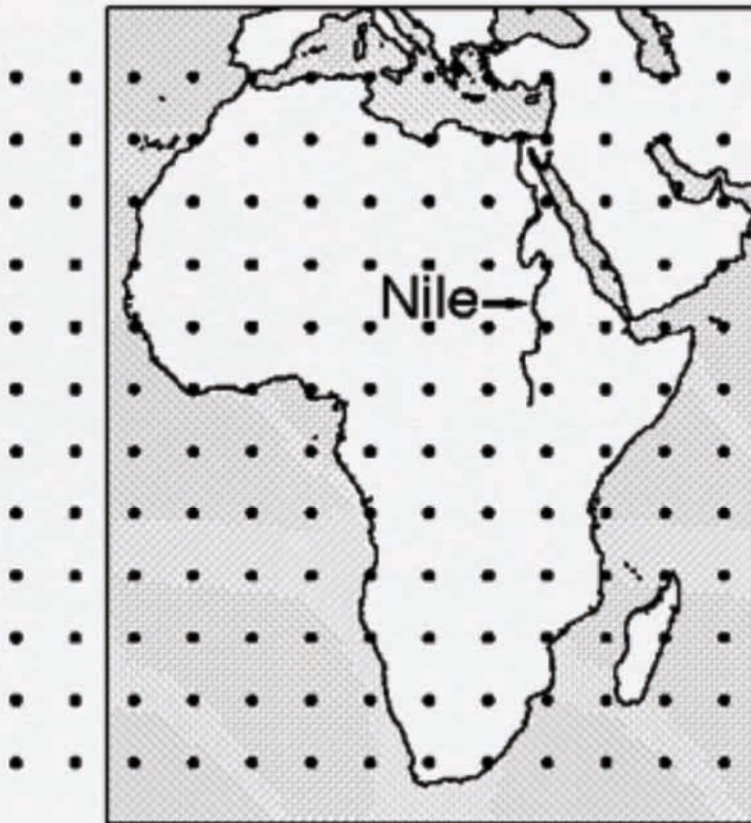
$$\mu_i = \left(\frac{\partial F}{\partial N_i} \right)_{T, V, N_j}$$

$$\beta \mu^{IG} = \ln \Lambda^3 + \ln \rho + 1$$

$$\beta \mu^{IG} = \beta \mu^0 + \ln \rho$$

Monte Carlo simulation

What is the difference between $\langle A \rangle$ and F ?



The price of importance sampling

$$\begin{aligned}\langle A \rangle_{NVT} &= \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)] \\ &= \int d\mathbf{r}^N A(\mathbf{r}^N) P(\mathbf{r}^N) = \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) P(\mathbf{r}^N)}{\int d\mathbf{r}^N P(\mathbf{r}^N)}\end{aligned}$$

$$P(\mathbf{r}^N) = \frac{\exp[-\beta U(\mathbf{r}^N)]}{Q_{NVT} \Lambda^{3N} N!}$$

$$= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) C \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N C \exp[-\beta U(\mathbf{r}^N)]} = \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]}$$

Generate configuration using MC:

$$\{\mathbf{r}_1^N, \mathbf{r}_2^N, \mathbf{r}_3^N, \mathbf{r}_4^N, \dots, \mathbf{r}_M^N\} \quad \bar{A} = \frac{1}{M} \sum_{i=1}^M A(\mathbf{r}_i^N) = \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) P^{MC}(\mathbf{r}^N)}{\int d\mathbf{r}^N P^{MC}(\mathbf{r}^N)}$$

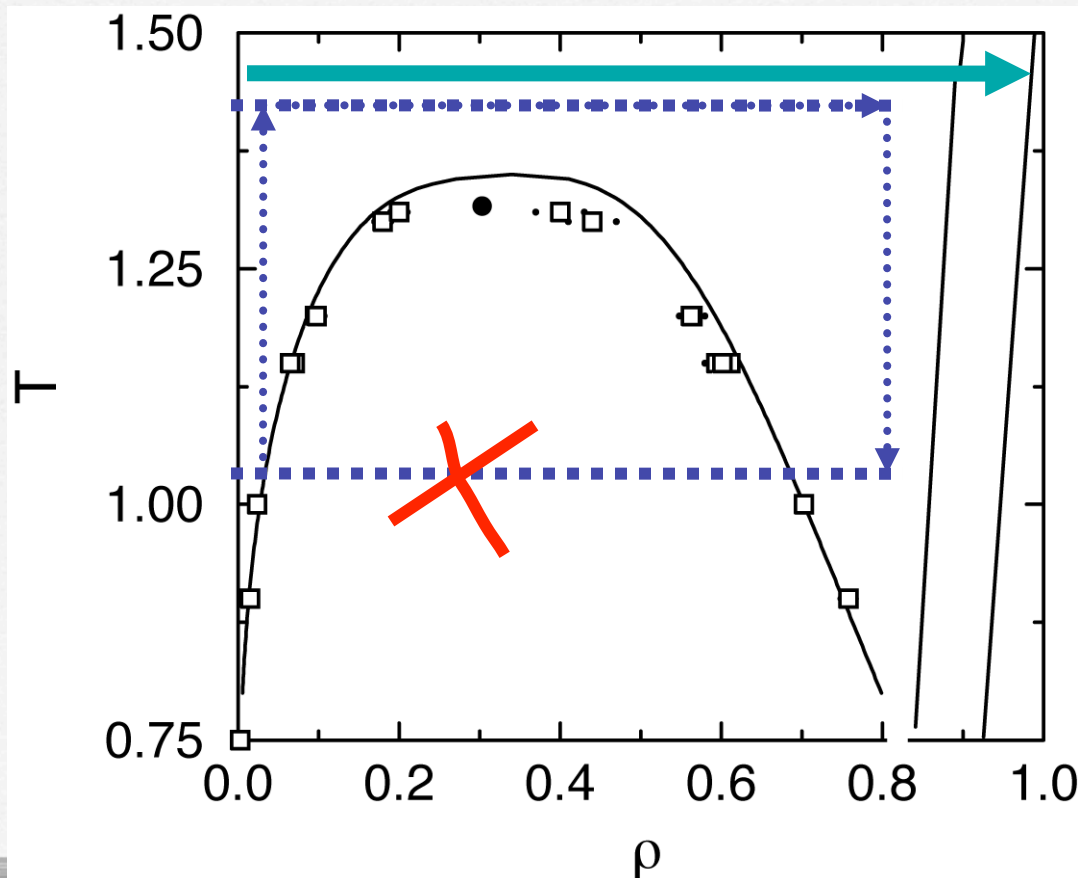
with

$$P^{MC}(\mathbf{r}^N) = C^{MC} \exp[-\beta U(\mathbf{r}^N)]$$

$$\begin{aligned}&= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) C^{MC} \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N C^{MC} \exp[-\beta U(\mathbf{r}^N)]} \\ &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]} \neq\end{aligned}$$

Experimental (2)

$$\left(\frac{\partial F}{\partial V}\right)_{N,T} = -P \quad F(P) - F(P_0) = -\int_{V_0}^V \left(\frac{\partial F}{\partial V}\right)_{N,T} dV$$



?

- Works always for vapor-liquid
- Requires a large number of simulations to fit the equation of state
- Does not work for solid-liquid

Thermodynamic integration

Coupling parameter

$$U(\lambda) = (1 - \lambda)U_I + \lambda U_{II}$$

$$U(\lambda) = U^{LJ} + \lambda U^{\text{dipole-dipole}}$$

$$U(0) = U^{LJ} \quad \text{Lennard-Jones}$$

$$U(1) = U^{\text{Stockm}} \quad \text{Stockmayer}$$

$$Q_{NVT}(\lambda) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\lambda)]$$

$$Q_{NVT}(\lambda) = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\lambda)]$$

$$\left(\frac{\partial F(\lambda)}{\partial \lambda} \right)_{N,T} = -\frac{1}{\beta} \frac{\partial}{\partial \lambda} \ln(Q) = -\frac{1}{\beta} \frac{1}{Q} \frac{\partial Q}{\partial \lambda}$$

$$= \frac{\int d\mathbf{r}^N (\partial U(\lambda)/\partial \lambda) \exp[-\beta U(\lambda)]}{\int d\mathbf{r}^N \exp[-\beta U(\lambda)]} = \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_{\lambda}$$

Free energy as
ensemble average!

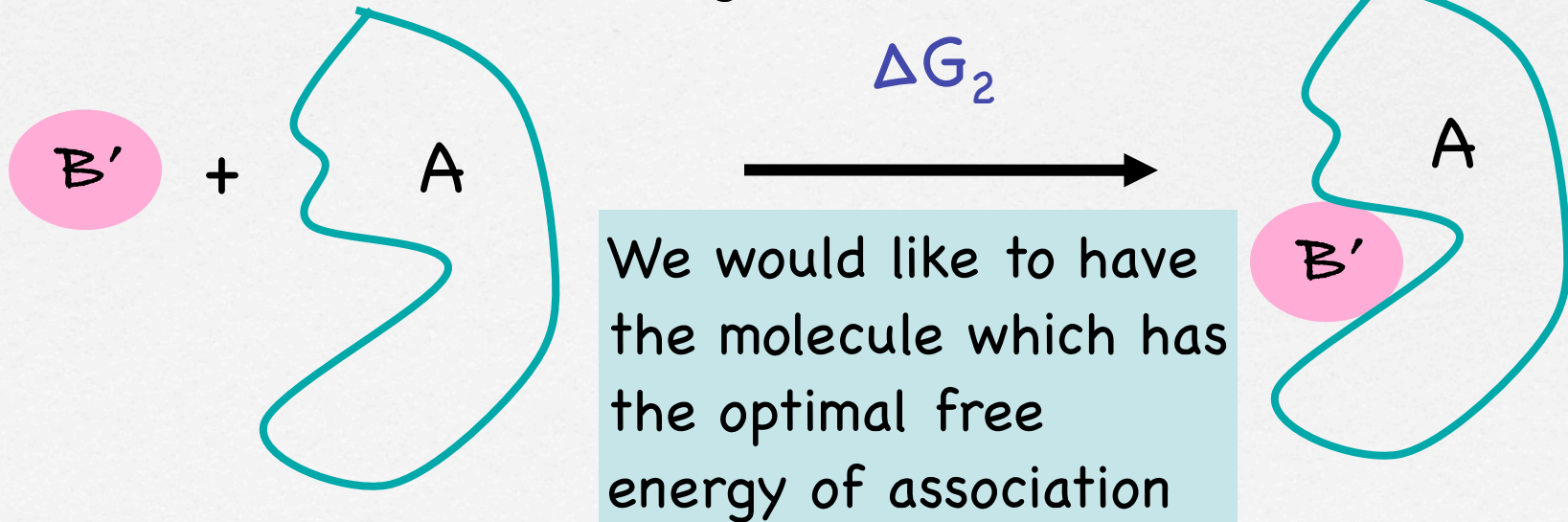
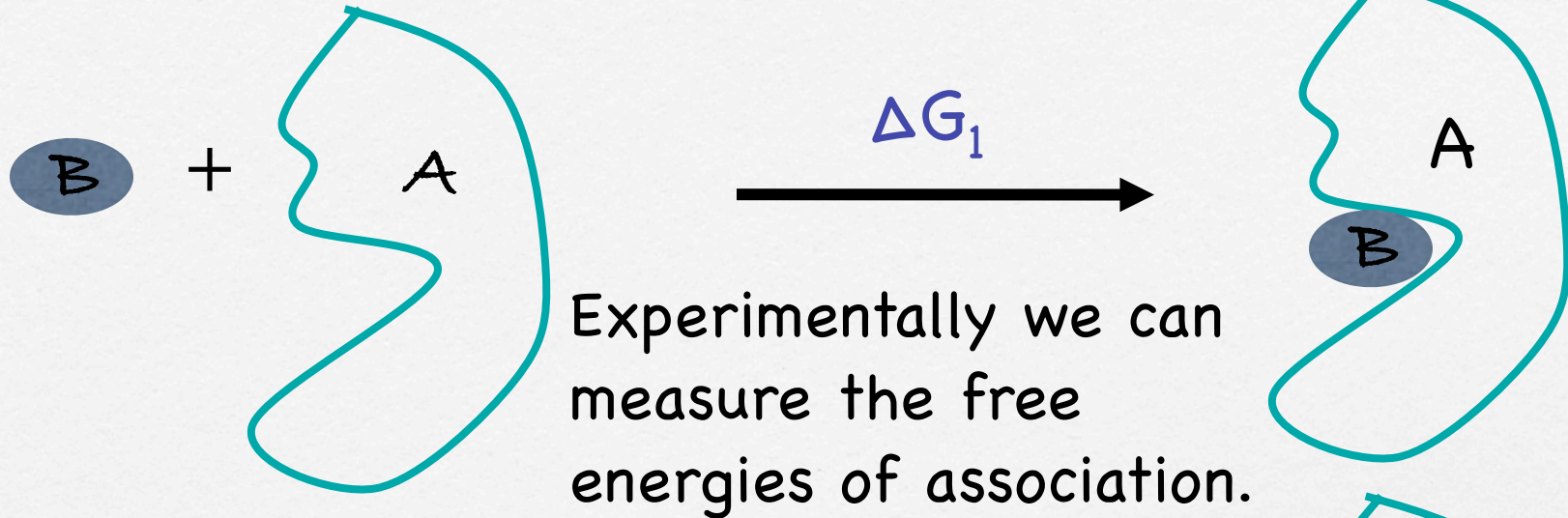
$$F(\lambda=1) - F(\lambda=0) = \int d\lambda \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_{\lambda}$$

$$F = -\frac{1}{\beta} \ln(Q_{NVT})$$

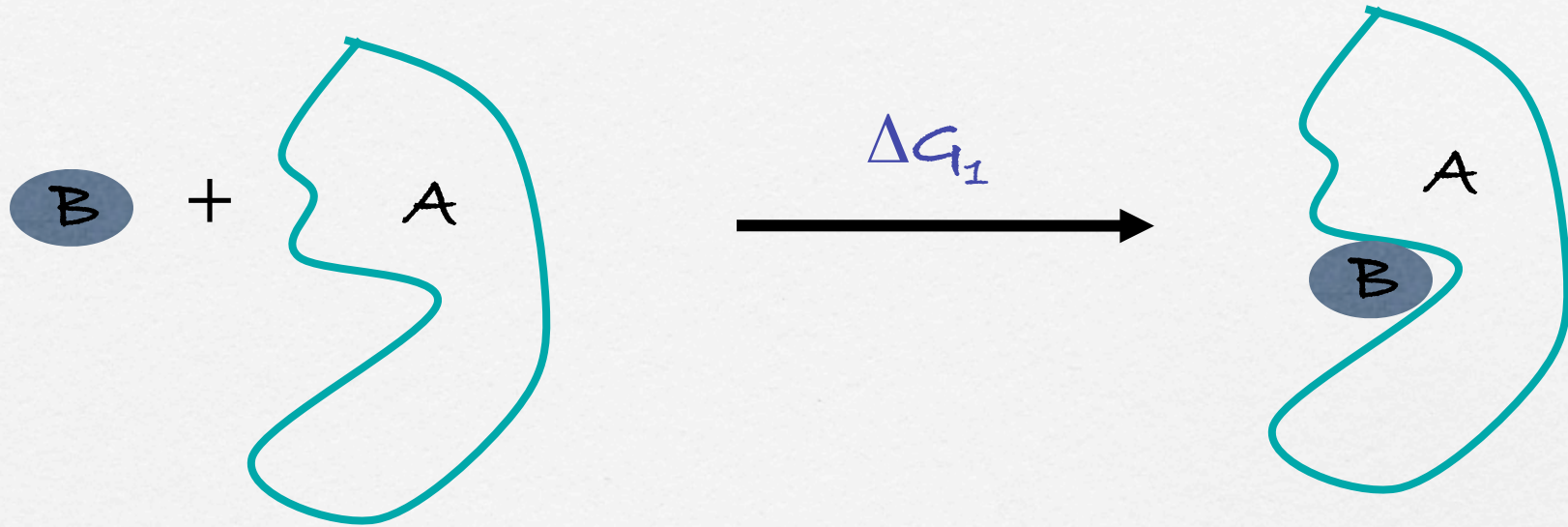
$$U(\lambda) = (1-\lambda)U_I + \lambda U_{II}$$

$$\left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_{\lambda} = \langle U_{II} - U_I \rangle_{\lambda}$$

Molecular association



How to compute the free energy of association?



Solution: compute the **potential of mean force** =
the reversible work required to bring B to the host

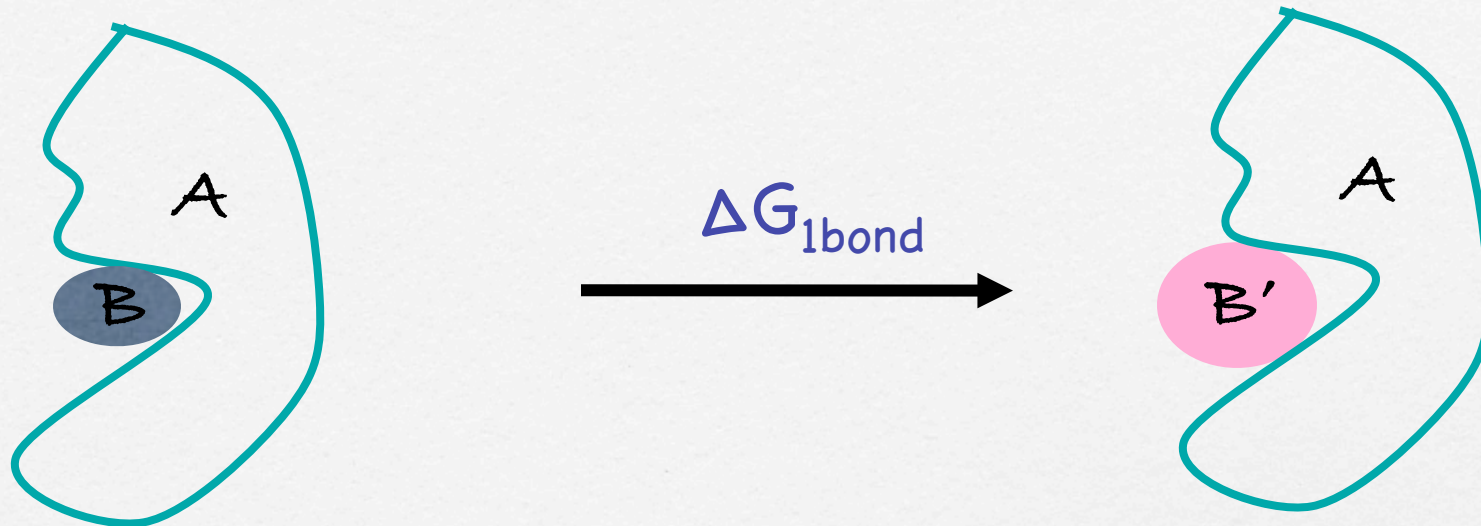


We do not need to know the absolute value, but to see whether B or B' is better we need to know only:

$$\Delta\Delta G = \Delta G_1 - \Delta G_2 \quad U(\lambda) = U(B) + \lambda[U(B') - U(B)]$$

In a simulation it is easier to compute:

$$\Delta\Delta G = \Delta G_{\text{bind}} - \Delta G_{\text{solv}}$$



$$U(\lambda) = U(B) + \lambda[U(B') - U(B)]$$

$$F(\lambda = 1) - F(\lambda = 0) = \int d\lambda \left\langle \frac{\partial U(\lambda)}{\partial \lambda} \right\rangle_{\lambda}$$

When will this computation be accurate?

Chemical potential

$$Q_{NVT} = \frac{1}{\Lambda^{3N} N!} \int dr^N \exp[-\beta U(r^N)]$$

Scaled coordinates: $s=r/L$

$$Q_{NVT} = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp[-\beta U(s^N; L)]$$

$$\beta F = -\ln(Q_{NVT})$$

$$= -\ln\left(\frac{V^N}{\Lambda^{3N} N!}\right) - \ln\left(\int ds^N \exp[-\beta U(s^N; L)]\right)$$

$$= -N \ln\left(\frac{1}{\Lambda^3 \rho}\right) - \ln\left(\int ds^N \exp[-\beta U(s^N; L)]\right)$$

Chemical potential: Widom test particle method

$$\beta F = -\ln(Q_{NVT})$$

$$= -N \ln\left(\frac{1}{\Lambda^3 \rho}\right) - \ln\left(\int ds^N \exp[-\beta U(s^N; L)]\right)$$

$$\beta F = \beta F^{IG} + \beta F^{ex} \quad \beta \mu = \beta \mu^{IG} + \beta \mu^{ex}$$

$$\mu \equiv \left(\frac{\partial F}{\partial N}\right)_{V,T}$$

$$\beta \mu^{IG} \equiv \left(\frac{\partial \beta F^{IG}}{\partial N}\right)_{V,T} \quad \beta \mu^{ex} \equiv \left(\frac{\partial \beta F^{ex}}{\partial N}\right)_{V,T}$$

$$\begin{aligned}
\beta\mu &\equiv \left(\frac{\partial \beta F}{\partial N} \right)_{V,T} & \beta\mu &= \frac{\beta F(N+1) - \beta F(N)}{N+1 - N} \\
& & &= -\ln \frac{Q(N+1)}{Q(N)} \\
& & &= -\ln \left(\frac{\frac{V^{N+1}}{\Lambda^{3N+3}(N+1)!}}{\frac{V^N}{\Lambda^{3N}N!}} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right) \\
& & &= -\ln \left(\frac{V}{\Lambda^3(N+1)} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right) \\
\beta\mu &= \beta\mu^{IG} + \beta\mu^{ex} \\
\beta\mu^{ex} &= -\ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)
\end{aligned}$$

$$\beta\mu^{ex} = -\ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$U(s^{N+1}; L) = \Delta U^+ + U(s^N; L)$$

$$\beta\mu^{ex} = -\ln \left(\frac{\int ds^N \int ds_{N+1} \exp[-\beta(\Delta U^+ + U(s^N; L))]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$= -\ln \left(\frac{\int ds_{N+1} \int ds^N \{ \exp[-\beta \Delta U^+] \} \exp[-\beta U(s^N; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$= -\ln \left(\int ds_{N+1} \langle \exp[-\beta \Delta U^+] \rangle_{NVT} \right)$$

Ghost particle!

Algorithm 16 (Widom Test Particle Insertion)

```
subroutine Widom  
  
xtest=box*ranf()  
call ener(xtest,entest)  
wtest=wtest  
+      +exp(-beta*entest)  
return  
end
```

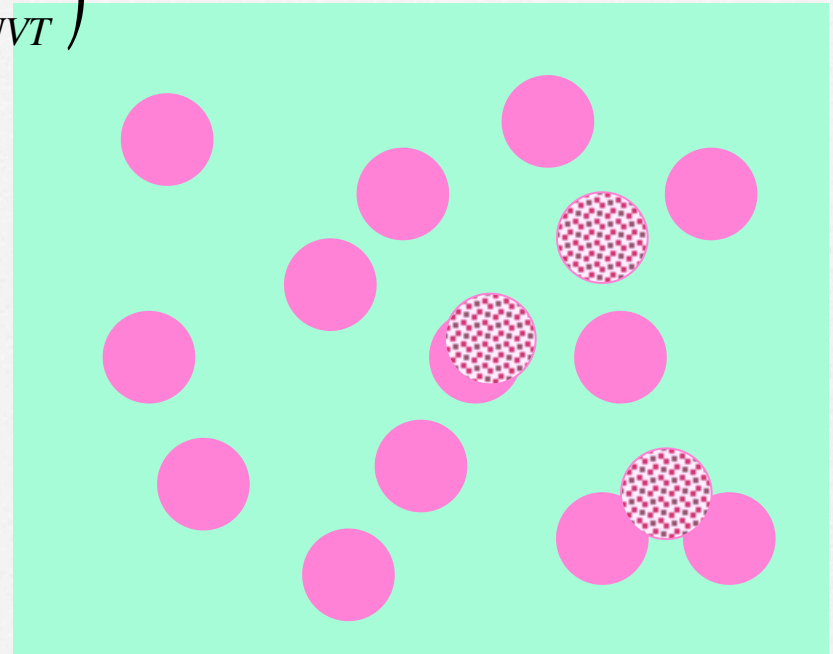
excess chemical potential
via the addition of test particles
generate a random position
determine energy
update Boltzmann factor in (7.2.5)

Hard spheres

$$\beta\mu^{ex} = -\ln\left(\int ds_{N+1} \left\langle \exp[-\beta\Delta U^+] \right\rangle_{NVT}\right)$$

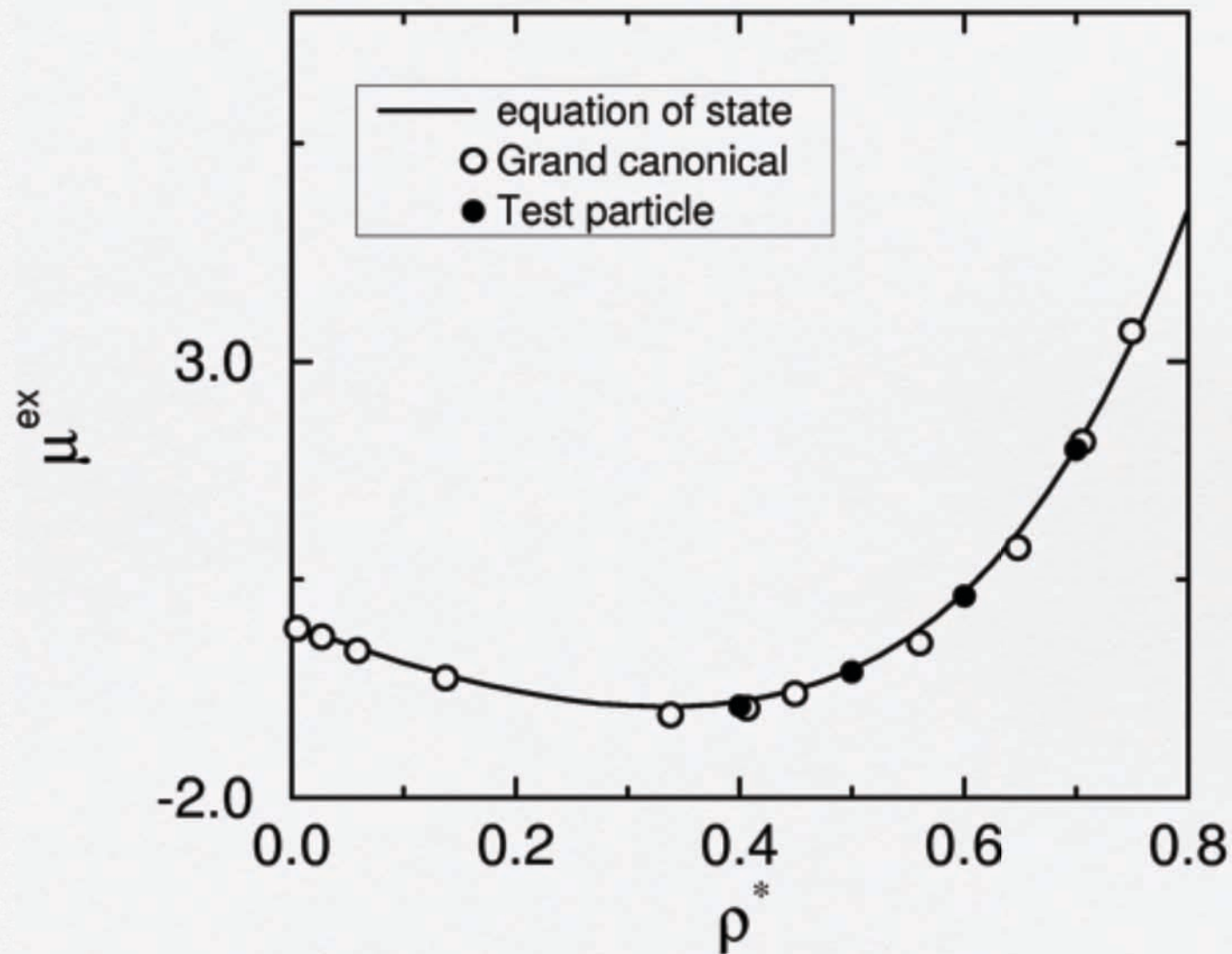
$$U(r) = \begin{cases} \infty & r \leq \sigma \\ 0 & r > \sigma \end{cases}$$

$$\left\langle \exp[-\beta\Delta U^+] \right\rangle = \begin{cases} 0 & \text{if overlap} \\ 1 & \text{no overlap} \end{cases}$$



Probability to insert a test particle!

Lennard-Jones fluid



Real-particle method

$$\beta\mu \equiv \left(\frac{\partial \beta F}{\partial N} \right)_{V,T}$$

$$\begin{aligned} \beta\mu &= \frac{\beta F(N+1) - \beta F(N)}{N+1 - N} \\ &= -\ln \frac{Q(N+1)}{Q(N)} \end{aligned}$$

$$= -\ln \left(\frac{\frac{V^{N+1}}{\Lambda^{3N+3} (N+1)!}}{\frac{V^N}{\Lambda^{3N} N!}} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$= -\ln \left(\frac{V}{\Lambda^3 (N+1)} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = \beta\mu^{IG} + \beta\mu^{ex}$$

$$\beta\mu^{ex} = -\ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu^{ex} = -\ln\left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]}\right)$$

$$U(s^{N+1}; L) = \Delta U^+ + U(s^N; L)$$

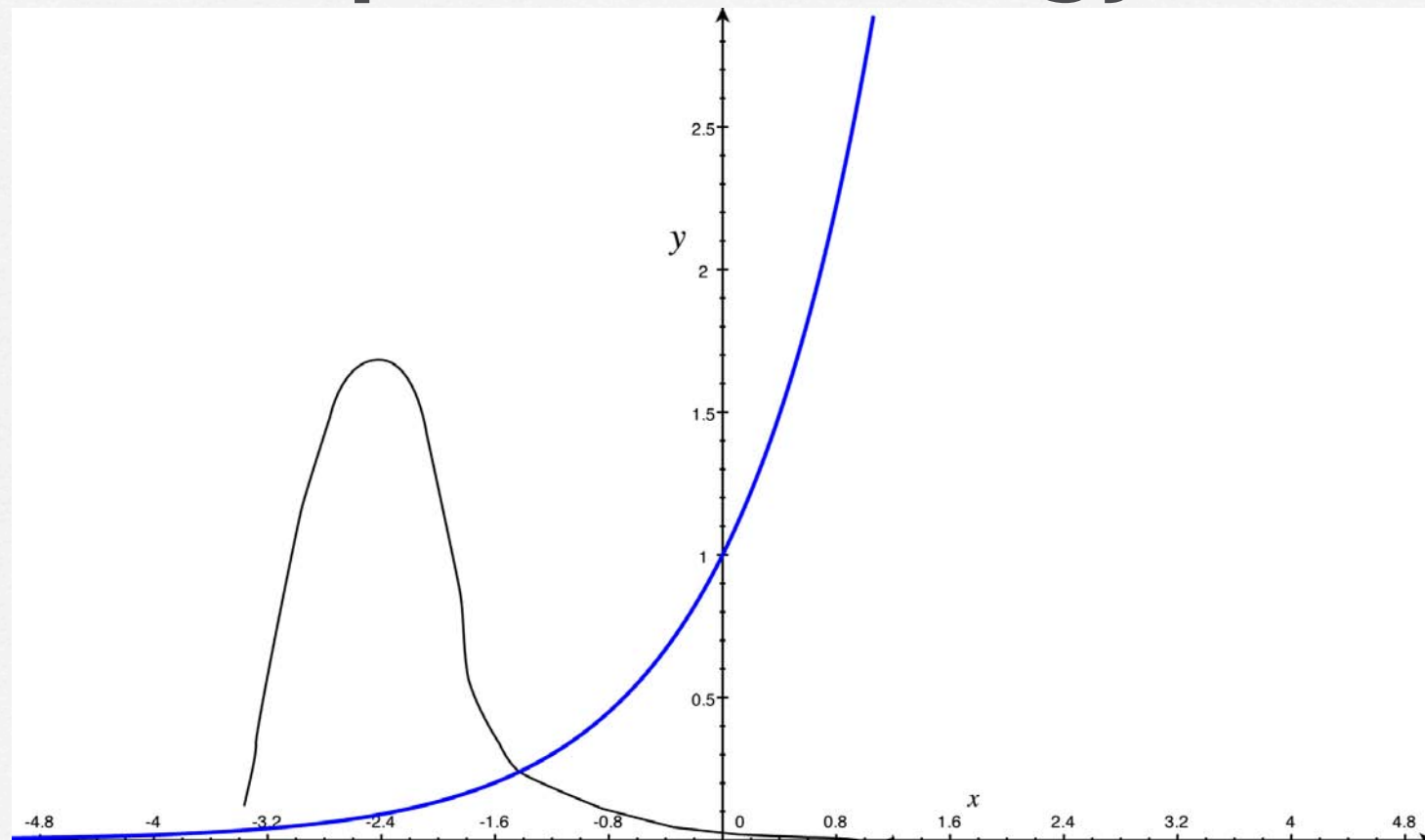
$$U(s^N; L) = U(s^{N+1}; L) - \Delta U^+$$

$$\begin{aligned}\beta\mu^{ex} &= -\ln\left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^{N+1} \exp[+\beta\Delta U^+ - \beta U(s^{N+1}; L)]}\right) \\ &= -\ln\left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^{N+1} \exp[+\beta\Delta U^+] \exp[-\beta U(s^{N+1}; L)]}\right)\end{aligned}$$

$$= +\ln\left(\langle \exp[+\beta\Delta U^+] \rangle_{N+1VT}\right)$$

real particle!

Real particle energy



Other ensembles: NPT

NVT: Helmholtz free energy

NPT: Gibbs free energy

$$\mu \equiv \left(\frac{\partial F}{\partial N} \right)_{V,T}$$

$$\mu \equiv \left(\frac{\partial G}{\partial N} \right)_{P,T}$$

$$\beta G = -\ln(Q_{NPT})$$

$$Q_{NPT} = \frac{1}{\Lambda^{3N} N!} \int dV V^N \exp(-\beta V P) \int ds^N \exp[-\beta U(s^N; L)]$$

$$\beta \mu = \frac{\beta G(N+1) - \beta G(N)}{N+1 - N}$$

$$\beta \mu = -\ln \frac{Q(N+1)}{Q(N)}$$

$$\beta\mu = -\ln \frac{Q(N+1)}{Q(N)}$$

$$\beta\mu = -\ln \frac{\frac{1}{\Lambda^{3N+3} (N+1)!}}{\frac{1}{\Lambda^{3N} N!}} \left(\frac{\int dV V^{N+1} \exp(-\beta VP) \int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = -\ln \frac{1}{\Lambda^3 (N+1)} \left(\frac{\int dV V^N \exp(-\beta VP) V \int ds^N \exp[-\beta U(s^N; L)] \int ds_{N+1} \exp(-\beta \Delta U^+)}{\int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = -\ln \frac{1}{\Lambda^3 (N+1)} \left(\frac{\int dV V^N \exp(-\beta V P) V \int ds^N \exp[-\beta U(s^N; L)] \int ds_{N+1} \exp(-\beta \Delta U^+)}{\int dV V^N \exp(-\beta V P) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = \ln(\Lambda^3 \beta P) - \ln \left(\frac{\beta P}{N+1} \frac{\int dV V^N \exp(-\beta V P) \int ds^N \exp[-\beta U(s^N; L)] \int ds_{N+1} V \exp(-\beta \Delta U^+)}{\int dV V^N \exp(-\beta V P) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

The volume fluctuates!

$$\left\langle \frac{\beta P V}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle \neq \left\langle \frac{\beta P V}{N+1} \right\rangle \left\langle \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle$$

$$\beta\mu = \ln(\Lambda^3 \beta P) - \ln \left\langle \frac{\beta P V}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle$$

NVT versus NPT

NVT:

$$\beta\mu = \beta \ln(\rho) - \ln \left\langle \int ds_{N+1} \exp[-\beta \Delta U^+] \right\rangle_{NVT}$$

NPT:

$$\beta\mu = \ln(\Lambda^3 \beta P) - \ln \left\langle \frac{\beta P V}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle$$

Overlapping Distribution Method

Two systems:

System 0: N, V, T, U_0

System 1: N, V, T, U_1

$$Q_0 = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp(-\beta U_0)$$

$$Q_1 = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp(-\beta U_1)$$

$$\Delta\beta F = \beta F_1 - \beta F_0 = -\ln(Q_1/Q_0)$$

$$= -\ln \left(\frac{\int ds^N \exp(-\beta U_1)}{\int ds^N \exp(-\beta U_0)} \right) = -\ln \left(\frac{q_1}{q_0} \right)$$

$$\int ds^N \exp(-\beta U_1) = \frac{q_1}{q_0} \int ds^N \exp(-\beta U_0)$$

System 0: N, V, T, U_0

System 1: N, V, T, U_1

Let us define

$$p_1(\Delta U) \equiv \frac{\int d\mathbf{s}^N \exp(-\beta U_1) \delta(U_1 - U_0 - \Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_1)}$$

and

$$p_0(\Delta U) \equiv \frac{\int d\mathbf{s}^N \exp(-\beta U_0) \delta(U_1 - U_0 - \Delta U)}{\int d\mathbf{s}^N \exp(-\beta U_0)}$$

$$\Delta\beta F = -\ln\left(\frac{\int ds^N \exp(-\beta U_1)}{\int ds^N \exp(-\beta U_0)}\right) = -\ln\left(\frac{q_1}{q_0}\right)$$

$$p_1(\Delta U) = \frac{\int ds^N \exp(-\beta U_1) \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_1)} \quad p_0(\Delta U) = \frac{\int ds^N \exp(-\beta U_0) \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_0)}$$

= ΔU : Because of the δ function it can be taken outside the integration

$$\begin{aligned} p_1(\Delta U) &= \frac{\int ds^N \exp[-\beta(U_1 - U_0)] \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_1)} \\ &= \frac{\exp[-\beta \Delta U] \int ds^N \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_1)} \end{aligned}$$

$$\int ds^N \exp(-\beta U_1) = \frac{q_1}{q_0} \int ds^N \exp(-\beta U_0)$$

Overlapping Distribution Method

$$p_1(\Delta U) = \frac{\exp[-\beta\Delta U] \int ds^N \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_1)}$$

$$= \frac{q_0}{q_1} \exp(-\beta\Delta U) \frac{\int ds^N \exp[-\beta U_0] \delta(U_1 - U_0 - \Delta U)}{\int ds^N \exp(-\beta U_0)}$$

$$\int ds^N \exp(-\beta U_1) = \frac{q_1}{q_0} \int ds^N \exp(-\beta U_0)$$

$$\frac{q_0}{q_1} = \exp(\beta\Delta F)$$

$$p_1(\Delta U) = \frac{q_0}{q_1} \exp(-\beta\Delta U) p_0(\Delta U)$$

$$\ln p_1(\Delta U) = \beta(\Delta F - \Delta U) + \ln p_0(\Delta U)$$

$$\ln p_1(\Delta U) = \beta(\Delta F - \Delta U) + \ln p_0(\Delta U)$$

Let us define two new functions:

$$f_0(\Delta U) \equiv \ln p_0(\Delta U) - 0.5\beta\Delta U$$

$$\beta\Delta F = f_1(\Delta U) - f_0(\Delta U)$$

$$f_1(\Delta U) \equiv \ln p_1(\Delta U) + 0.5\beta\Delta U$$

Fit f_0 and f_1 to two polynomials that only differ by the offset.

$$f_1(\Delta U) \equiv C_1 + a\Delta U + b\Delta U^2 + c\Delta U^3$$

$$\beta\Delta F = C_1 - C_0$$

$$f_0(\Delta U) \equiv C_0 + a\Delta U + b\Delta U^2 + c\Delta U^3$$

Simulate system 0: compute f_0

Simulate system 1: compute f_1

System 0: $N-1, V, T, U$

+ 1 ideal gas

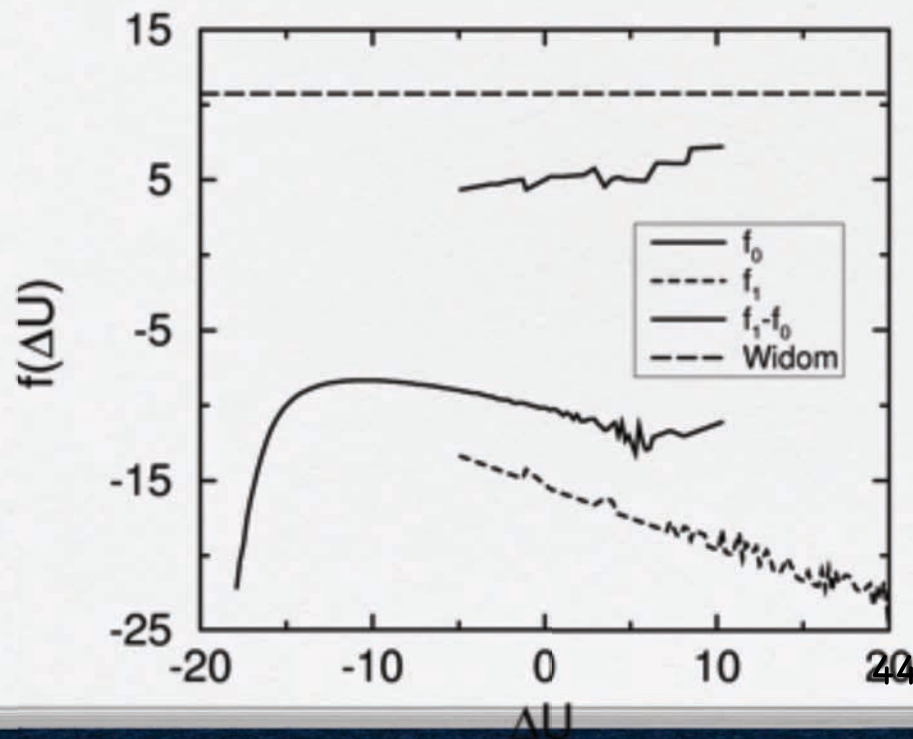
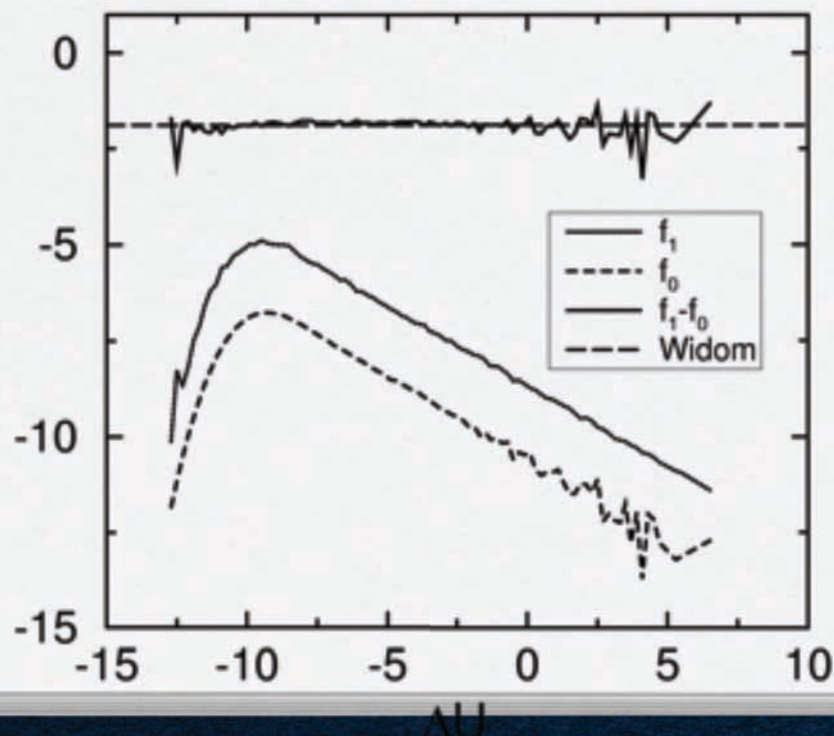
$$\Delta\beta F = \beta F_1 - \beta F_0 \equiv \beta\mu^{ex}$$

System 1: N, V, T, U

$$\Delta U = U_1 - U_0$$

System 0: test particle energy System 1: real particle energy

$$\beta\mu^{ex} = f_1(\Delta U) - f_0(\Delta U)$$

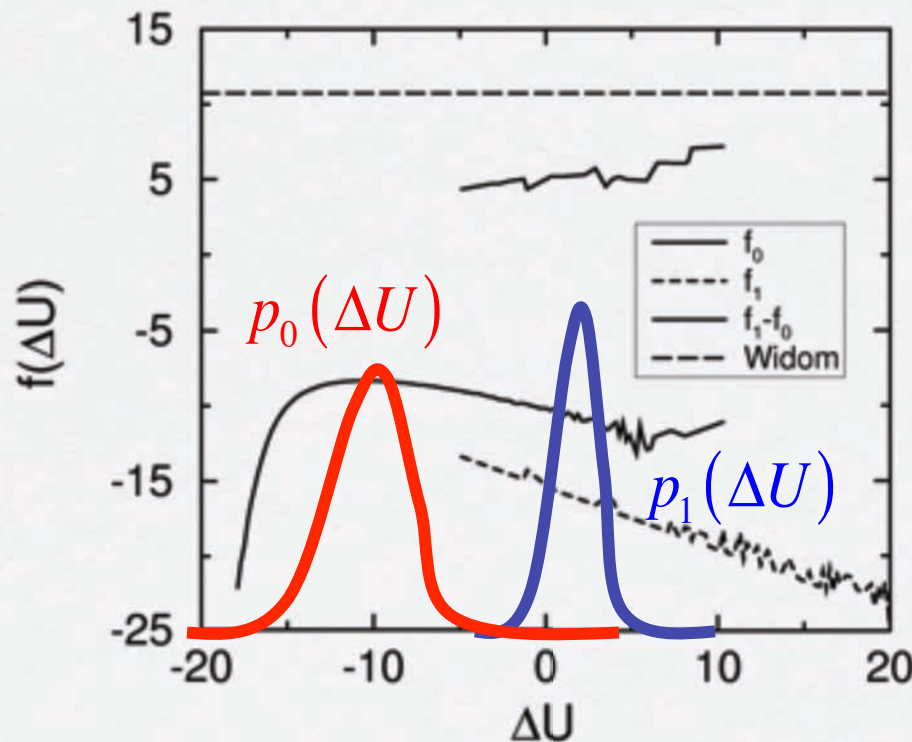


System 0: $N-1, V, T, U$
+ 1 ideal gas

$$f_0(\Delta U) \equiv \ln p_0(\Delta U) - 0.5\beta\Delta U$$

System 1: N, V, T, U

$$f_1(\Delta U) \equiv \ln p_1(\Delta U) + 0.5\beta\Delta U$$



- Accurate sampling only if there is overlap between the two distributions
- How to create this overlap?

Intermezzo: order parameters and Landau Free energies (1)

Landau free energy density:

$$\beta f(q) = -\ln \frac{1}{\Lambda^{3N} N!} \int \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] d\mathbf{r}^N$$

Probability to find the system
with order parameter q

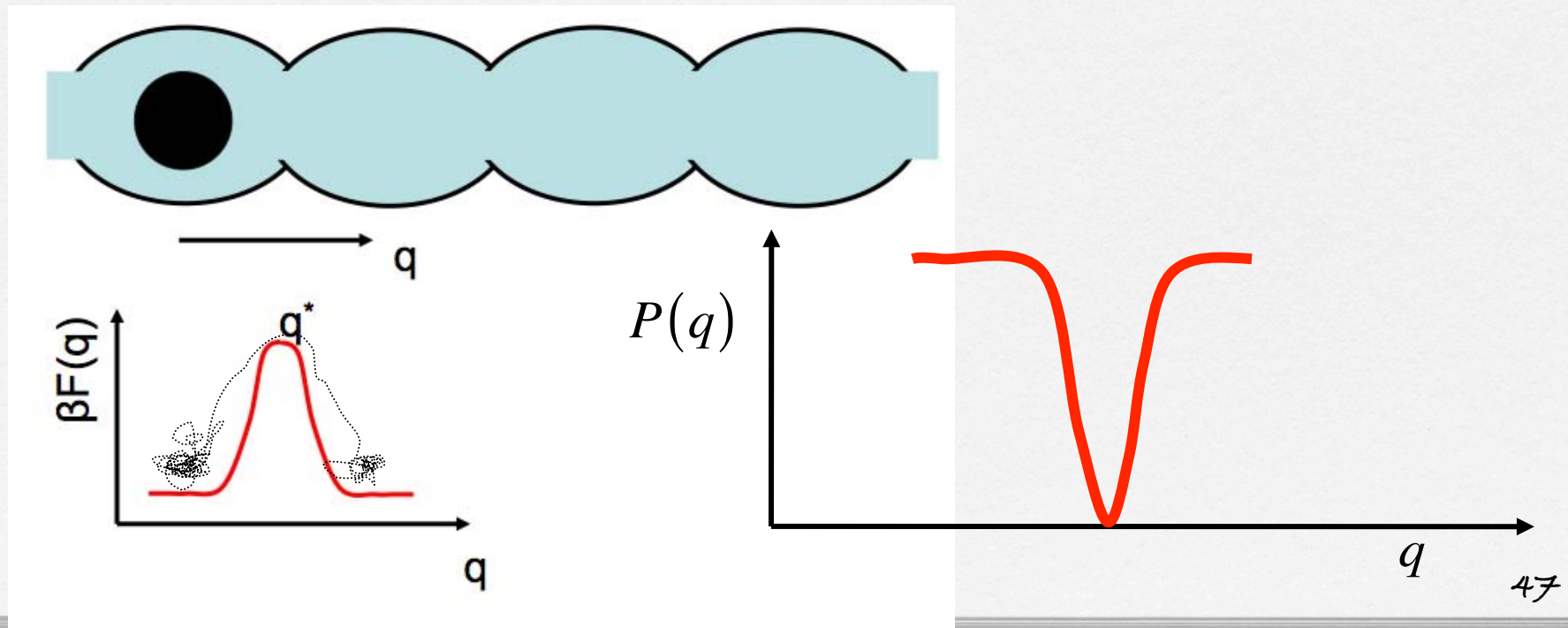
$$P(q) dq \propto \exp[-\beta f(q)] dq$$

Free energy:

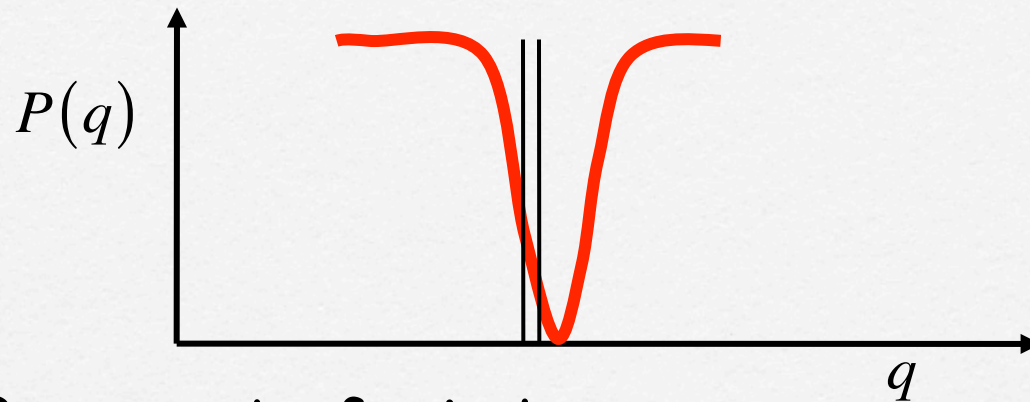
$$F = \int f(q) dq$$

Probability to find the system with order parameter q

$$P(q) dq \propto \exp[-\beta f(q)] dq$$



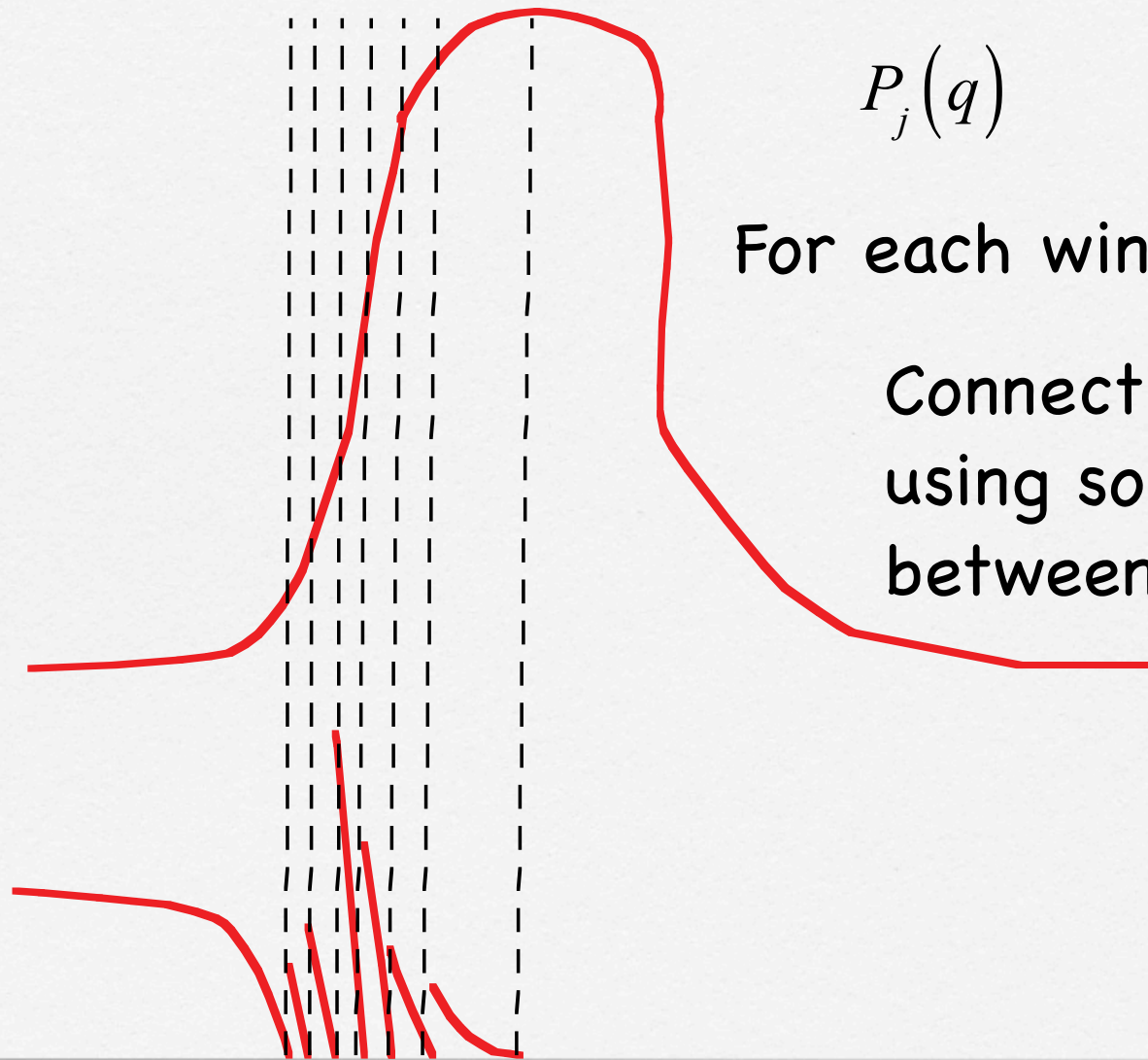
Histogram method (1)



Define a set of windows:

$$U_j(q) = \begin{cases} \infty & \text{if } q < q_j^{\min} \\ 0 & \text{if } q_j^{\min} < q < q_j^{\max} \\ \infty & \text{if } q > q_j^{\max} \end{cases}$$

For each window determine: $P_j(q)$



$$P_j(q)$$

For each window separately

Connect the windows
using some overlap
between Windows

Umbrella sampling (1)

$$\exp[-\beta f(q)] = \frac{1}{\Lambda^{3N} N!} \int \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] d\mathbf{r}^N$$

$$\Delta f(q) = f(q) - f^{IG}(q)$$

$$\exp[-\beta \Delta f(q)] = \frac{\int \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] d\mathbf{r}^N}{\int \delta(q - q(\mathbf{r}^N)) d\mathbf{r}^N}$$

$$\exp[-\beta \Delta f(q)] = \frac{\int \pi(\mathbf{r}^N) \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] d\mathbf{r}^N}{\int \pi(\mathbf{r}^N) \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) d\mathbf{r}^N}$$

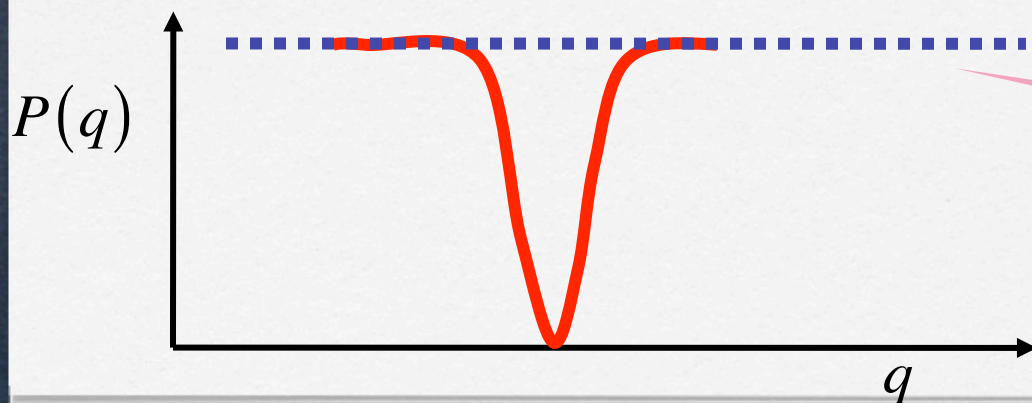
$$\exp[-\beta \Delta f(q)] = \frac{\int \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] \pi(\mathbf{r}^N) d\mathbf{r}^N}{\int \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) \pi(\mathbf{r}^N) d\mathbf{r}^N}$$

$$\exp[-\beta\Delta f(q)] = \frac{\int \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] \pi(\mathbf{r}^N) d\mathbf{r}^N}{\int \pi^{-1}(\mathbf{r}^N) \delta(q - q(\mathbf{r}^N)) \pi(\mathbf{r}^N) d\mathbf{r}^N}$$

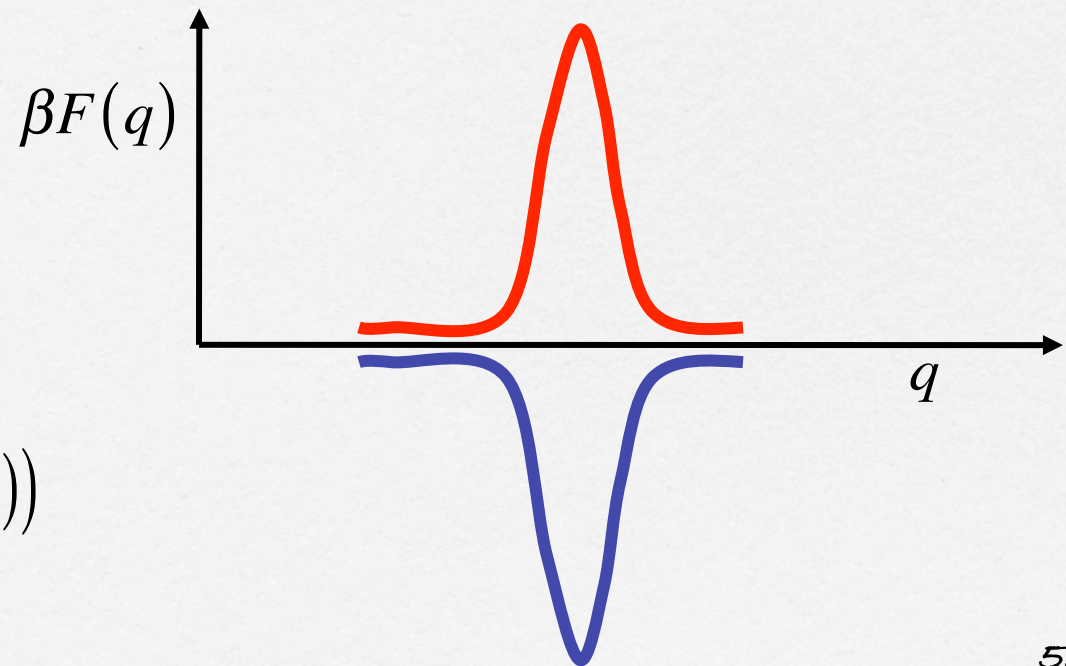
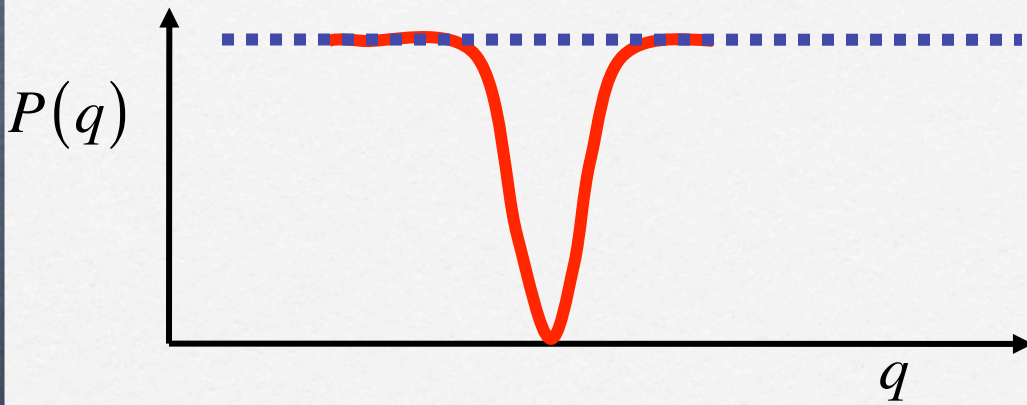
This is NOT a Boltzmann distribution

$$\exp[-\beta\Delta f(q)] = \frac{\langle \delta(q - q(\mathbf{r}^N)) \exp[-\beta U(\mathbf{r}^N)] \pi^{-1}(\mathbf{r}^N) \rangle_{\pi}}{\langle \delta(q - q(\mathbf{r}^N)) \pi^{-1}(\mathbf{r}^N) \rangle_{\pi}}$$

How to choose π ?



Ideal sampling



$$\pi(\mathbf{r}^N) = \exp(-\beta \tilde{F}^{est}(q))$$