

Ensembles II

Free energies and phase equilibria 7

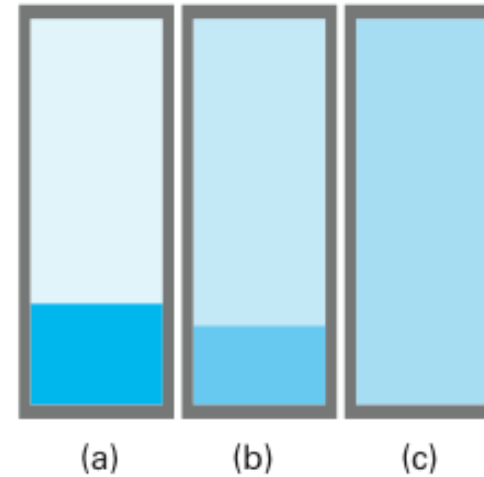
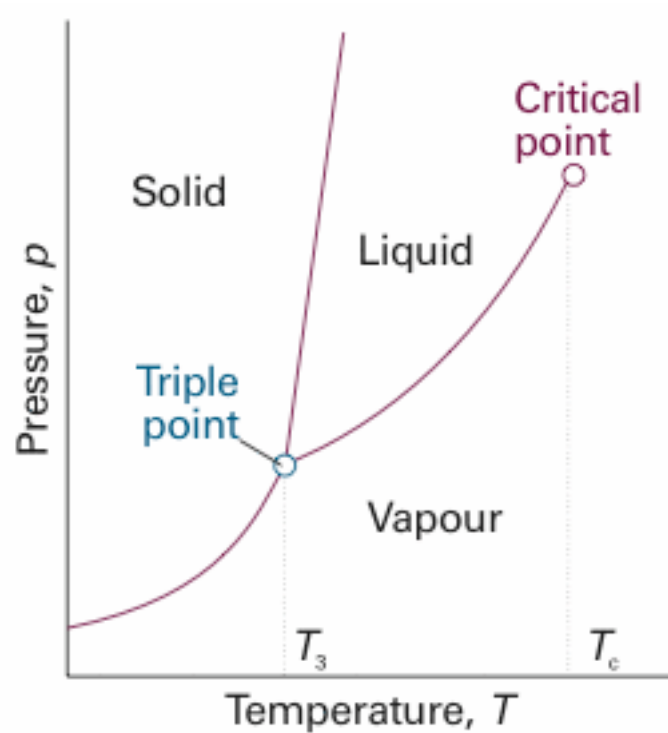
Chemical potentials 7.2

Gibbs ensemble 8

Phase behavior

- Simulations are good for predicting phase behavior
 - Prediction of thermodynamic stability of phases,
 - Coexistence lines
 - Critical points
 - Triple points
 - First order/second order phase transitions

Phase diagrams



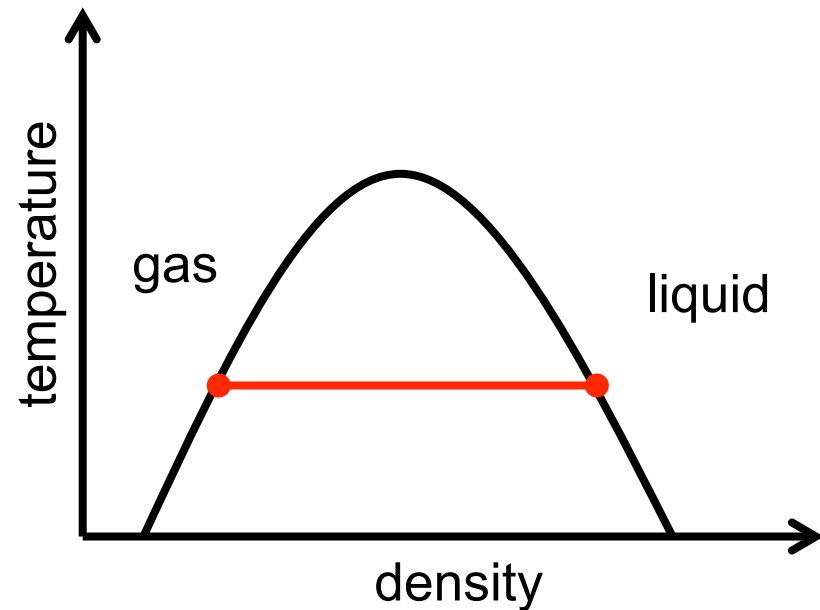
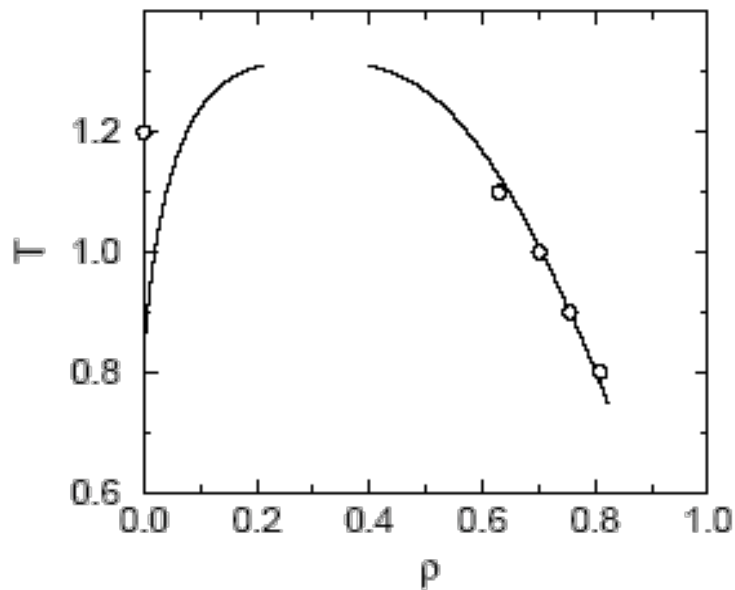
Critical point: no difference between liquid and vapor

Triple point: liquid, vapor and solid in equilibrium.

How do we compute these lines?

Phase diagram

Along the liquid gas coexistence line increasing the pressure and temperature at constant volume the liquid density becomes lower and the vapor density higher.



To determine the phase diagram of a given system we need to know the coexistence densities at given temperature (and pressure).

Easy in experiment, but what about simulations?

Phase equilibrium

Chemical potential

$$\mu = \left(\frac{\partial F}{\partial N} \right)_{V,T} = \left(\frac{\partial G}{\partial N} \right)_{P,T} = G_m$$

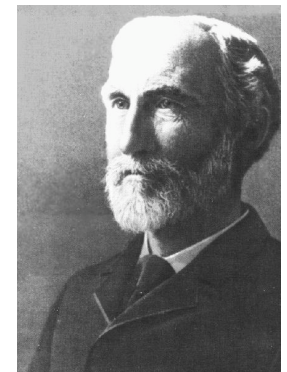
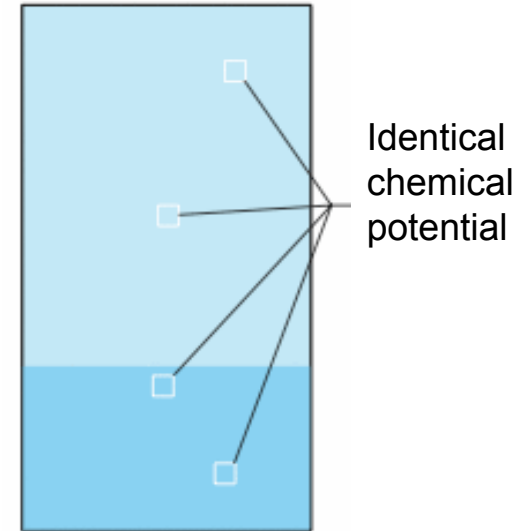
Criterion for equilibrium (for single component)

$$T_\alpha = T_\beta \quad P_\alpha = P_\beta \quad \mu_\alpha = \mu_\beta$$

If $\mu_\alpha > \mu_\beta$ transport of particles from phase α to phase β .

Stable phase: lowest chemical potential

For single component stable phase has lowest Gibbs free energy $G=\mu$



Relation between thermodynamic potentials

$$F = U - TS$$

$$G = F + PV$$

Suppose we have $F(n, V, T)$

Then we can find G from F by realizing

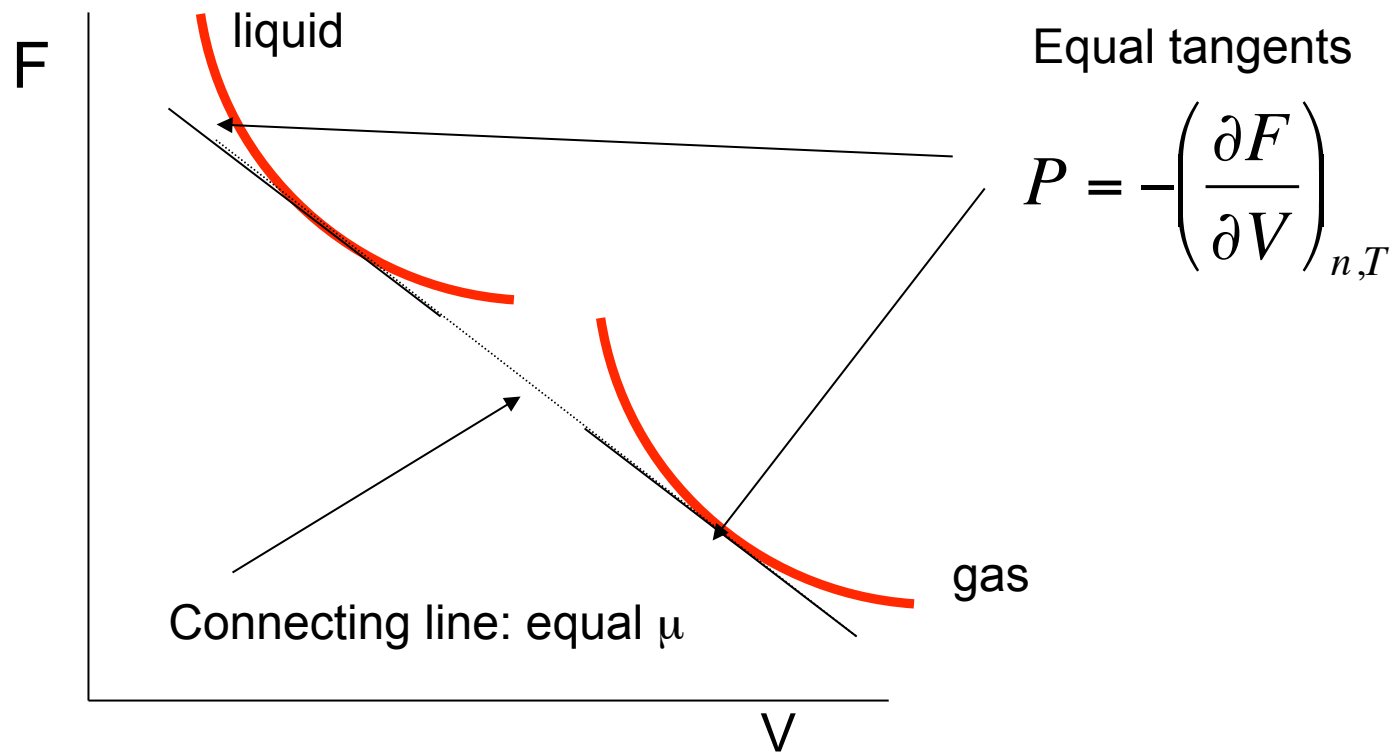
$$P = -\left(\frac{\partial F}{\partial V}\right)_{n,T}$$

$$G = F - V\left(\frac{\partial F}{\partial V}\right)_{n,T}$$

All thermodynamic quantities can be derived from F and its derivatives

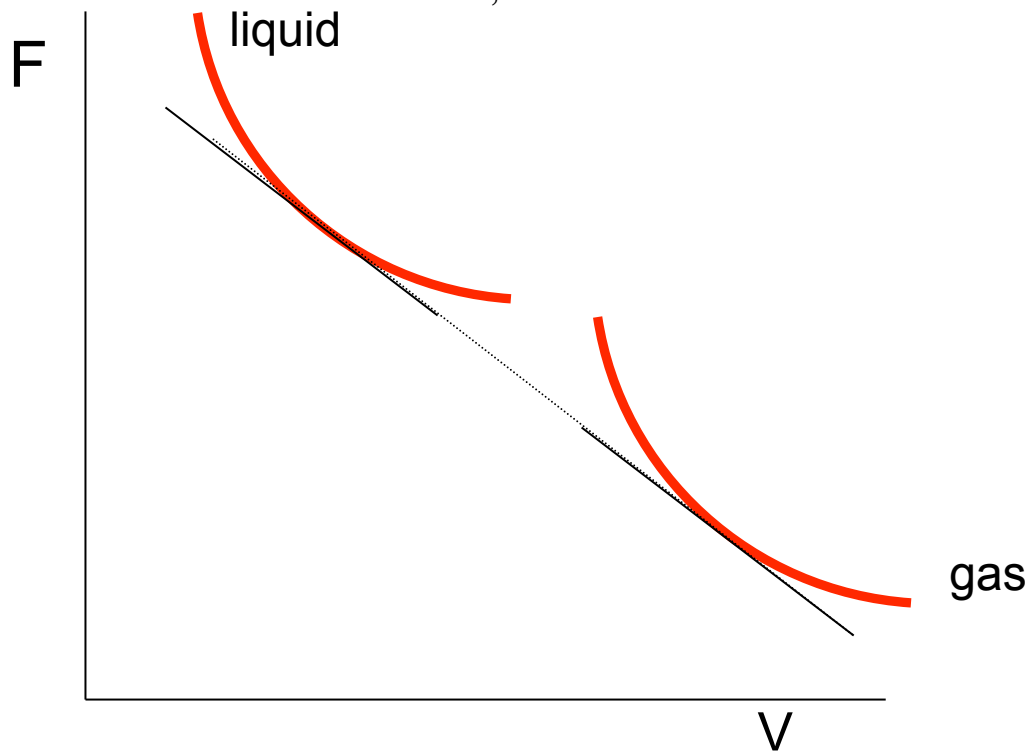
Common tangent construction

- Phase equilibria follows from $F(V,T)$



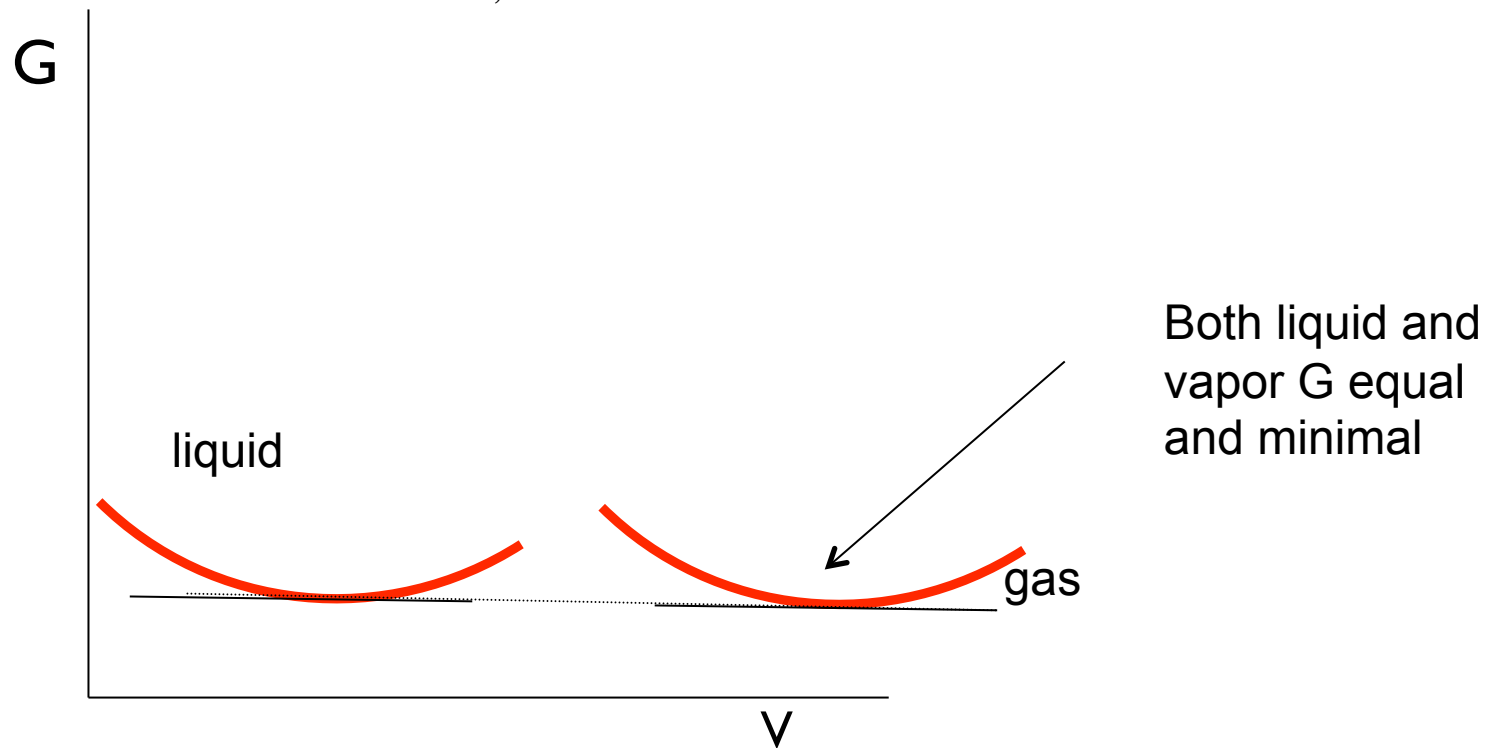
Common tangent construction

$$G = F - V \left(\frac{\partial F}{\partial V} \right)_{n,T}$$



Common tangent construction

$$G = F - V \left(\frac{\partial F}{\partial V} \right)_{n,T}$$



Only equilibrium when P, T is on coexistence line.

We need F or μ

- So equilibrium from $F(V)$ alone or from P and μ

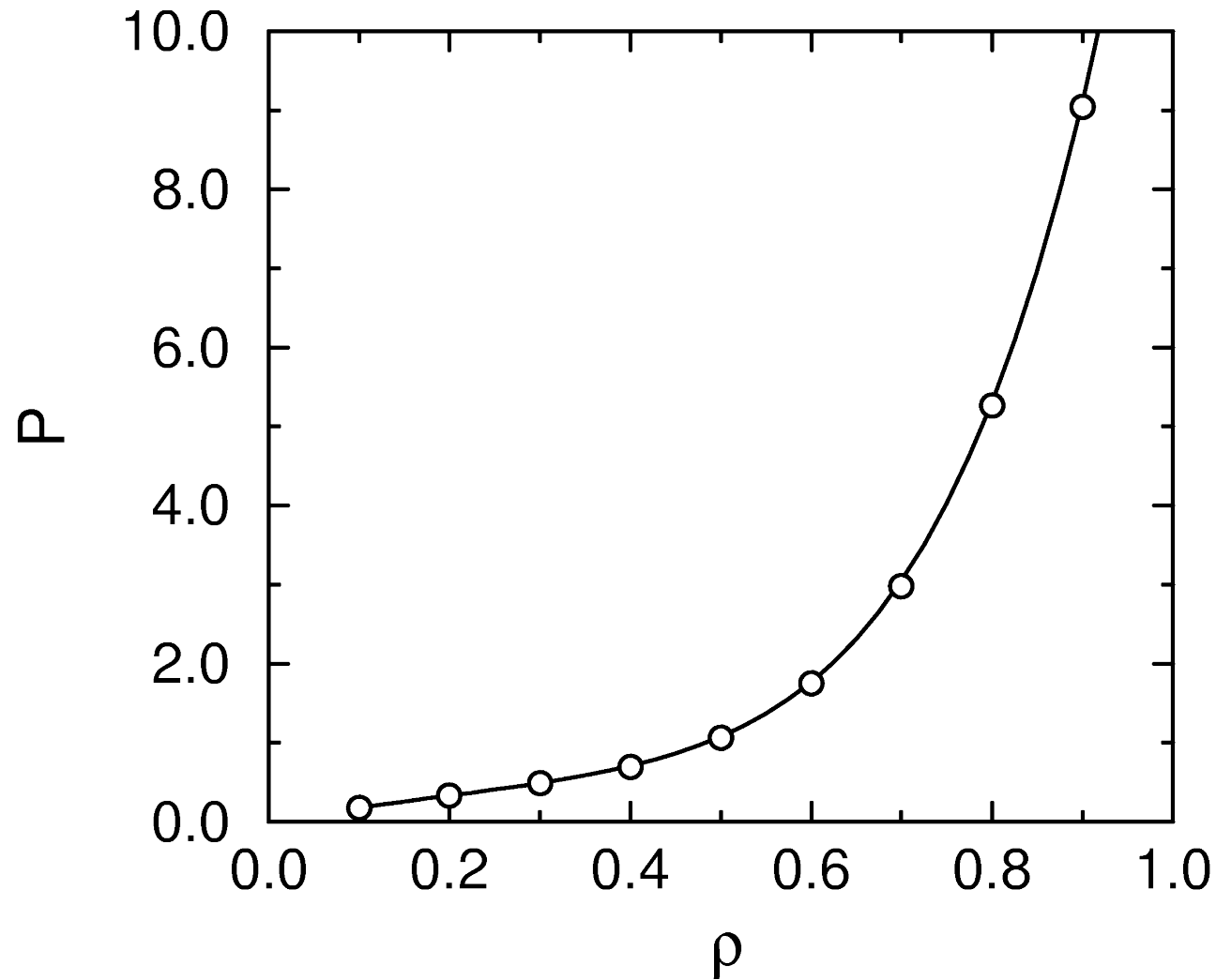
$$F(V) = F(V_0) + \int_{V_0}^V \left(\frac{\partial F}{\partial V} \right)_{N,T} dV = F(V_0) - \int P dV$$

$$F(\rho) = F(\rho_0) + N \int_{\rho_0}^{\rho} \frac{P(\rho')}{\rho'^2} d\rho'$$

- So in fact for only 1 point of the equation of state the F is needed
- For liquid e.o.s even from ideal gas

$$\beta F(\rho)/N = \beta F^{id}(\rho)/N + \int_0^{\rho} \frac{\beta P(\rho') - \rho'}{\rho'^2} d\rho'$$

Equation of state



$$P = P(\rho, T)$$

$$\left(\frac{\partial F}{\partial V} \right)_{N, T} = -P$$

$$F(\rho) = F(\rho_0) + N \int_{\rho_0}^{\rho} \frac{P(\rho')}{\rho'^2} d\rho'$$

Statistical Thermodynamics

Partition function

$$Q_{NVT} = \frac{1}{\Lambda^{3N} N!} \int dr^N \exp[-\beta U(r^N)]$$

Ensemble average

$$\langle A \rangle_{NVT} = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int dr^N A(r^N) \exp[-\beta U(r^N)]$$

Probability to find a particular configuration

$$P(r^N) = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int dr'^N \delta(r'^N - r^N) \exp[-\beta U(r'^N)] \propto \exp[-\beta U(r^N)]$$

Free energy

$$\beta F = -\ln(Q_{NVT})$$

Problem: Q is in general inaccessible

Chemical potential

$$Q_{NVT} = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp[-\beta U(s^N; L)]$$

$$\beta F = -\ln(Q_{NVT})$$

$$= -\ln\left(\frac{V^N}{\Lambda^{3N} N!}\right) - \ln\left(\int ds^N \exp[-\beta U(s^N; L)]\right)$$

$$= -N \ln\left(\frac{1}{\Lambda^3 \rho}\right) + N - \ln\left(\int ds^N \exp[-\beta U(s^N; L)]\right)$$

$$\beta F = \beta F^{IG} + \beta F^{ex}$$

$$\mu \equiv \left(\frac{\partial F}{\partial N}\right)_{V,T}$$

}

$$\beta \mu^{IG} \equiv \left(\frac{\partial \beta F^{IG}}{\partial N}\right)_{V,T}$$

$$\beta \mu^{ex} \equiv \left(\frac{\partial \beta F^{ex}}{\partial N}\right)_{V,T}$$

$$\beta \mu = \beta \mu^{IG} + \beta \mu^{ex}$$

$$\beta F^{IG} = N[\ln(\Lambda^3 \rho) - 1]$$

$$\beta \mu^{IG} \equiv \left(\frac{\partial \beta F^{ig}}{\partial N}\right)_{V,T}$$

$$= \ln(\Lambda^{3N} \rho)$$

$$\beta \mu^{IG} = \beta \mu^0 + \ln(\rho)$$

Widom test particle insertion

$$\beta\mu \equiv \left(\frac{\partial \beta F}{\partial N} \right)_{V,T}$$

$$\begin{aligned} \beta\mu &= \frac{\beta F(N+1) - \beta F(N)}{N+1 - N} \\ &= -\ln \frac{Q(N+1)}{Q(N)} \\ &= -\ln \left(\frac{\frac{V^{N+1}}{\Lambda^{3N+3}(N+1)!}}{\frac{V^N}{\Lambda^{3N}N!}} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right) \\ &= -\ln \left(\frac{V}{\Lambda^3(N+1)} \right) - \ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right) \end{aligned}$$

$$\beta\mu = \beta\mu^{IG} + \beta\mu^{ex}$$

$$\beta\mu^{ex} = -\ln \left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]} \right)$$

Widom test particle insertion

$$\beta\mu^{ex} = -\ln\left(\frac{\int ds^{N+1} \exp[-\beta U(s^{N+1}; L)]}{\int ds^N \exp[-\beta U(s^N; L)]}\right)$$

$$U(s^{N+1}; L) = \Delta U^+ + U(s^N; L)$$

$$\beta\mu^{ex} = -\ln\left(\frac{\int ds^N \int ds_{N+1} \exp[-\beta(\Delta U^+ + U(s^N; L))]}{\int ds^N \exp[-\beta U(s^N; L)]}\right)$$

$$= -\ln\left(\frac{\int ds_{N+1} \int ds^N \left\{ \exp[-\beta \Delta U^+] \right\} \exp[-\beta U(s^N; L)]}{\int ds^N \exp[-\beta U(s^N; L)]}\right)$$

$$= -\ln\left(\int ds_{N+1} \left\langle \exp[-\beta \Delta U^+] \right\rangle_{NVT}\right)$$

Ghost particle!

Algorithm 16 (Widom Test Particle Insertion)

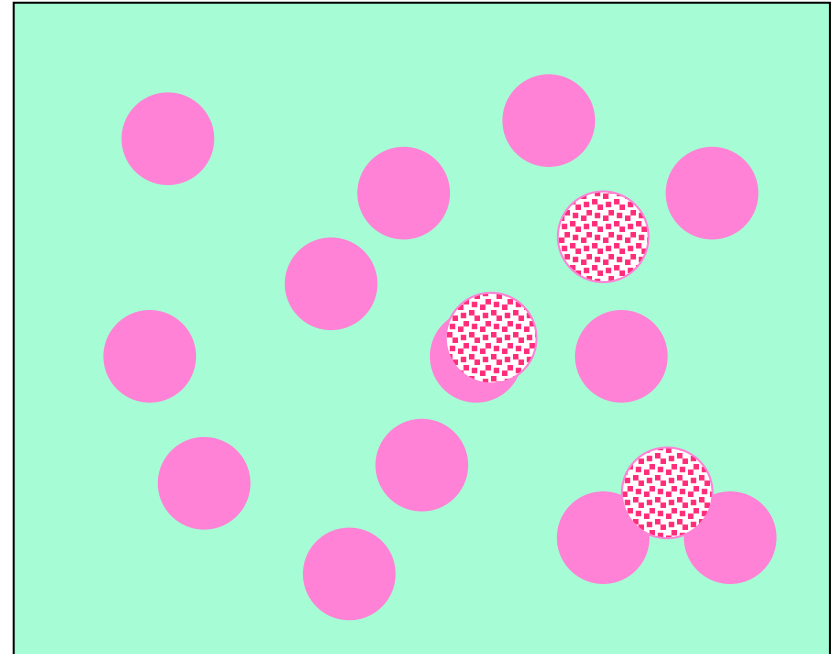
subroutine Widom	excess chemical potential
	via the addition of test particles
xtest=box*ranf()	generate a random position
call ener(xtest,entest)	determine energy
wtest=wtest	update Boltzmann factor in (7.2.5)
+ +exp(-beta*entest)	
return	
end	

Hard spheres

$$\beta\mu^{ex} = -\ln\left(\int ds_{N+1} \left\langle \exp[-\beta\Delta U^+] \right\rangle_{NVT}\right)$$

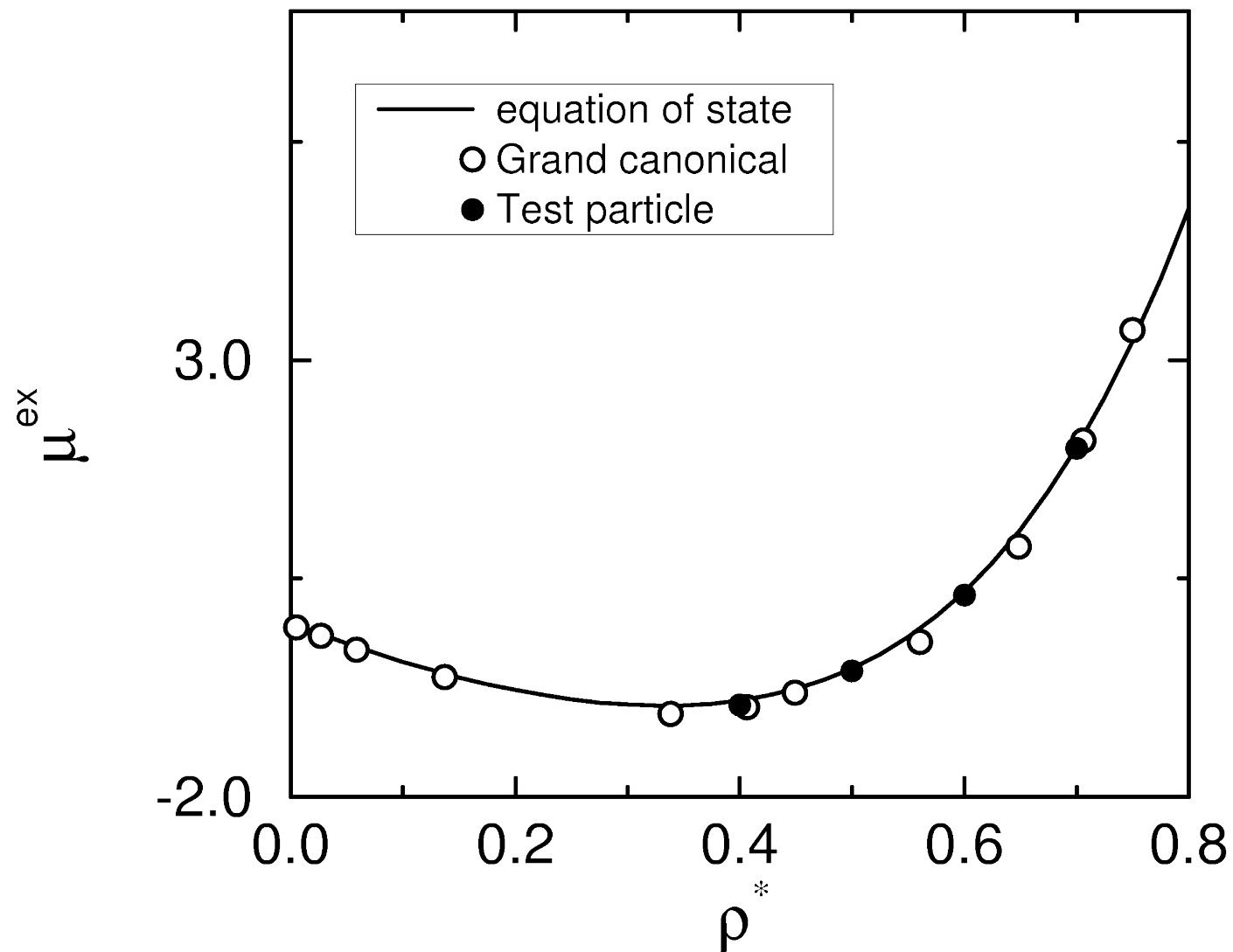
$$U(r) = \begin{cases} \infty & r \leq \sigma \\ 0 & r > \sigma \end{cases}$$

$$\left\langle \exp[-\beta\Delta U^+] \right\rangle = \begin{cases} 0 & \text{if overlap} \\ 1 & \text{no overlap} \end{cases}$$



Probability to insert a test particle!

Lennard-Jones fluid



Other ensembles: *NPT*

NVT: Helmholtz free energy

NPT: Gibbs free energy

$$\mu \equiv \left(\frac{\partial F}{\partial N} \right)_{V,T}$$

$$\mu \equiv \left(\frac{\partial G}{\partial N} \right)_{P,T}$$

$$\beta G = -\ln(Q_{NPT}) \quad Q_{NPT} = \frac{1}{\Lambda^{3N} N!} \int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)]$$

$$\beta\mu = \frac{\beta G(N+1) - \beta G(N)}{N+1 - N}$$

$$\beta\mu = -\ln \frac{Q(N+1)}{Q(N)}$$

$$\beta\mu = -\ln \frac{\frac{1}{\Lambda^{3N+3} (N+1)!} \left(\int dV V^{N+1} \exp(-\beta VP) \int ds^{N+1} \exp[-\beta U(s^{N+1}; L)] \right)}{\frac{1}{\Lambda^{3N} N!} \left(\int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)] \right)}$$

The volume fluctuates!

$$\beta\mu = -\ln \left(\frac{\left\langle \frac{\beta PV}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle}{\int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = \ln \left(\frac{\left\langle \frac{\beta PV}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle}{\int dV V^N \exp(-\beta VP) \int ds^N \exp[-\beta U(s^N; L)]} \right)$$

$$\beta\mu = \ln(\Lambda^3 \beta P) - \ln \left\langle \frac{\beta PV}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle$$

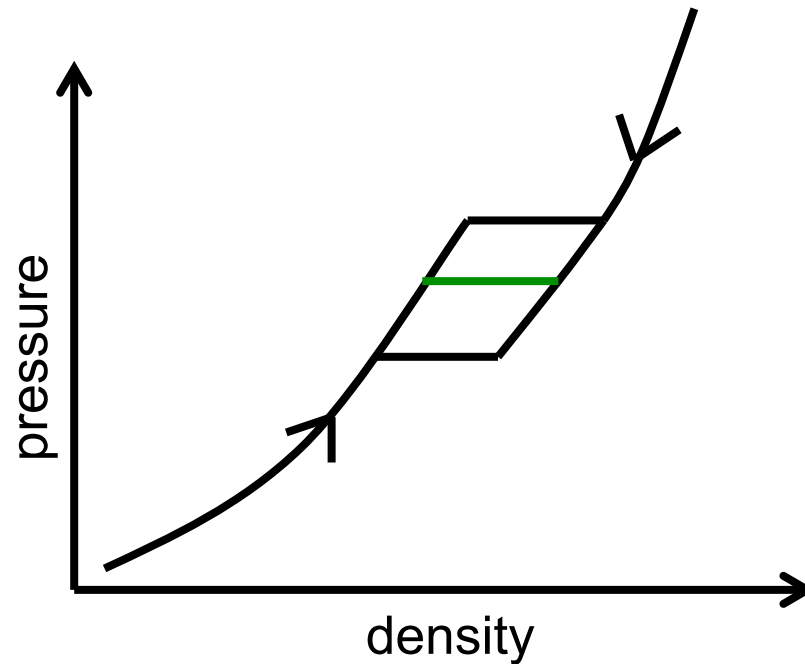
NVT:

$$\beta\mu = \beta \ln(\rho) - \ln \left\langle \int ds_{N+1} \exp[-\beta \Delta U^+] \right\rangle_{NVT}$$

NPT:

$$\beta\mu = \ln(\Lambda^3 \beta P) - \ln \left\langle \frac{\beta P V}{N+1} \int ds_{N+1} \exp(-\beta \Delta U^+) \right\rangle$$

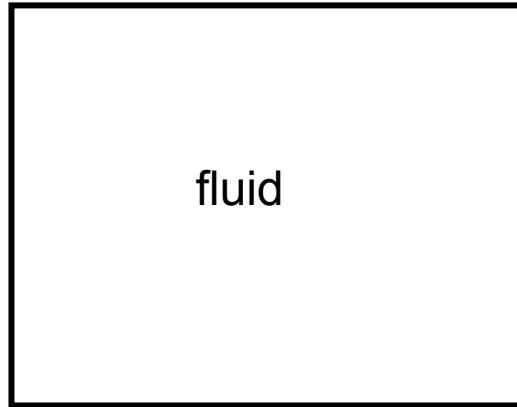
Phase equilibrium and hysteresis



solution : determine equilibrium from chemical potentials

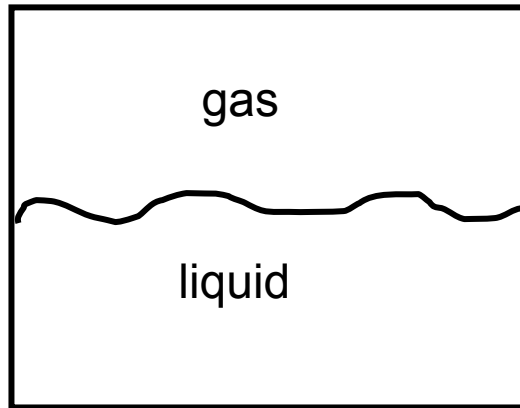
Idea: can't we compute the equilibrium directly ?

NVT Ensemble



Let's lower the temperature...

NVT Ensemble



Problem:

The systems we study are usually small $\Rightarrow$
large fraction of all particles resides in/near interface.

leads to bad estimate of phase coexistence

Possible solution #1

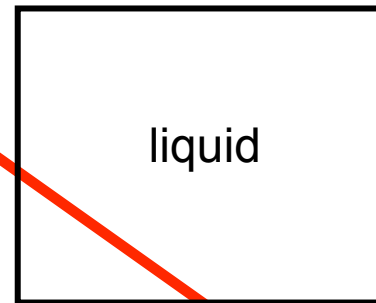
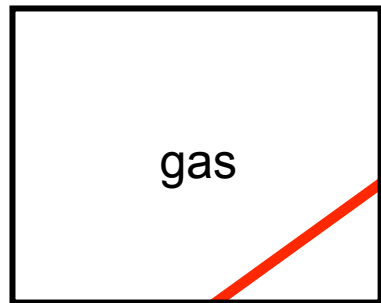
larger systems:

particles	% of part. in interface
1 000	49%
64 000	14%
1 million	6%

⇒ we need huge systems ⇒ computationally expensive

Possible solution #2: “ μPT ”-Ensemble

If we could compute liquid and gas in two separate simulations at constant μ , P and T

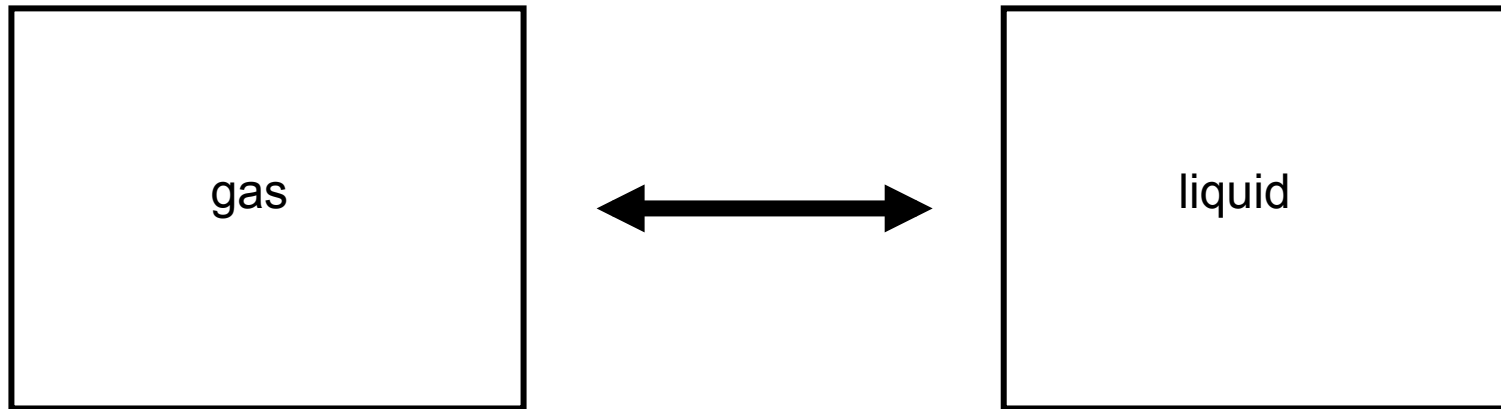


Problem: no such ensemble exists!

- μ , P and T are intensive parameters
- extensive ones unbounded

We have to fix at least one extensive variable (such as N or V)

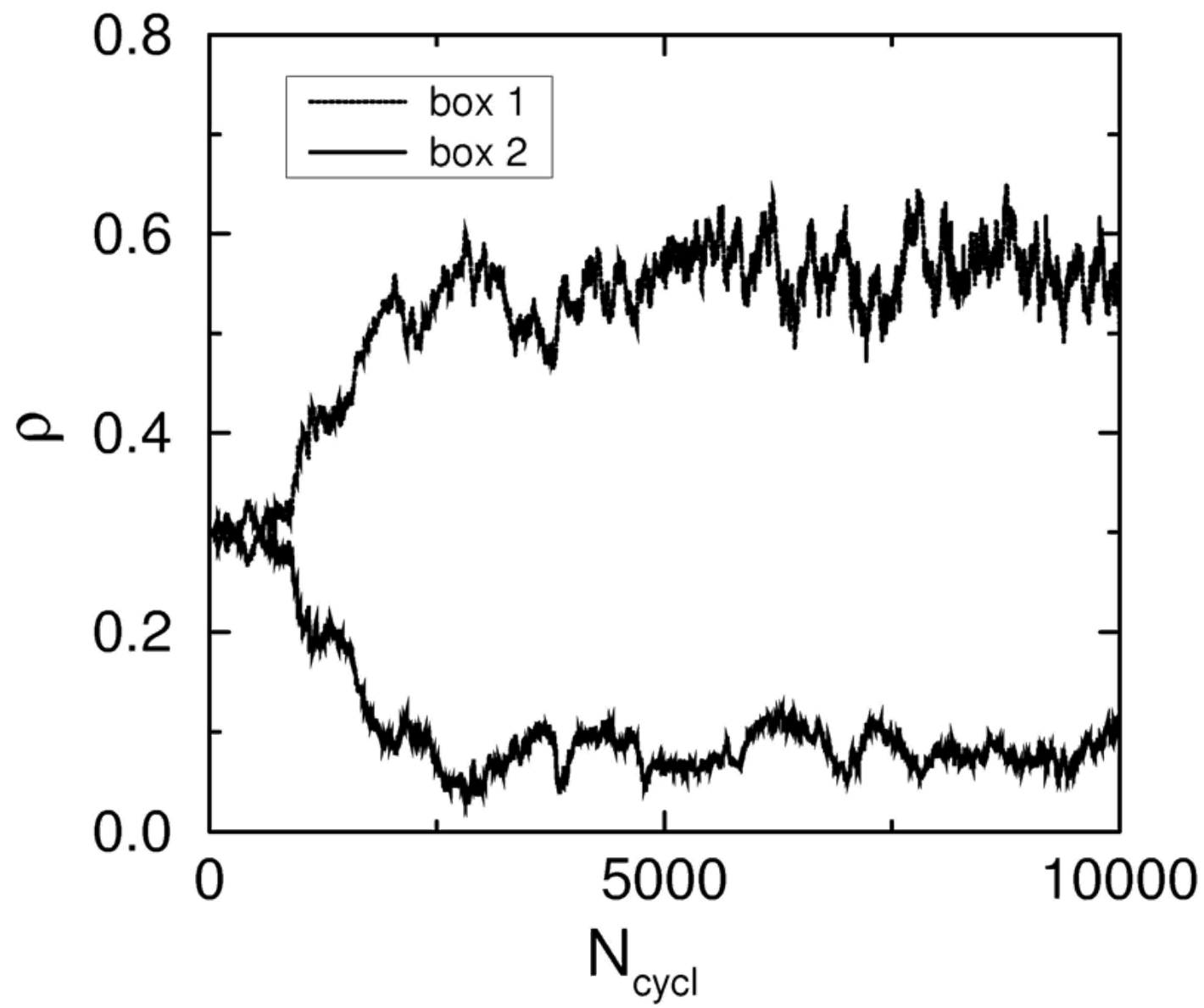
Possible solution #3: The Gibbs ensemble



achieve equilibrium
by coupling them

A. Z. Panagiotopoulos, 1987.



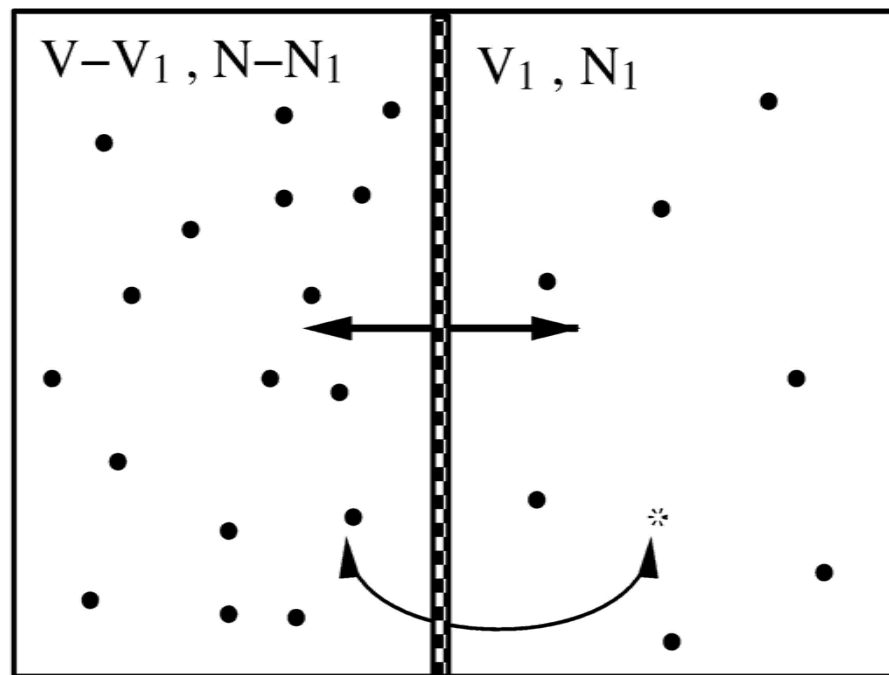


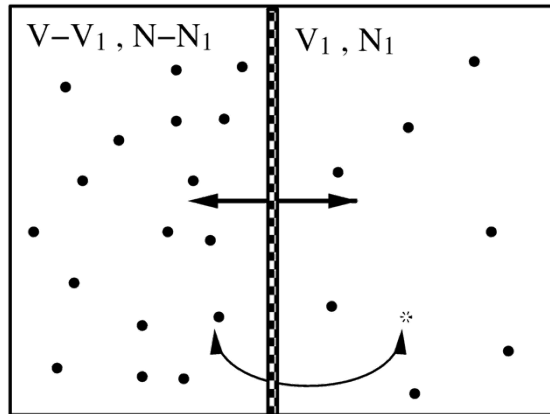
Overall system: NVT ensemble

$$N = N_1 + N_2$$

$$V = V_1 + V_2$$

$$T_1 = T_2$$

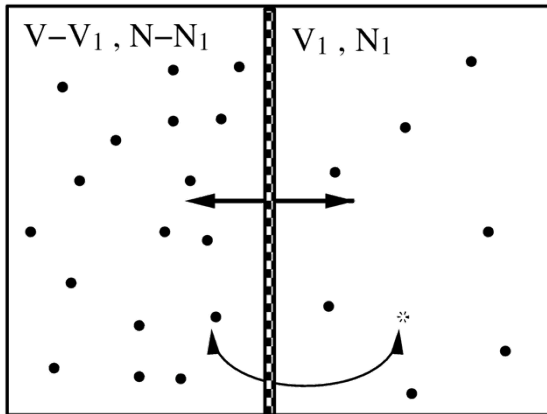




- distribute N_1 particles
- change the volume V_1
- displace the particles

partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1!)} \int_0^V dV_1 V_1^{N_1} (V - V_1)^{N-N_1} \int ds_1^{N_1} \exp[-\beta U(s_1^{N_1})] \int ds_2^{N-N_1} \exp[-\beta U(s_2^{N-N_1})]$$

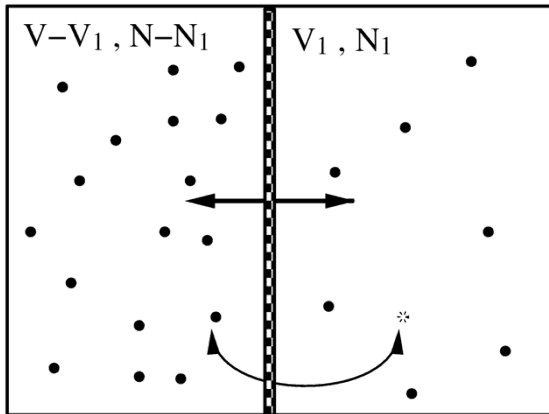


partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1)!} \int_0^V dV_1 V_1^{N_1} (V - V_1)^{N - N_1} \int ds_1^{N_1} \exp[-\beta U(s_1^{N_1})] \int ds_2^{N - N_1} \exp[-\beta U(s_2^{N - N_1})]$$

Distribute N_1 particles over two volumes:

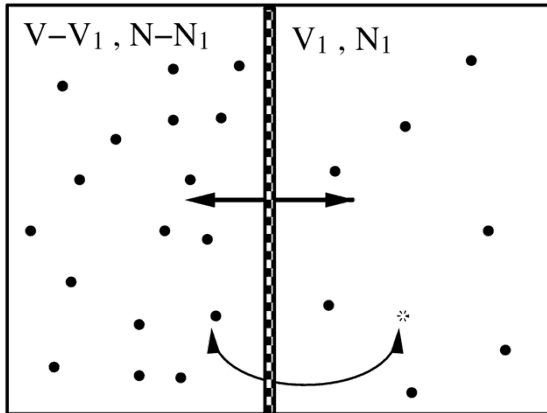
$$\binom{N}{N_1} = \frac{N!}{N_1! (N - N_1)!}$$



partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1!)} \int_0^V dV_1 V_1^{N_1} (V - V_1)^{N - N_1} \int ds_1^{N_1} \exp[-\beta U(s_1^{N_1})] \int ds_2^{N - N_1} \exp[-\beta U(s_2^{N - N_1})]$$

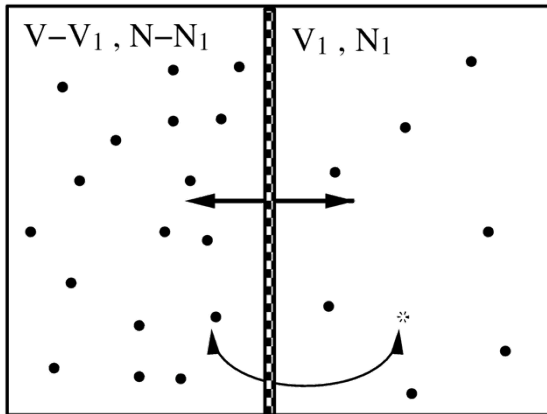
Integrate volume V_1



partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1!)} \int_0^V dV_1 V_1^{N_1} (V - V_1)^{N-N_1} \\ \int ds_1^{N_1} \exp [-\beta U(s_1^{N_1})] \int ds_2^{N-N_1} \exp [-\beta U(s_2^{N-N_1})]$$

Displace the particles in box1 and box2

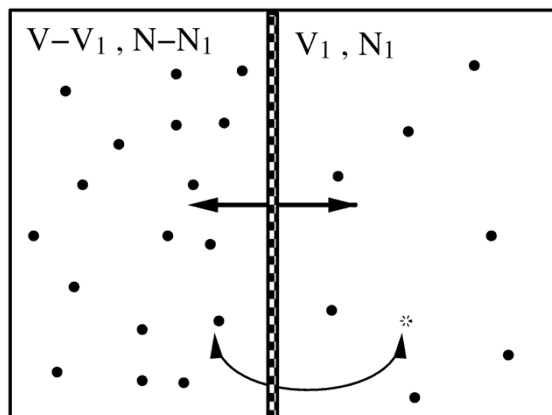


partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1!)} \int_0^V dV_1 \boxed{V_1^{N_1} (V - V_1)^{N - N_1}} \\ \int ds_1^{N_1} \exp[-\beta U(s_1^{N_1})] \int ds_2^{N - N_1} \exp[-\beta U(s_2^{N - N_1})]$$

Displace the particles in box 1 and box2

scaled coordinates in $[0,1)$



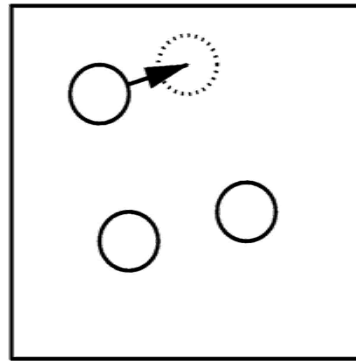
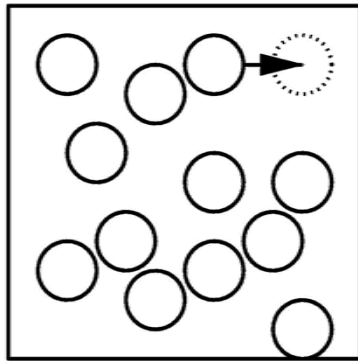
partition function:

$$Q_G(N, V, T) = \sum_{N_1=0}^N \frac{1}{V \Lambda^{3N} N_1! (N - N_1)!} \int_0^V dV_1 V_1^{N_1} (V - V_1)^{N - N_1} \\ \int ds_1^{N_1} \exp[-\beta U(s_1^{N_1})] \int ds_2^{N - N_1} \exp[-\beta U(s_2^{N - N_1})]$$

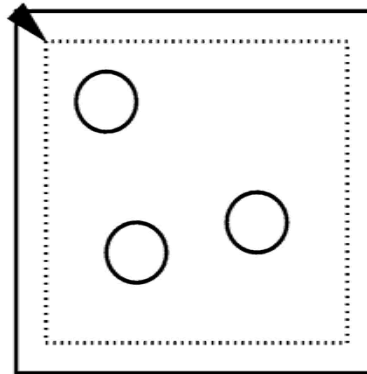
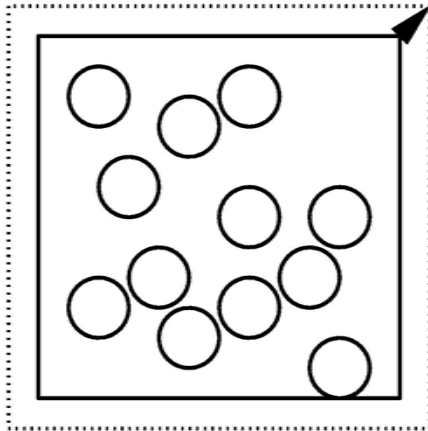
probability distribution:

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N - N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N - N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N - N_1}) \right] \right\}$$

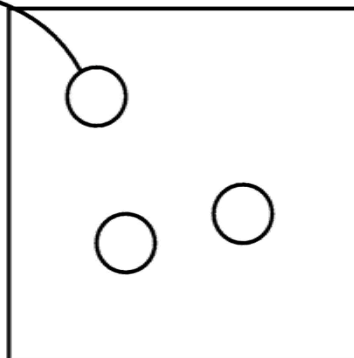
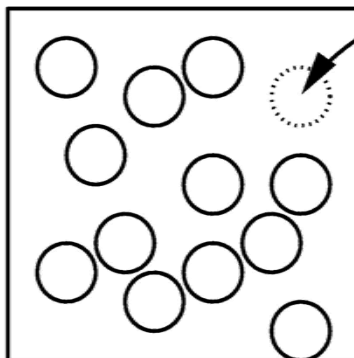
3 different kinds of trial moves:



Particle displacement



Volume change
 $\Rightarrow$ equal P



Particle exchange
 $\Rightarrow$ equal μ

Acceptance rules

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N-N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N-N_1}) \right] \right\}$$

Detailed Balance:

$$K(o \rightarrow n) = K(n \rightarrow o)$$

$$\mathcal{N}(o) \times \alpha(o \rightarrow n) \times \text{acc}(o \rightarrow n) = \mathcal{N}(n) \times \alpha(n \rightarrow o) \times \text{acc}(n \rightarrow o)$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\mathcal{N}(n) \times \alpha(n \rightarrow o)}{\mathcal{N}(o) \times \alpha(o \rightarrow n)}$$

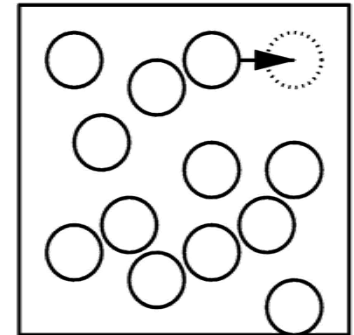
$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\mathcal{N}(n)}{\mathcal{N}(o)}$$

Displacement of a particle in box1

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N-N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N-N_1}) \right] \right\}$$

$$\mathcal{N}(n) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(n) + U(s_2^{N-N_1}) \right] \right\}$$

$$\mathcal{N}(o) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(o) + U(s_2^{N-N_1}) \right] \right\}$$



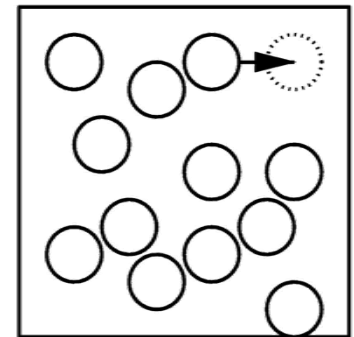
$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\cancel{\frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!}} \exp \left\{ -\beta \left[U(n) + \cancel{U(s_2^{N-N_1})} \right] \right\}}{\cancel{\frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!}} \exp \left\{ -\beta \left[U(o) + \cancel{U(s_2^{N-N_1})} \right] \right\}}$$

Displacement of a particle in box1

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N-N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N-N_1}) \right] \right\}$$

$$\mathcal{N}(n) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(n) + U(s_2^{N-N_1}) \right] \right\}$$

$$\mathcal{N}(o) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(o) + U(s_2^{N-N_1}) \right] \right\}$$

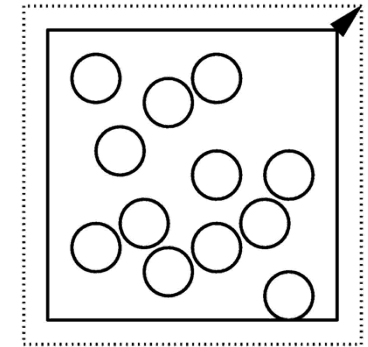


$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\exp \{ -\beta U_1(n) \}}{\exp \{ -\beta U_1(o) \}}$$

Volume change

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N-N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N-N_1}) \right] \right\}$$

$$V_1^n = V_1^o + \Delta V$$



$$\frac{\mathcal{N}(n)}{\mathcal{N}(o)} = \frac{\frac{V_{1,n}^{N_1} (V - V_{1,n})^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_{1,n}^{N_1}) + U(s_{2,n}^{N-N_1}) \right] \right\}}{\frac{V_{1,o}^{N_1} (V - V_{1,o})^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_{1,o}^{N_1}) + U(s_{2,o}^{N-N_1}) \right] \right\}}$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{V_{1,n}^{N_1} (V - V_{1,n})^{N-N_1} \exp \left\{ -\beta U(s_n^N) \right\}}{V_{1,o}^{N_1} (V - V_{1,o})^{N-N_1} \exp \left\{ -\beta U(s_o^N) \right\}}$$

Volume change

More efficient: random walk in $\ln [V_1/(V-V_1)]$

$$Q_G(N, V, T) = \frac{1}{\Lambda^{3N} N!} \sum_{N_1=0}^N \binom{N}{N_1} \\ \times \int_{-\infty}^{\infty} d \ln \left(\frac{V_1}{V - V_1} \right) \frac{V_1(V - V_1)}{V} V_1^{N_1} (V - V_1)^{N-N_1} \\ \times \int ds_1^{N_1} \exp \left[-\beta U(s_1^{N_1}) \right] \int ds_2^{N-N_1} \exp \left[-\beta U(s_2^{N-N_1}) \right]$$

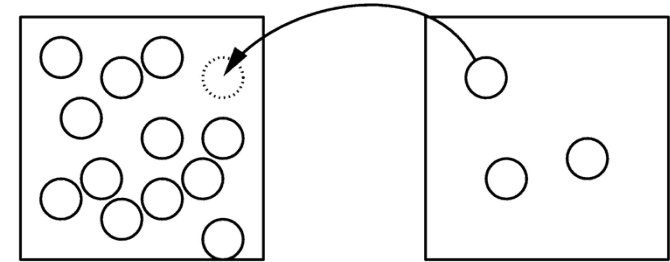
$$\mathcal{N}(n) \propto \frac{V_{1,n}^{N_1+1} (V - V_{1,n})^{N-(N_1+1)}}{V N_1! (N - N_1)!} \exp \left[-\beta U(s_n^N) \right]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{V_{1,n}^{N_1+1} (V - V_{1,n})^{N-N_1+1} \exp \left\{ -\beta U(s_n^N) \right\}}{V_{1,o}^{N_1+1} (V - V_{1,o})^{N-N_1+1} \exp \left\{ -\beta U(s_o^N) \right\}}$$

Moving a particle from box1 to box2

$$\mathcal{N}(N_1, V_1, s_1^{N_1}, s_2^{N-N_1}) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta \left[U(s_1^{N_1}) + U(s_2^{N-N_1}) \right] \right\}$$

acceptance rule:



$$\mathcal{N}(n) \propto \frac{V_1^{N_1-1} (V - V_1)^{N-(N_1-1)}}{(N_1-1)! (N - (N_1-1))!} \exp \left\{ -\beta [U_1(n) + U_2(n)] \right\}$$

$$\mathcal{N}(o) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta [U_1(o) + U_2(o)] \right\}$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\frac{V_1^{N_1-1} (V - V_1)^{N-(N_1-1)}}{(N_1-1)! (N - (N_1-1))!} \exp \left\{ -\beta [U_1(n) + U_2(n)] \right\}}{\frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \left\{ -\beta [U_1(o) + U_2(o)] \right\}}$$

Moving a particle from box1 to box2

$$\mathcal{N}(n) \propto \frac{V_1^{N_1-1} (V - V_1)^{N-(N_1-1)}}{(N_1-1)! (N - (N_1-1))!} \exp \{-\beta [U_1(n) + U_2(n)]\}$$

$$\mathcal{N}(o) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \{-\beta [U_1(o) + U_2(o)]\}$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\frac{V_1^{N_1-1} (V - V_1)^{N-(N_1-1)}}{(N_1-1)! (N - (N_1-1))!} \exp \{-\beta [U_1(n) + U_2(n)]\}}{\frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \{-\beta [U_1(o) + U_2(o)]\}}$$

Moving a particle from box1 to box2

$$\mathcal{N}(n) \propto \frac{V_1^{N_1-1} (V - V_1)^{N-(N_1-1)}}{(N_1-1)! (N - (N_1-1))!} \exp \{ -\beta [U_1(n) + U_2(n)] \}$$

$$\mathcal{N}(o) \propto \frac{V_1^{N_1} (V - V_1)^{N-N_1}}{N_1! (N - N_1)!} \exp \{ -\beta [U_1(o) + U_2(o)] \}$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N_1 (V - V_1)}{(N - N_1 + 1) V_1} \frac{\exp \{ -\beta U(s_n^N) \}}{\exp \{ -\beta U(s_o^N) \}}$$

Algorithm 17 (Basic Gibbs Ensemble Simulation)

```
PROGRAM mc_Gibbs
```

```
do icycl=1,ncycl
```

```
  ran=ranf()*(npart+nvol+nswap)
```

```
  if (ran.le.npart) then
```

```
    call mcmove
```

```
  else if (ran.le.(npart+nvol))
```

```
    call mcvol
```

```
  else
```

```
    call mcswap
```

```
  endif
```

```
  call sample
```

```
enddo
```

```
end
```

Gibbs ensemble simulation

perform `ncycl` MC cycles

detailed balance!!!

attempt to displace a particle

attempt to change the volume

attempt to swap a particle

sample averages

Algorithm 18 (Attempt to Change the Volume in the Gibbs Ensemble)

SUBROUTINE mcvol	attempt to change the volume
call toterg(box1,en1o)	energy old conf. box 1
call toterg(box2,en2o)	and 2 (box1: box length)
vol=box1**3	old volume box 1 and 2
vo2=v-vol	
lnvn=log(vo1/vo2)+	random walk in $\ln V_1/V_2$
+ (ranf()-0.5)*vmax	
v1n=v*exp(lnvn)/(1+exp(lnvn))	new volume box 1 and 2
v2n=v-v1n	
box1n=v1n**(1/3)	new box length box 1
box2n=v2n**(1/3)	new box length box 2
do i=1,npart	
if (ibox(i).eq.1) then	determine which box
fact=box1n/box1o	
else	
fact=box2n/box2o	
endif	
x(i)=x(i)*fact	rescale positions
enddo	
call toterg(box1n,en1n)	total energy box 1
call toterg(box2n,en2n)	total energy box 2
arg1=-beta*((en1n-en1o)+	
+ (npbox(1)+1)*log(v1n/v1o)/beta)	appropriate weight function
arg2=-beta*((en2n-en2o)+	acceptance rule (2.2.2)

<code>v1n=v*exp(lnvn)/(1+exp(lnvn))</code>	new volume box 1 and 2
<code>v2n=v-v1n</code>	
<code>box1n=v1n**(1/3)</code>	new box length box 1
<code>box2n=v2n**(1/3)</code>	new box length box 2
<code>do i=1,npart</code>	
<code>if (ibox(i).eq.1) then</code>	determine which box
<code>fact=box1n/box1o</code>	
<code>else</code>	
<code>fact=box2n/box2o</code>	
<code>endif</code>	
<code>x(i)=x(i)*fact</code>	rescale positions
<code>enddo</code>	
<code>call toterg(box1n,en1n)</code>	total energy box 1
<code>call toterg(box2n,en2n)</code>	total energy box 2
<code>arg1=-beta*((en1n-en1o)+</code>	
<code>+ (npbox(1)+1)*log(v1n/v1o)/beta)</code>	appropriate weight function
<code>arg2=-beta*((en2n-en2o)+</code>	acceptance rule (8.3.3)
<code>+ (npbox(2)+1)*log(v2n/v2o)/beta)</code>	
<code>if (ranf().gt.exp(arg1+arg2)) then</code>	
<code>do i=1,npart</code>	REJECTED
<code>if (ibox(i).eq.) then</code>	determine which box
<code>fact=box1o/box1n</code>	
<code>else</code>	
<code>fact=box2o/box2n</code>	
<code>endif</code>	
<code>x(i)=x(i)*fact</code>	restore old configuration
<code>enddo</code>	
<code>endif</code>	
<code>return</code>	
<code>end</code>	

Algorithm 19 (Attempt to Swap a Particle between the Two Boxes)

```
SUBROUTINE mcswap

if (ranf().lt.0.5) then
  in=1
  out=2
else
  in=2
  out=1
endif
xn=ranf()*box(in)
call ener(xn,enn,in)
w(in)=w(in)+vol(in)*
+ exp(-beta*enn)/(npbox(in)+1)
if (npbox(out).eq.0) return
ido=0
do while (ido.ne.out)
  o=int(npart*ranf())+1
  ido=ibox(o)
enddo
call ener(x(o),eno,out)
arg=exp(-beta*(enn-eno +
+ log(vol(out)*(npbox(in)+1)/
+ (vol(in)*npbox(out))))/beta)
if (ranf().lt.arg) then
  x(o)=xn
```

attempts to swap a particle
between the two boxes
which box to add or remove

new particle at a random position
energy new particle in box *in*
update chemical potential (8.3.5)

if box empty return
find a particle to be removed

energy particle *o* in box *out*

acceptance rule (8.3.4)

add new particle to box *in*

```

        out=2
    else
        in=2
        out=1
    endif
    xn=ranf()*box(in)
    call ener(xn,enn,in)
    w(in)=w(in)+vol(in)*
+ exp(-beta*enn)/(npbox(in)+1)
    if (npbox(out).eq.0) return
    ido=0
    do while (ido.ne.out)
        o=int(npart*ranf())+1
        ido=ibox(o)
    enddo
    call ener(x(o),eno,out)
    arg=exp(-beta*(enn-eno +
+ log(vol(out)*(npbox(in)+1)/
+ (vol(in)*npbox(out)))/beta))
    if (ranf().lt.arg) then
        x(o)=xn
        ibox(o)=in
        nbox(out)=npbox(out)-1
        nbox(in)=npbox(in)+1
    endif
    return
end

```

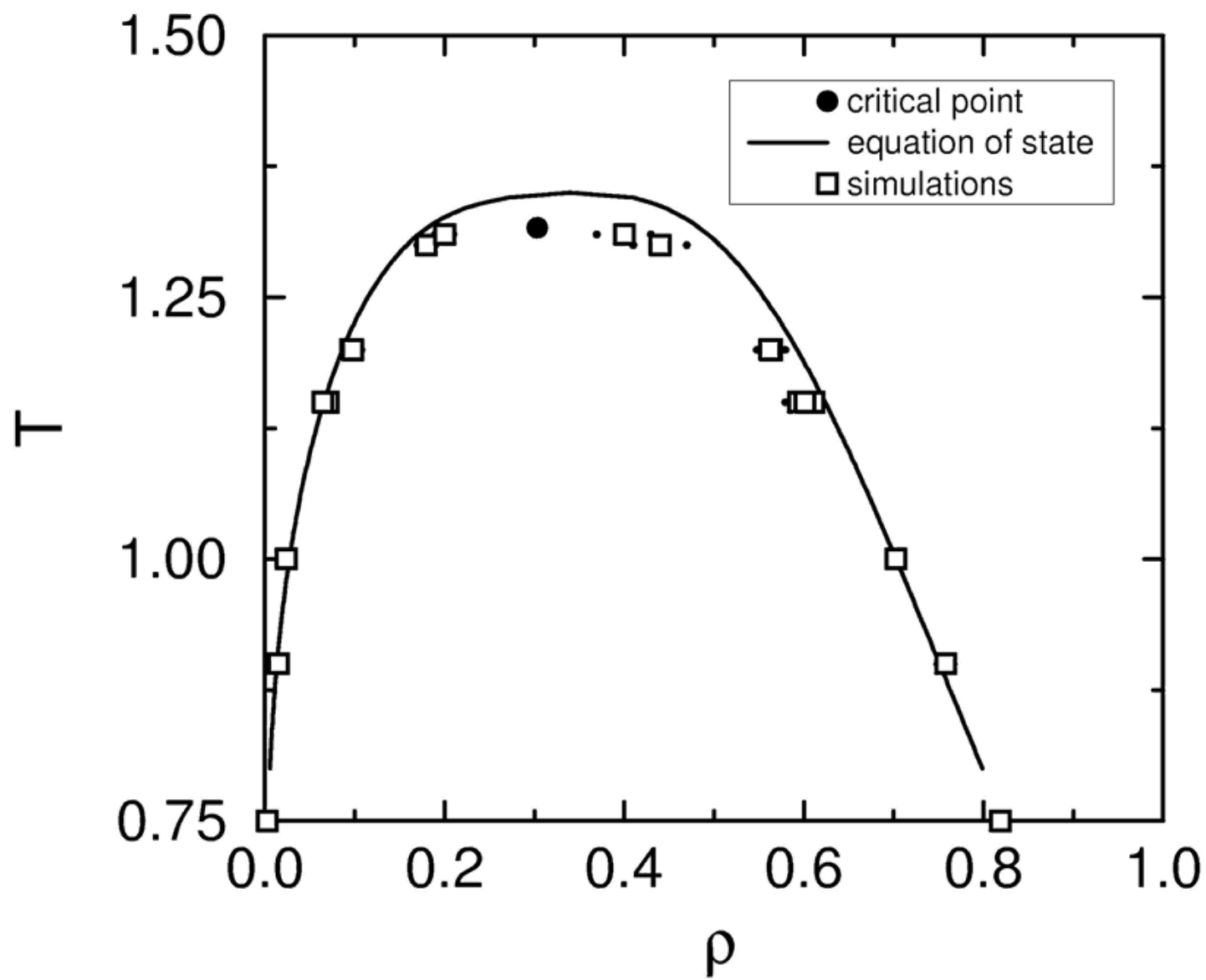
new particle at a random position
 energy new particle in box in
 update chemical potential (8.3.5)

if box empty return
 find a particle to be removed

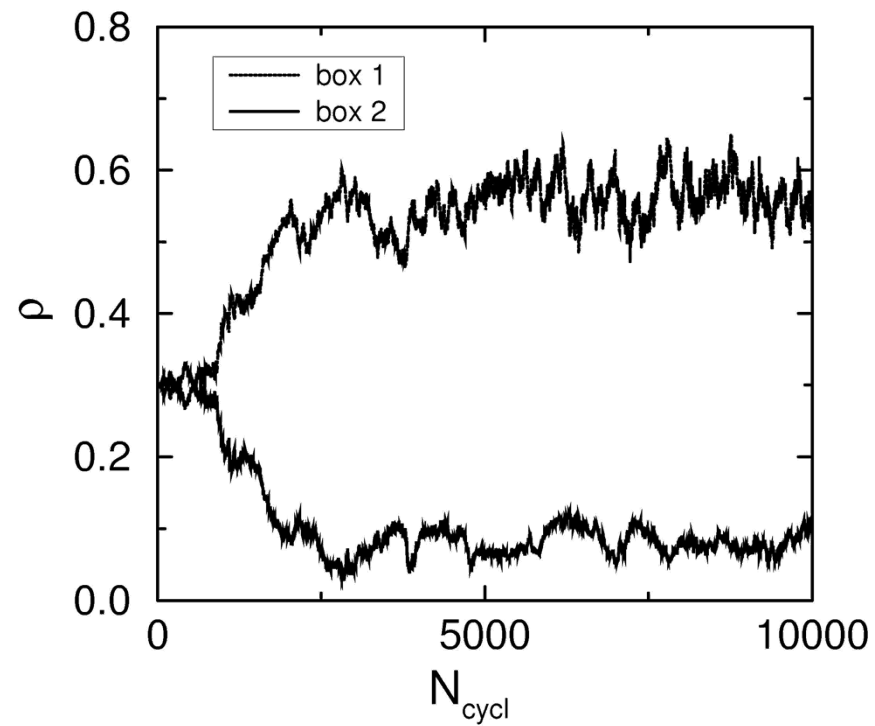
energy particle o in box out

acceptance rule (8.3.4)

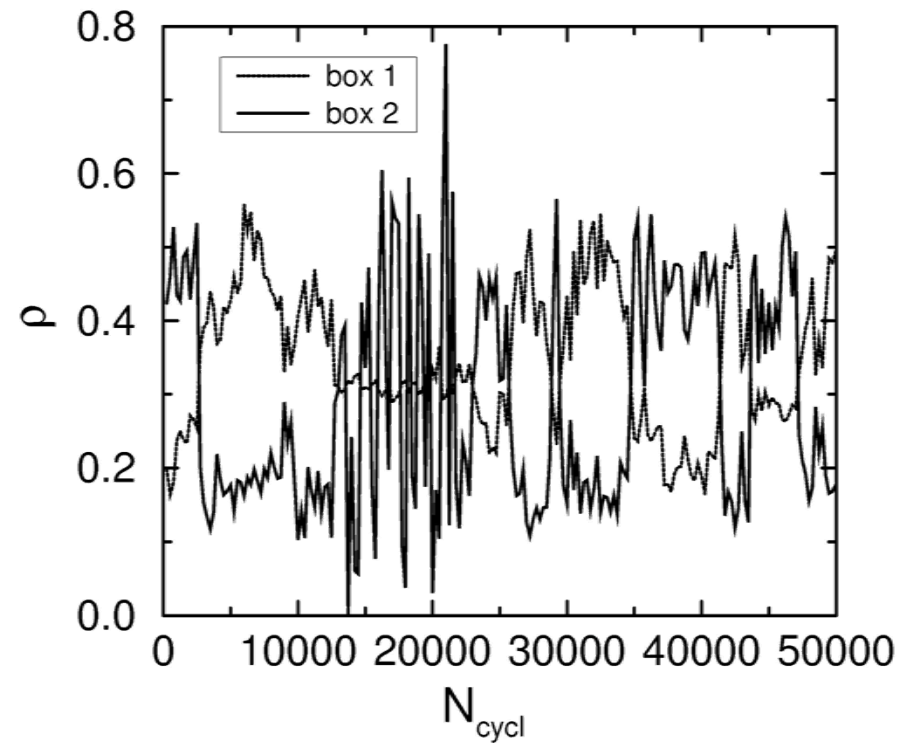
add new particle to box in



Analyzing the results (1)

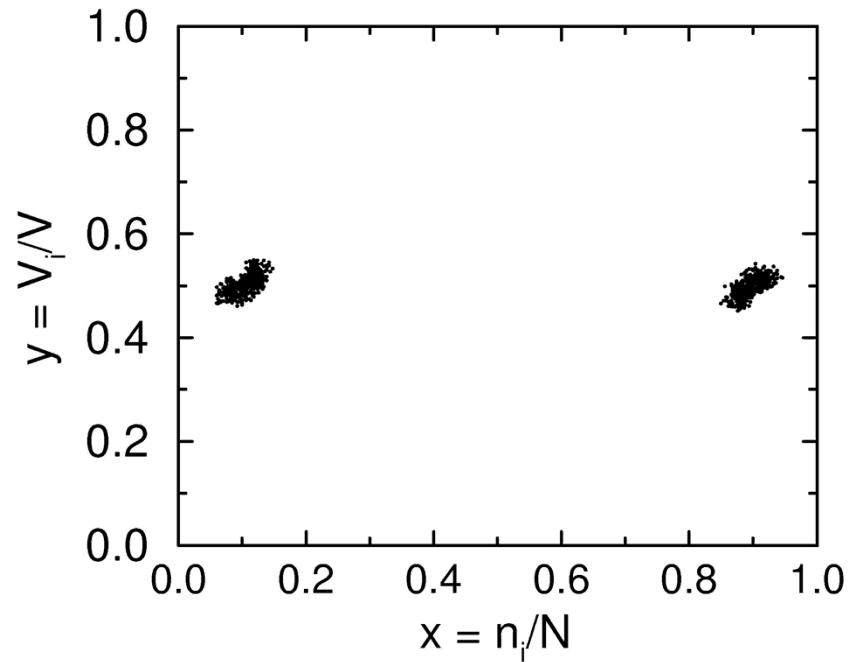


well below T_c

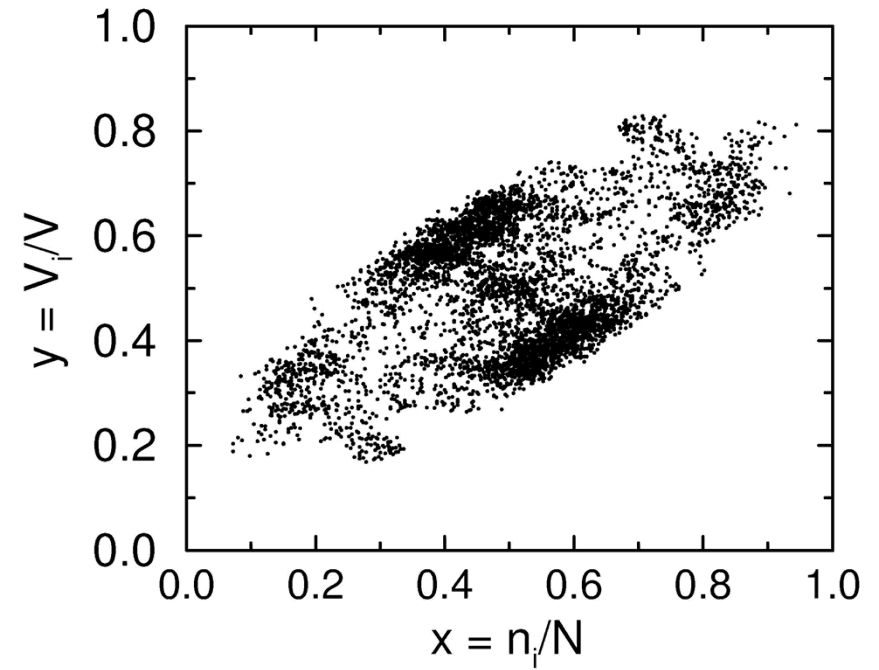


approaching T_c

Analyzing the results (2)

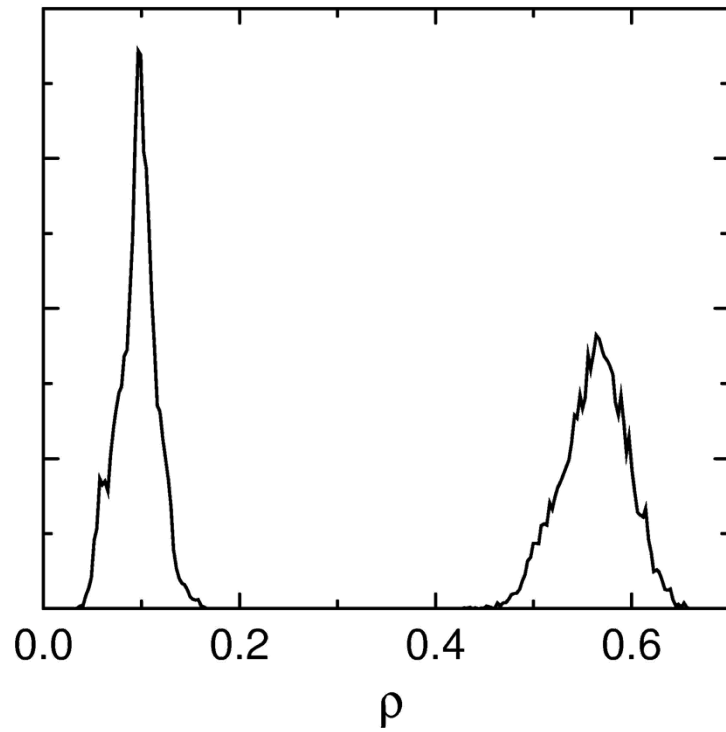


well below T_c

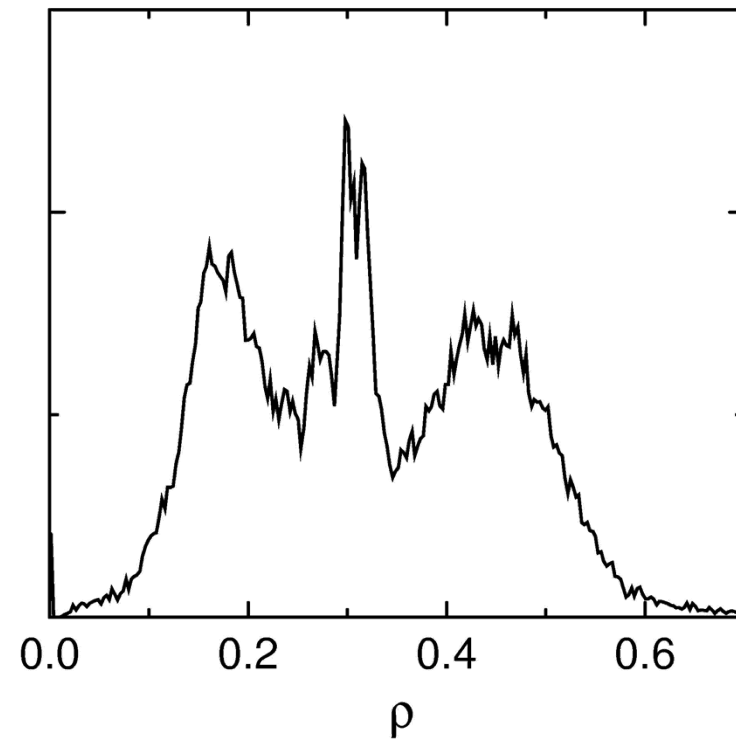


approaching T_c

Analyzing the results (3)



well below T_c



approaching T_c

Advantages

- single simulation to study phase coexistence: system „finds“ the densities of coexisting phases
- free energies/chem. potentials need not be calculated
- significant reduction of computer time

Disadvantages

-
- only for vapor-liquid and liquid-liquid coexistence
- not very successful for dense phases (particle insertion!)