

Monte Carlo in different ensembles

Chapter 5

NVT ensemble

NPT ensemble

Grand-canonical (μVT) ensemble

Exotic ensembles: semigrand, isotension

Different Ensembles

Ensemble	Name	Constant (Imposed)	Fluctuating (Measured)	Thermo- dynamics
NVT	Canonical	N,V,T	P	F
NPT	Isobaric-isothermal	N,P,T	V	G
μ VT	Grand-canonical	μ ,V,T	N	Ω

NVT ensemble: Statistical Thermodynamics

Partition function

$$Q_{NVT} = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]$$

Ensemble average

$$\langle A \rangle_{NVT} = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]$$

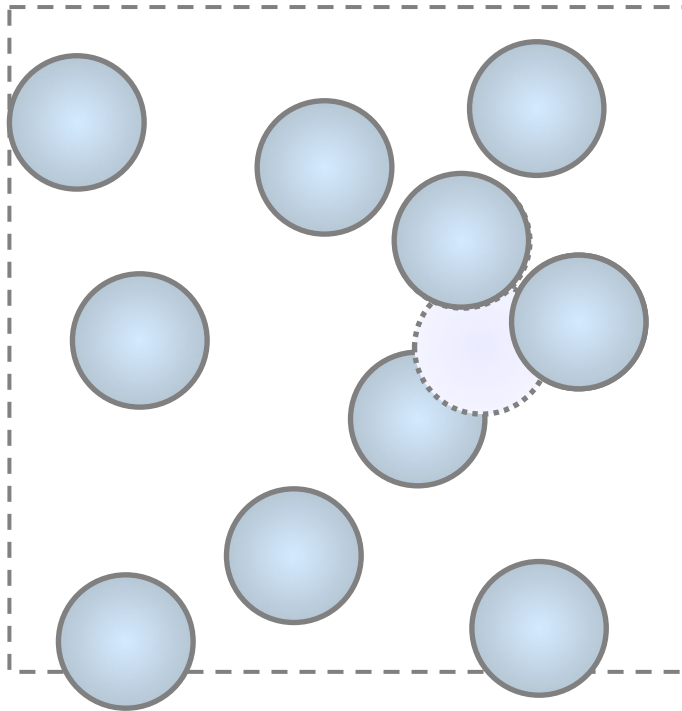
Probability to find a particular configuration

$$N(\mathbf{r}^N) = \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}'^N \delta(\mathbf{r}'^N - \mathbf{r}^N) \exp[-\beta U(\mathbf{r}'^N)] \propto \exp[-\beta U(\mathbf{r}^N)]$$

Free energy

$$\beta F = -\ln(Q_{NVT})$$

Monte Carlo sampling



Samples configuration space according to the Boltzmann distribution

$$\langle A \rangle = \frac{\int dr^N A e^{-U(r^N)/k_B T}}{\int dr^N e^{-U(r^N)/k_B T}}$$

**Monte Carlo
simulation**

Ensemble average

$$\begin{aligned}
 \langle A \rangle_{NVT} &= \frac{1}{Q_{NVT}} \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)] \\
 &= \int d\mathbf{r}^N A(\mathbf{r}^N) P(\mathbf{r}^N) = \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) P(\mathbf{r}^N)}{\int d\mathbf{r}^N P(\mathbf{r}^N)} \\
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) C \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N C \exp[-\beta U(\mathbf{r}^N)]} = \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]}
 \end{aligned}$$

$$P(\mathbf{r}^N) = \frac{\exp[-\beta U(\mathbf{r}^N)]}{Q_{NVT} \Lambda^{3N} N!}$$

Generate configuration using MC:

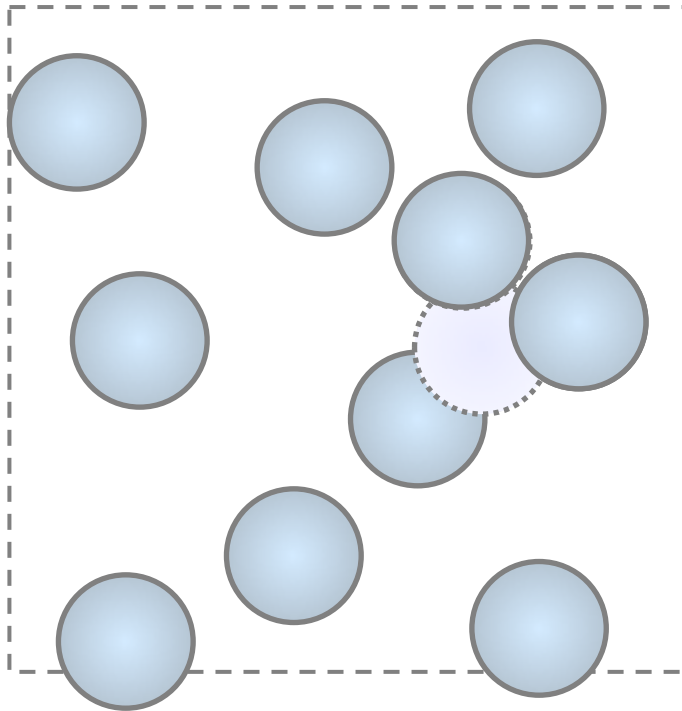
$$\{\mathbf{r}_1^N, \mathbf{r}_2^N, \mathbf{r}_3^N, \mathbf{r}_4^N \dots, \mathbf{r}_M^N\} \quad \bar{A} = \frac{1}{M} \sum_{i=1}^M A(\mathbf{r}_i^N) \approx \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) P^{MC}(\mathbf{r}^N)}{\int d\mathbf{r}^N P^{MC}(\mathbf{r}^N)}$$

with

$$P^{MC}(\mathbf{r}^N) = C^{MC} \exp[-\beta U(\mathbf{r}^N)]$$

$$\begin{aligned}
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) C^{MC} \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N C^{MC} \exp[-\beta U(\mathbf{r}^N)]} \\
 &= \frac{\int d\mathbf{r}^N A(\mathbf{r}^N) \exp[-\beta U(\mathbf{r}^N)]}{\int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]}
 \end{aligned}$$

Monte Carlo sampling



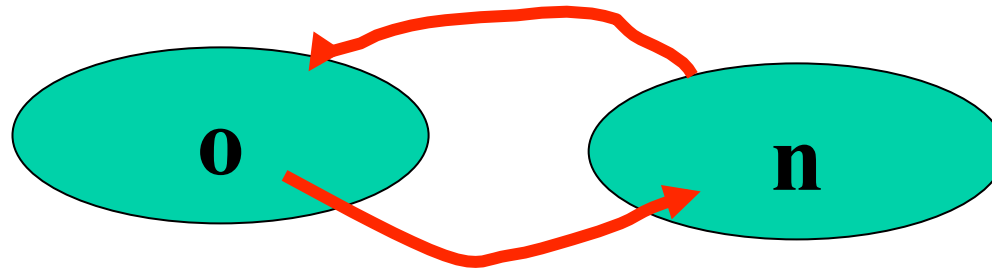
**Monte Carlo
simulation**

Samples configuration space
according to the Boltzmann
distribution

$$\langle A \rangle = \frac{\int dr^N A e^{-U(r^N)/k_B T}}{\int dr^N e^{-U(r^N)/k_B T}}$$

How do we derive monte carlo
algorithm?

Detailed balance



$$K(o \rightarrow n) = K(n \rightarrow o)$$

$$K(o \rightarrow n) = N(o) \times \alpha(o \rightarrow n)$$

$$K(n \rightarrow o) = N(n) \times \alpha(n \rightarrow o)$$

Symmetric generation probability

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n) \times \alpha(n \rightarrow o)}{N(o) \times \alpha(o \rightarrow n)} = \frac{N(n)}{N(o)}$$

NVT-ensemble

$$N(n) \propto \exp[-\beta U(n)]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n)}{N(o)}$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \exp[-\beta [U(n) - U(o)]]$$

Choice of metropolis et al.

$$\text{acc}(o \rightarrow n) = \min[1, \exp(-\beta \Delta U(o \rightarrow n))]$$

Algorithm 2 (Attempt to Displace a Particle)

<pre>SUBROUTINE mcmove o=int(ranf()*npart)+1 call ener(x(o), eno) xn=x(o)+(ranf()-0.5)*delx call ener(xn, enn) if (ranf().lt.exp(-beta + *(enn-eno)) x(o)=xn return end</pre>	attempts to displace a particle select a particle at random energy old configuration give particle random displacement energy new configuration acceptance rule (3.2.1) accepted: replace $x(o)$ by xn
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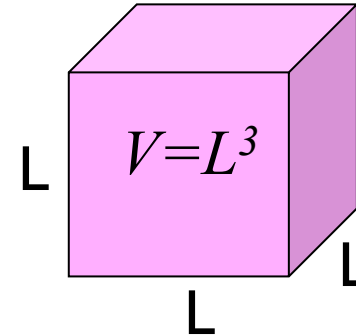
Comments to this algorithm:

1. Subroutine `ener` calculates the energy of a particle at the given position.
2. Note that, if a configuration is rejected, the old configuration is retained.
3. The `ranf()` is a random number uniform in $[0, 1]$.

Scaled coordinates

Partition function

$$Q_{NVT} = \frac{1}{\Lambda^{3N} N!} \int d\mathbf{r}^N \exp[-\beta U(\mathbf{r}^N)]$$



Scaled coordinates

$$\mathbf{s}_i = \mathbf{r}_i / L$$

This gives for the partition function

$$\begin{aligned} Q_{NVT} &= \frac{L^{3N}}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)] \\ &= \frac{V^N}{\Lambda^{3N} N!} \int d\mathbf{s}^N \exp[-\beta U(\mathbf{s}^N; L)] \end{aligned}$$

The energy depends on the real coordinates

Measuring the pressure

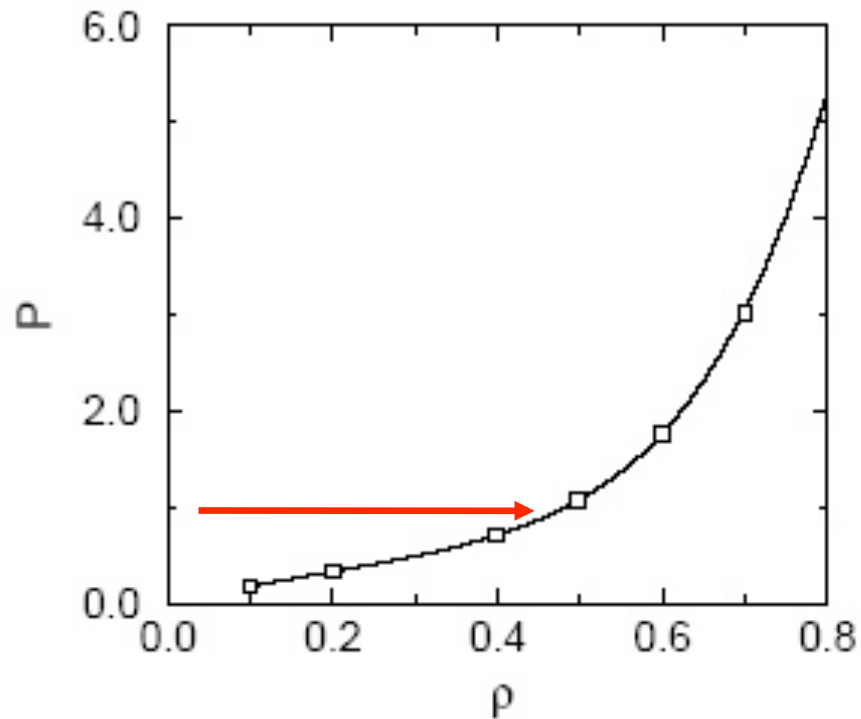
The pressure is the thermodynamic derivative of F

$$\begin{aligned}\beta P &= -\left(\frac{\partial \beta F}{\partial V}\right)_{N,T} = \left(\frac{\partial \ln Q_{NVT}}{\partial V}\right)_{N,T} = \frac{1}{Q_{NVT}} \left(\frac{\partial Q_{NVT}}{\partial V}\right)_{N,T} \\ &= \frac{\frac{N V^{N-1}}{\Lambda^{3N} N!} \int ds^N e^{-\beta U(r^N)} - \frac{V^N}{\Lambda^{3N} N!} \int ds^N e^{-\beta U(r^N)} \frac{d\beta U(r^N)}{dV}}{\frac{V^N}{\Lambda^{3N} N!} \int ds^N e^{-\beta U(s^N)}}\end{aligned}$$

this reduces to the virial pressure

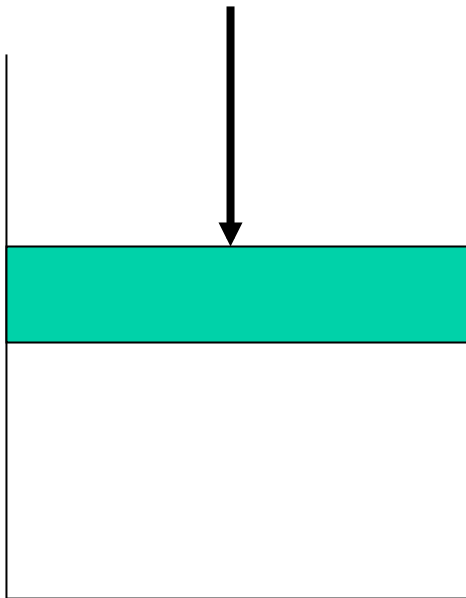
$$\beta P = \rho - \left\langle \frac{d\beta U(r^N)}{dV} \right\rangle$$

Constant pressure simulation?



- what if we are interested in one specific pressure P ?
- we need large part of the equation of state to find right volume V
 - needs many simulations.
- Better idea: constant pressure simulation.

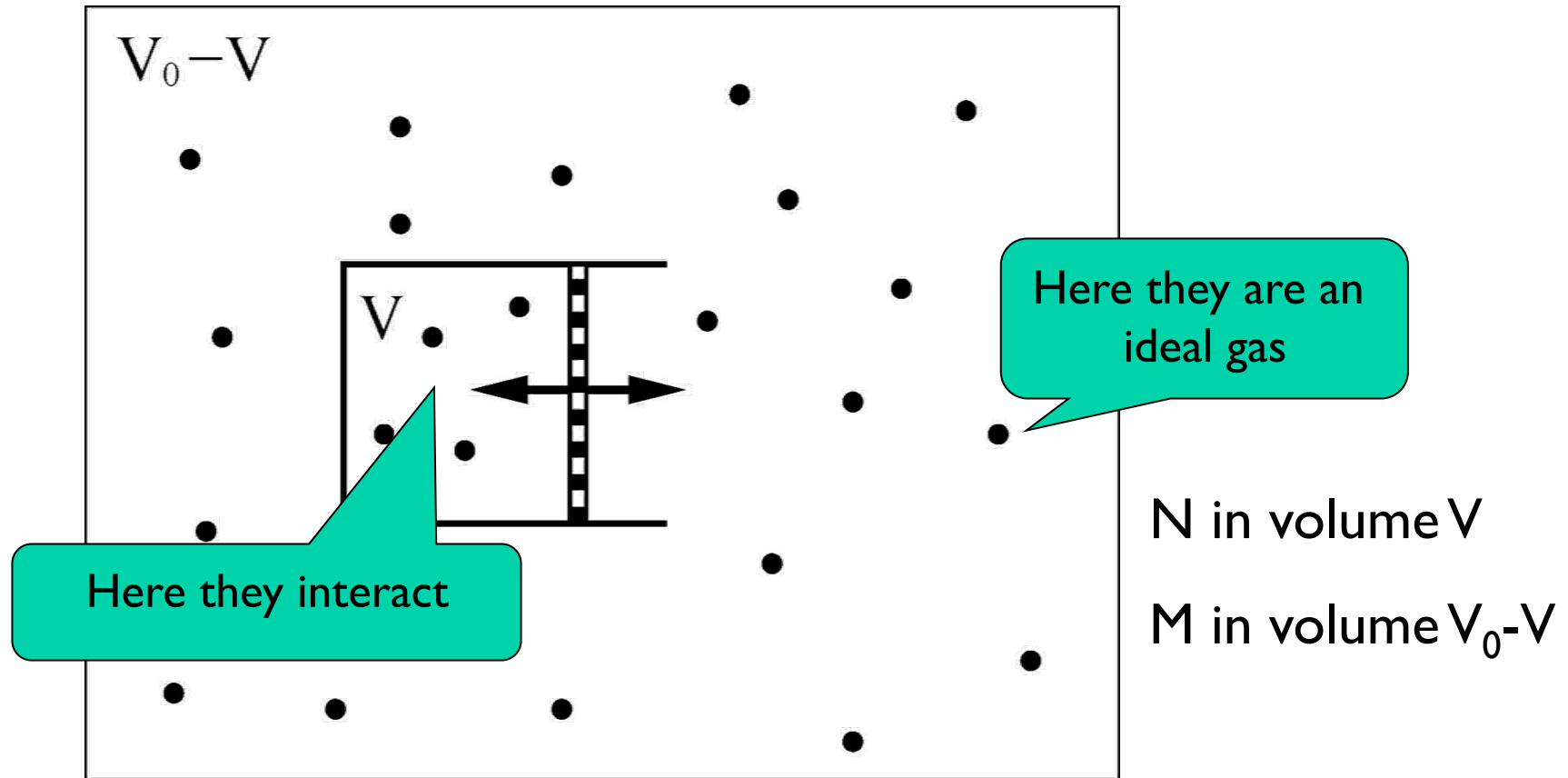
NPT ensemble



We control the

- Temperature (T)
- Pressure (P)
- Number of particles (N)

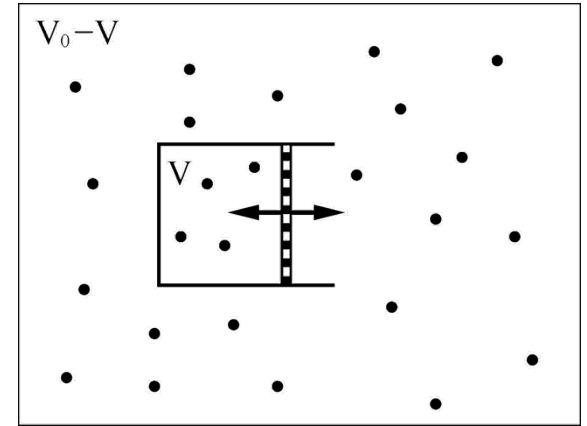
The perfect simulation ensemble



What is the statistical thermodynamics of this ensemble?

The perfect ensemble: partition function

$$Q_{NVT} = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$



$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \int ds^{M-N} \exp \left[-\beta U_0(s^{M-N}; L) \right] \frac{V^N}{\Lambda^{3N} N!} \times \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

To get the total partition function of this system, we have to integrate over all possible volumes:

$$Q_{MV_0, N, T} = \int dV \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

Now let us take the following limits:

$$\left. \begin{array}{l} M \rightarrow \infty \\ V_0 \rightarrow \infty \end{array} \right\} \rho = \frac{M}{V} \rightarrow \text{constant}$$

As the particles are an ideal gas in the big reservoir we have:

$$\rho = \beta P$$

big reservoir acts as manostat

$$Q_{MV_0, N, T} = \int dV \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp[-\beta U(s^N; L)]$$

We have

$$(V_0 - V)^{M-N} = V_0^{M-N} (1 - V/V_0)^{M-N} \approx V_0^{M-N} \exp[-(M-N)V/V_0]$$

$$(V_0 - V)^{M-N} \approx V_0^{M-N} \exp[-\rho V] = V_0^{M-N} \exp[-\beta P V]$$

This gives:

To make the partition function dimensionless (not trivial)

$$Q_{NPT} = \frac{\beta P}{N! \Lambda^{3N}} \int dV \exp[-\beta P V] V^N \int ds^N \exp[-\beta U(s^N; L)]$$

NPT Ensemble

Partition function:

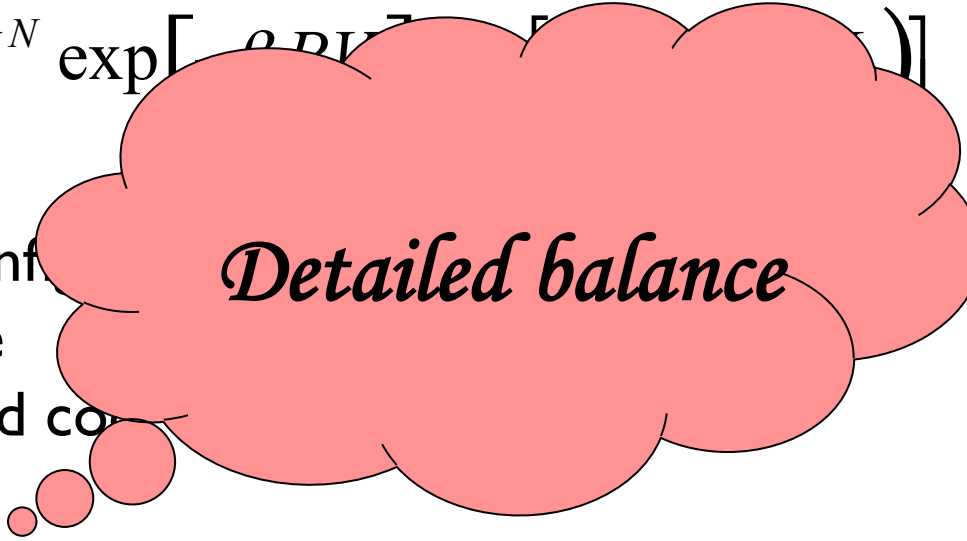
$$Q_{NPT} = \frac{\beta P}{N! \Lambda^{3N}} \int dV \exp[-\beta P V] V^N \int ds^N \exp[-\beta U(s^N; L)]$$

Probability to find a particular configuration:

$$N_{NPT}(V, \mathbf{s}^N) \propto V^N \exp[-\beta P V - \beta U(\mathbf{s}^N; L)]$$

Sample a particular configuration

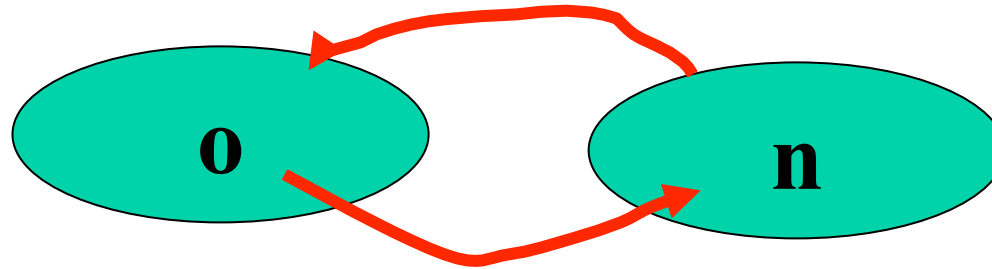
- change of volume
- change of reduced coordinates



Detailed balance

Acceptance rules ??

Detailed balance



$$K(o \rightarrow n) = K(n \rightarrow o)$$

$$K(o \rightarrow n) = N(o) \times \alpha(o \rightarrow n) \times \text{acc}(o \rightarrow n)$$

$$K(n \rightarrow o) = N(n) \times \alpha(n \rightarrow o) \times \text{acc}(n \rightarrow o)$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n) \times \alpha(n \rightarrow o)}{N(o) \times \alpha(o \rightarrow n)} = \frac{N(n)}{N(o)}$$

NPT-ensemble

$$N_{NPT} (V, \mathbf{s}^N) \propto V^N \exp[-\beta PV] \exp[-\beta U(\mathbf{s}^N; L)]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n)}{N(o)}$$

Suppose we change the position of a randomly selected particle

$$\begin{aligned} \frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} &= \frac{V^N \exp[-\beta PV] \exp[-\beta U(\mathbf{s}_n^N; L)]}{V^N \exp[-\beta PV] \exp[-\beta U(\mathbf{s}_o^N; L)]} \\ &= \frac{\exp[-\beta U(\mathbf{s}_n^N; L)]}{\exp[-\beta U(\mathbf{s}_o^N; L)]} = \exp \left\{ -\beta [U(n) - U(o)] \right\} \end{aligned}$$

NPT-ensemble

$$N_{NPT} (V, \mathbf{s}^N) \propto V^N \exp[-\beta P V] \exp[-\beta U(\mathbf{s}^N; L)]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n)}{N(o)}$$

Suppose we change the **volume** of the system

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{V_n^N \exp[-\beta P V_n] \exp[-\beta U(\mathbf{s}^N; L_n)]}{V_o^N \exp[-\beta P V_o] \exp[-\beta U(\mathbf{s}^N; L_o)]}$$

$$= \left(\frac{V_n}{V_o} \right)^N \exp[-\beta P (V_n - V_o)] \exp\{-\beta [U(n) - U(o)]\}$$

Algorithm: NPT

- Randomly change the position of a particle
- Randomly change the volume

Algorithm 10 (Basic NPT-Ensemble Simulation)

PROGRAM mc_npt	basic NPT ensemble simulation
do icycl=1,ncycl	perform ncycl MC cycles
ran=ranf()*(npart+1)+1	
if (ran.le.npart) then	
call mcmove	perform particle displacement
else	
call mcvol	perform volume change
endif	
if (mod(icycl,nsamp).eq.0)	
+ call sample	sample averages
enddo	
end	

Algorithm 2 (Attempt to Displace a Particle)

<pre>SUBROUTINE mcmove o=int(ranf()*npart)+1 call ener(x(o), eno) xn=x(o)+(ranf()-0.5)*delx call ener(xn, enn) if (ranf().lt.exp(-beta + *(enn-eno)) x(o)=xn return end</pre>	attempts to displace a particle select a particle at random energy old configuration give particle random displacement energy new configuration acceptance rule (3.2.1) accepted: replace $x(o)$ by xn
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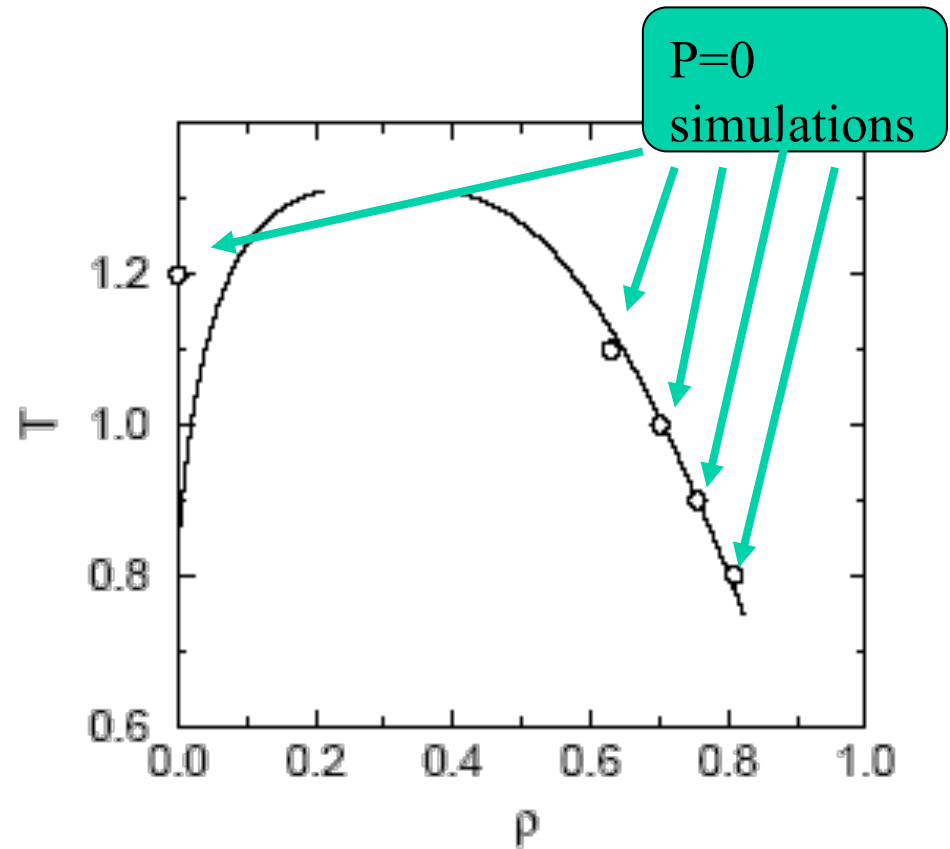
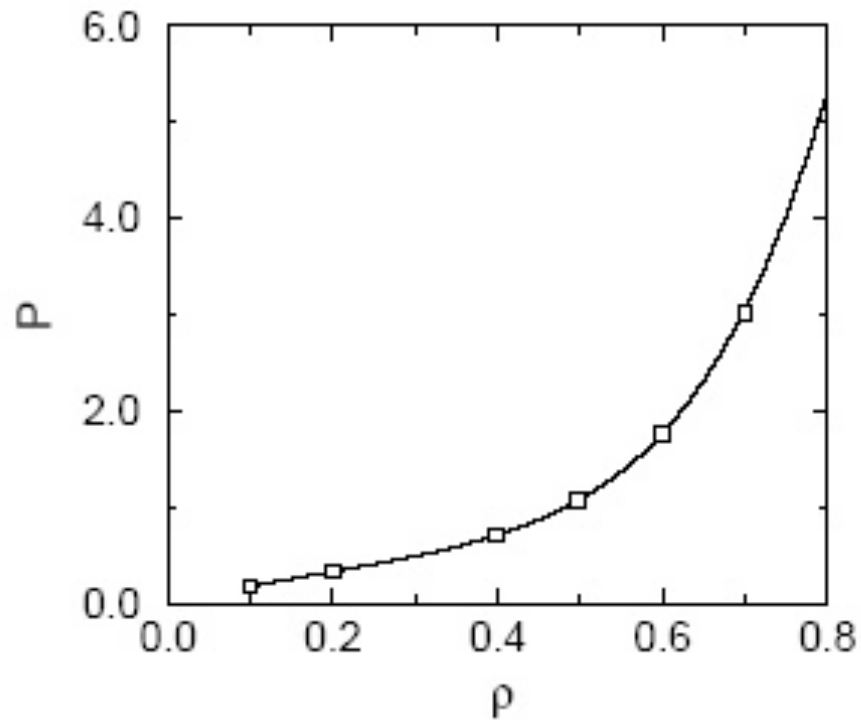
Comments to this algorithm:

1. Subroutine `ener` calculates the energy of a particle at the given position.
2. Note that, if a configuration is rejected, the old configuration is retained.
3. The `ranf()` is a random number uniform in $[0, 1]$.

Algorithm 11 (Attempt to Change the Volume)

SUBROUTINE mcvol	attempts to change the volume
call toterg(box, eno)	total energy old conf.
vo=box**3	determine old volume
lnvn=log(vo) + (ranf() - 0.5) * vmax	perform random walk in $\ln V$
vn=exp(lnvn)	
boxn=vn**(1/3)	new box length
do i=1, npart	
x(i) = x(i) * boxn / box	rescale center of mass
enddo	
call toterg(boxn, enn)	total energy new conf.
arg = -beta * ((enn - eno) + p * (vn - vo)	
+ - (npart + 1) * log(vn / vo) / beta)	appropriate weight function!
if (ranf() .gt. exp(arg)) then	acceptance rule (5.2.3)
do i=1, npart	REJECTED
x(i) = x(i) * box / boxn	restore the old positions
enddo	
endif	
return	
end	

NPT simulations



Setting pressure to zero will give liquid coexistence except close to critical point: why?

Measured and Imposed Pressure

- Imposed pressure P
- Measured pressure $\langle P_V \rangle$ from virial

$$P_V = - \left(\frac{\partial F}{\partial V} \right)_{N,T}$$

$$p(V) = \frac{\exp[-\beta(F(V) + PV)]}{Q_{NPT}}$$

$$Q_{NPT} = \beta P \int dV \exp[-\beta(F(V) + PV)]$$

In fact this is βG

$$\langle P_V \rangle = - \left\langle \left(\frac{\partial F}{\partial V} \right)_{N,T} \right\rangle$$

Measured and Imposed Pressure

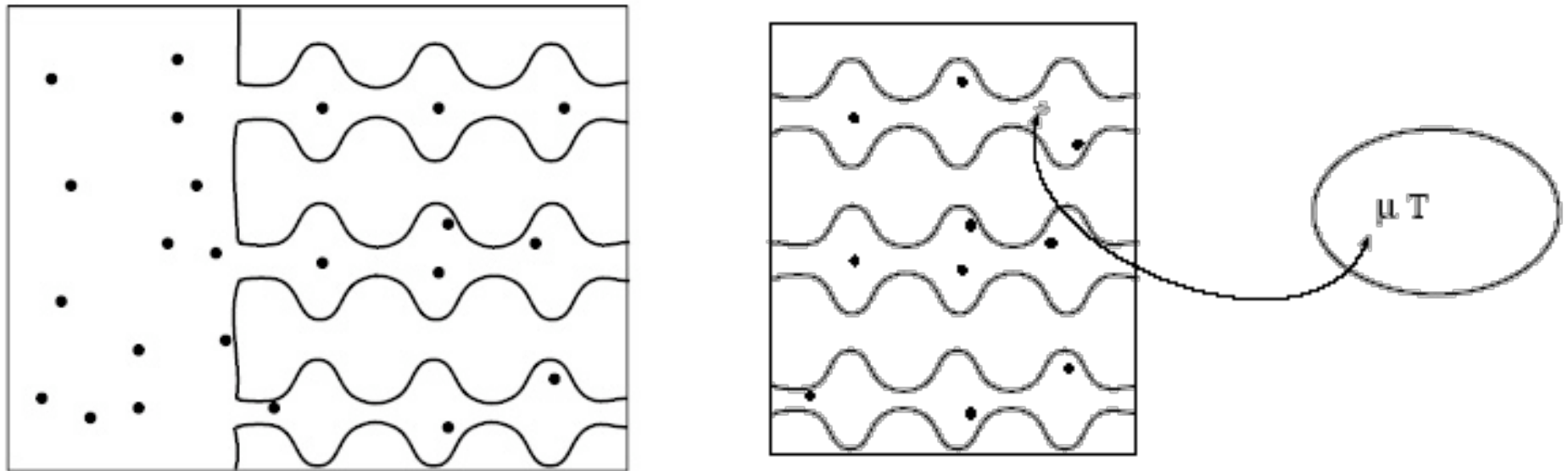
- Partial integration $\int_a^b f dg = [fg]_a^b - \int_a^b g df$
- For $V=0$ and $V=\infty$ $\exp[-\beta(F(V) + PV)] = 0$
- Therefore,

$$\langle P_V \rangle = \frac{\beta P}{Q(NPT)} \int dV \frac{\exp[-\beta PV]}{\beta} \frac{\partial \exp[-\beta F(V)]}{\partial V}$$

$$\langle P_V \rangle = \frac{\beta P}{Q(NPT)} \int dV P \exp[-\beta(F(V) + PV)] = P$$

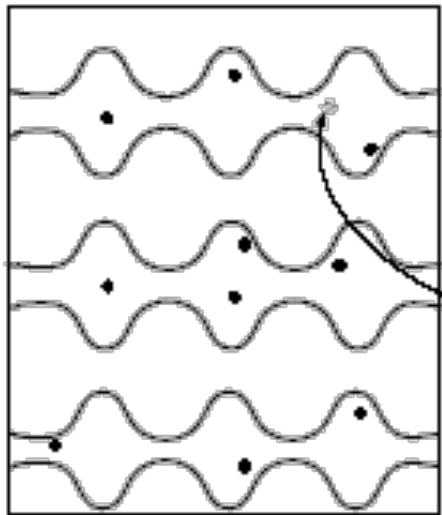
Grand-canonical ensemble

Suppose we have a gas in a porous material



What are the equilibrium conditions?

Grand-canonical ensemble

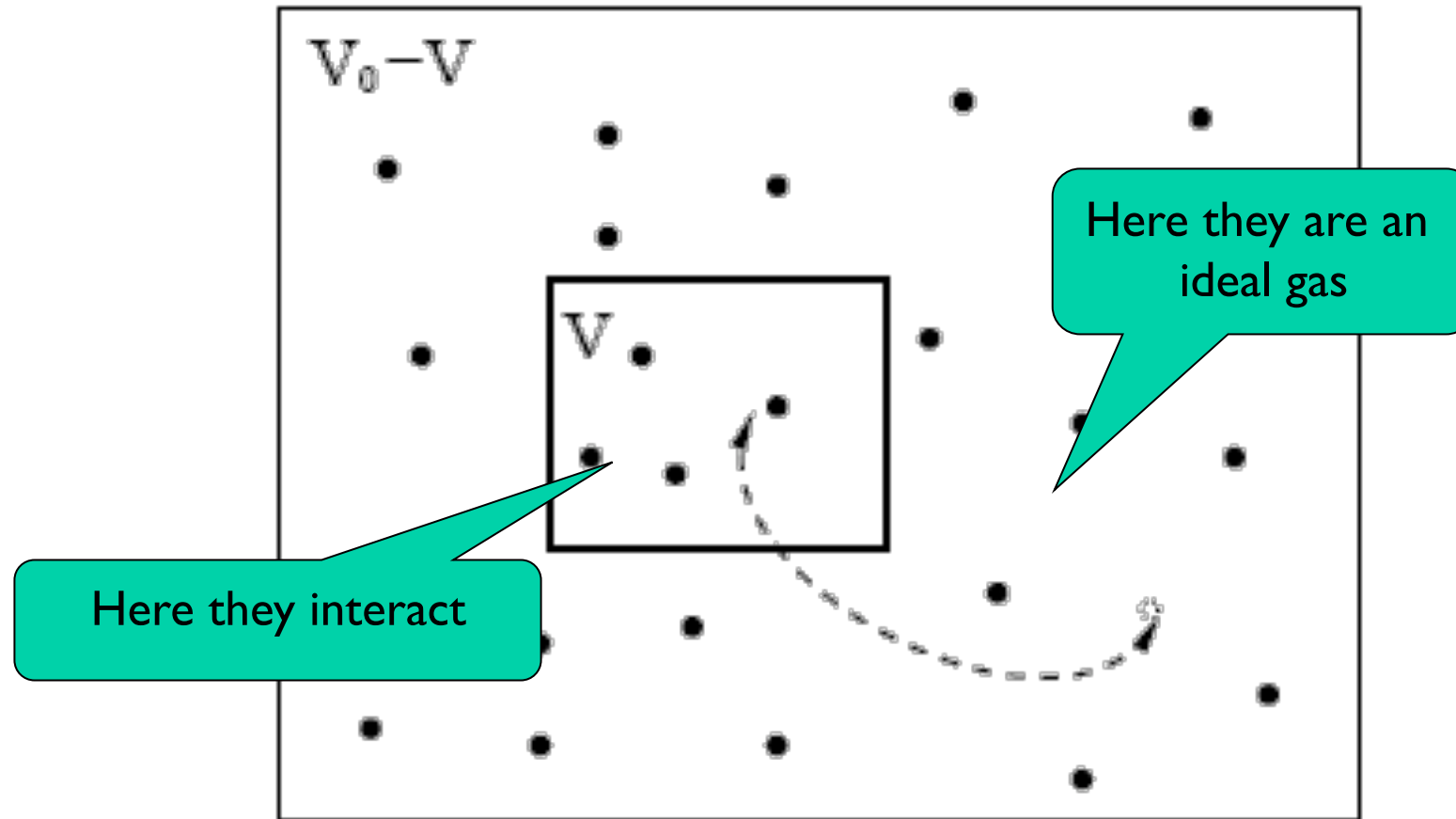


We impose:

- Temperature (T)
- Chemical potential (μ)
- Volume (V)

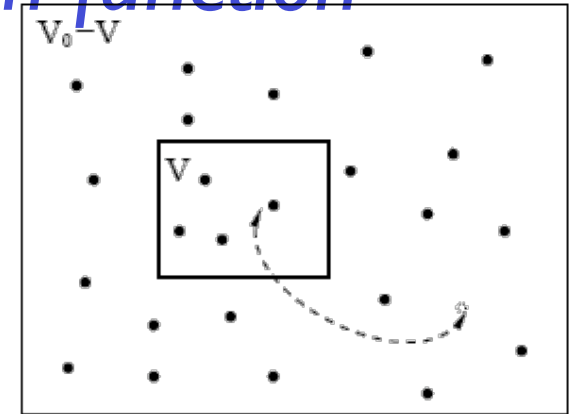
But **NOT** pressure

The μ perfect ensemble



What is the statistical thermodynamics of this ensemble?

The μ perfect ensemble: partition function



$$Q_{NVT} = \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \int ds^{M-N} \exp \left[-\beta U_0(s^{M-N}; L) \right] \frac{V^N}{\Lambda^{3N} N!} \\ \times \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

$$Q_{MV_0, NV, T} = \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

To get the total partition function of this system, we have to sum over all possible numbers of particles N

$$Q_{MV_0, N, T} = \sum_{N=0}^{N=M} \frac{(V_0 - V)^{M-N}}{\Lambda^{3M-N} (M-N)!} \frac{V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

Now let us take the limits:

$$\left. \begin{array}{l} M \rightarrow \infty \\ V_0 \rightarrow \infty \end{array} \right\} \rho = \frac{M}{V} \rightarrow \text{constant}$$

As the particles are an ideal gas in the big reservoir we have:

$$\mu = k_B T \ln (\Lambda^3 \rho)$$

$$Q_{\mu VT} = \sum_{N=0}^{N=\infty} \frac{\exp(\beta \mu N) V^N}{\Lambda^{3N} N!} \int ds^N \exp \left[-\beta U(s^N; L) \right]$$

μVT Ensemble

Partition function:

$$Q_{\mu VT} = \sum_{N=0}^{N=\infty} \frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!} \int ds^N \exp[-\beta U(s^N; L)]$$

Probability to find a particular configuration:

$$N_{\mu VT}(V, s^N) \propto \frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!} \exp[-\beta U(s^N; L)]$$

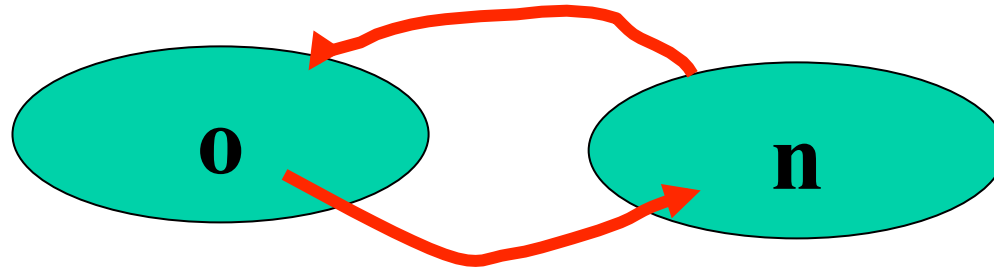
Sample a particular configuration

- Change of the number of particles
- Change of reduced coordinates

Detailed balance

Acceptance rules ??

Detailed balance



$$K(o \rightarrow n) = K(n \rightarrow o)$$

$$K(o \rightarrow n) = N(o) \times \alpha(o \rightarrow n) \times \text{acc}(o \rightarrow n)$$

$$K(n \rightarrow o) = N(n) \times \alpha(n \rightarrow o) \times \text{acc}(n \rightarrow o)$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n) \times \alpha(n \rightarrow o)}{N(o) \times \alpha(o \rightarrow n)} = \frac{N(n)}{N(o)}$$

μVT-ensemble

$$N_{\mu VT} (V, \mathbf{s}^N) \propto \frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!} \exp \left[-\beta U(\mathbf{s}^N; L) \right]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n)}{N(o)}$$

Suppose we change the position of a randomly selected particle

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{\cancel{\frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!}} \exp \left[-\beta U(\mathbf{s}_n^N; L) \right]}{\cancel{\frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!}} \exp \left[-\beta U(\mathbf{s}_o^N; L) \right]}$$

$$= \exp \left\{ -\beta [U(n) - U(o)] \right\}$$

μVT -ensemble

$$N_{\mu VT}(V, \mathbf{s}^N) \propto \frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!} \exp[-\beta U(\mathbf{s}^N; L)]$$

$$\frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} = \frac{N(n)}{N(o)}$$

Suppose we change the **number of particles** of the system

$$\begin{aligned} \frac{\text{acc}(o \rightarrow n)}{\text{acc}(n \rightarrow o)} &= \frac{\frac{\exp(\beta\mu(N+1)) V^{N+1}}{\Lambda^{3N+3} (N+1)!} \exp[-\beta U(\mathbf{s}^{N+1}; L_n)]}{\frac{\exp(\beta\mu N) V^N}{\Lambda^{3N} N!} \exp[-\beta U(\mathbf{s}^N; L_o)]} \\ &= \frac{\exp(\beta\mu) V}{\Lambda^3 (N+1)} \exp[-\beta \Delta U] \end{aligned}$$

Algorithm 12 (Basic Grand-Canonical Ensemble Simulation)

PROGRAM mc_gc	basic μ VT ensemble simulation
do icycl=1,ncycl	perform ncycl MC cycles
ran=int(ranf()*(npav+nexc))+1	
if (ran.le.npart) then	
call mcmove	displace a particle
else	
call mcexc	exchange a particle with the reservoir
endif	
if (mod(icycl,nsamp).eq.0)	
+ call sample	sample averages
enddo	
end	

Comments to this algorithm:

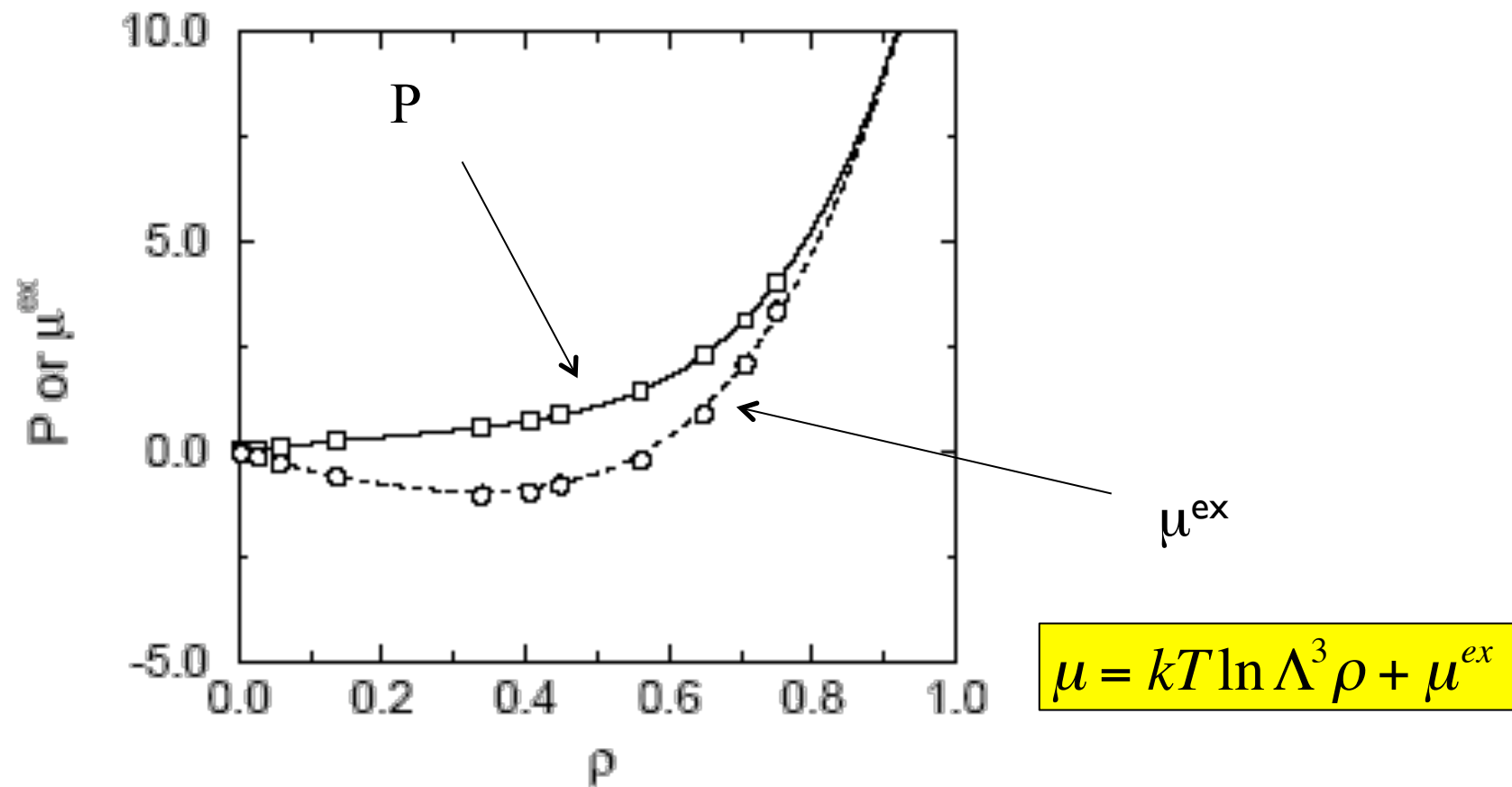
1. This algorithm ensures that, after each MC step, detailed balance is obeyed. Per cycle we perform on average $npav$ attempts^b to displace particles and $nexc$ attempts to exchange particles with the reservoir.
2. Subroutine `mcmove` attempts to displace a particle (Algorithm 2), subroutine `mcexc` attempts to exchange a particle with a reservoir (Algorithm 13), and subroutine `sample` samples quantities every $nsamp$ cycle.

Algorithm 13 (Attempt to Exchange a Particle with a Reservoir)

<pre>SUBROUTINE mcexc if (ranf().lt.0.5) then if (npart.eq.0) return o=int(npart*ranf())+1 call ener(x(o),eno) arg=npart*exp(beta*eno) + /(zz*vol) if (ranf().lt.arg) then x(o)=x(npart) npart=npart-1 endif else xn=ranf()*box call ener(xn,enn) arg=zz*vol*exp(-beta*enn) + /(npart+1) if (ranf().lt.arg) then x(npart+1)=xn npart=npart+1 endif endif return end</pre>	attempt to exchange a particle with a reservoir decide to remove or add a particle test whether there is a particle select a particle to be removed energy particle o acceptance rule (5.6.9) accepted: remove particle o new particle at a random position energy new particle acceptance rule (5.6.8) accepted: add new particle
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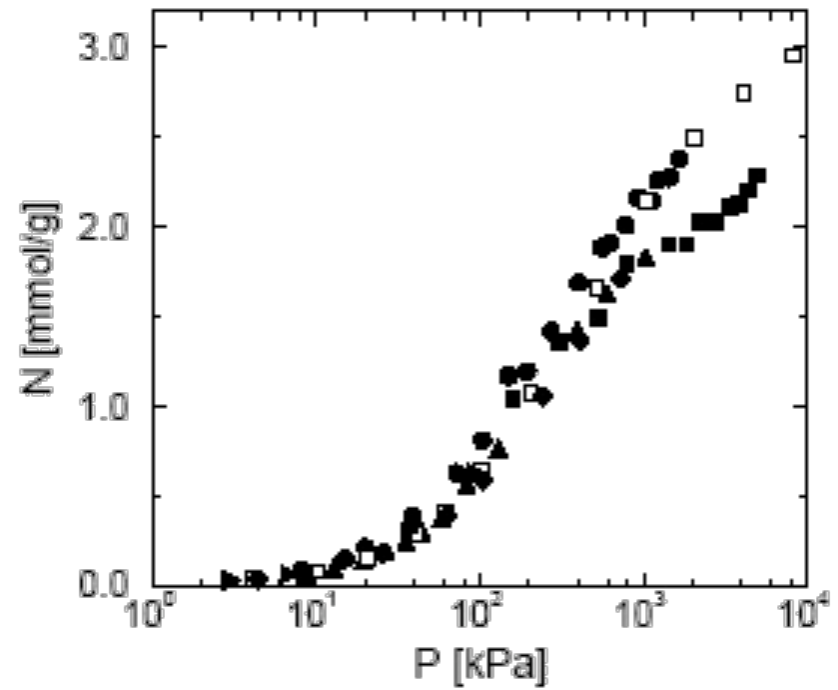
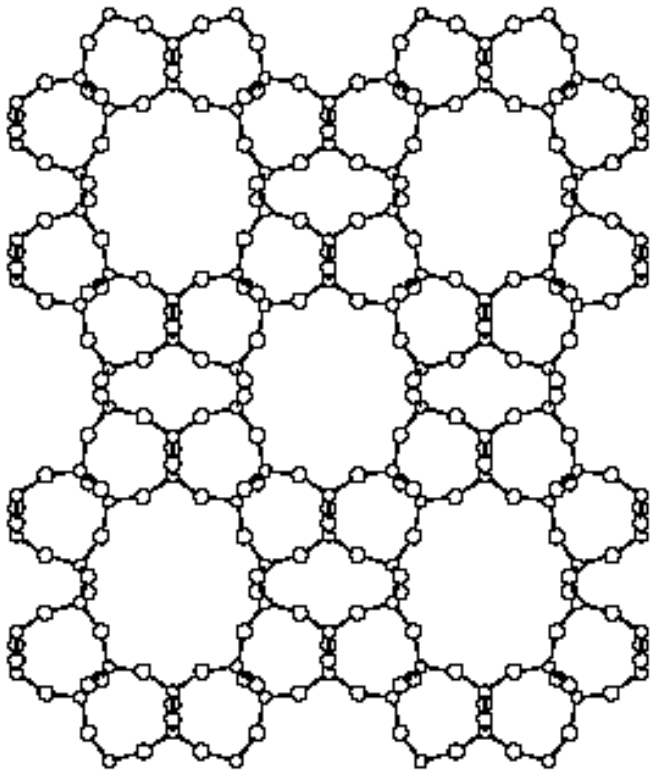
Comment to this algorithm:

Application: equation of state of Lennard-Jones at $T=2.0$



Potential problem: at high densities acceptance goes down.

Application: adsorption in zeolites



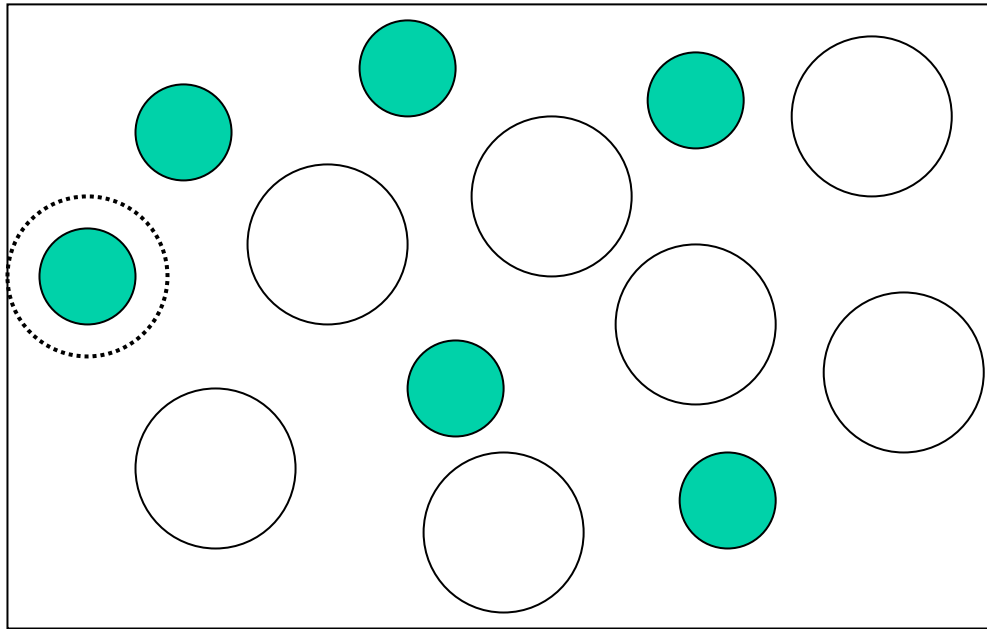
Semi grand ensemble

- For mixtures an additional ensemble exists: semigrand
- Constant $N\Delta\mu T$:
 - The difference between the chemical potentials of the components is fixed
 - total number of particles is fixed
- For a binary mixture
 - N_1, N_2 are allowed to change, but $N_1 + N_2 = N$
 - μ_2, μ_1 are allowed to change, but $\Delta\mu = \mu_2 - \mu_1$

$$\Xi = \frac{\beta P}{N! \Lambda^{3N}} \int dV V^N \exp[-\beta P V] \times \sum_{\text{identities}} \exp[\beta N_2 \Delta\mu] \times \int ds^N \exp[-\beta U(s^N)]$$

In binary mixture just 1,2

MC move changes identity



$$P_{acc}(1 \rightarrow 2) = \min[1, \exp[-\beta U + \beta \Delta \mu]]$$

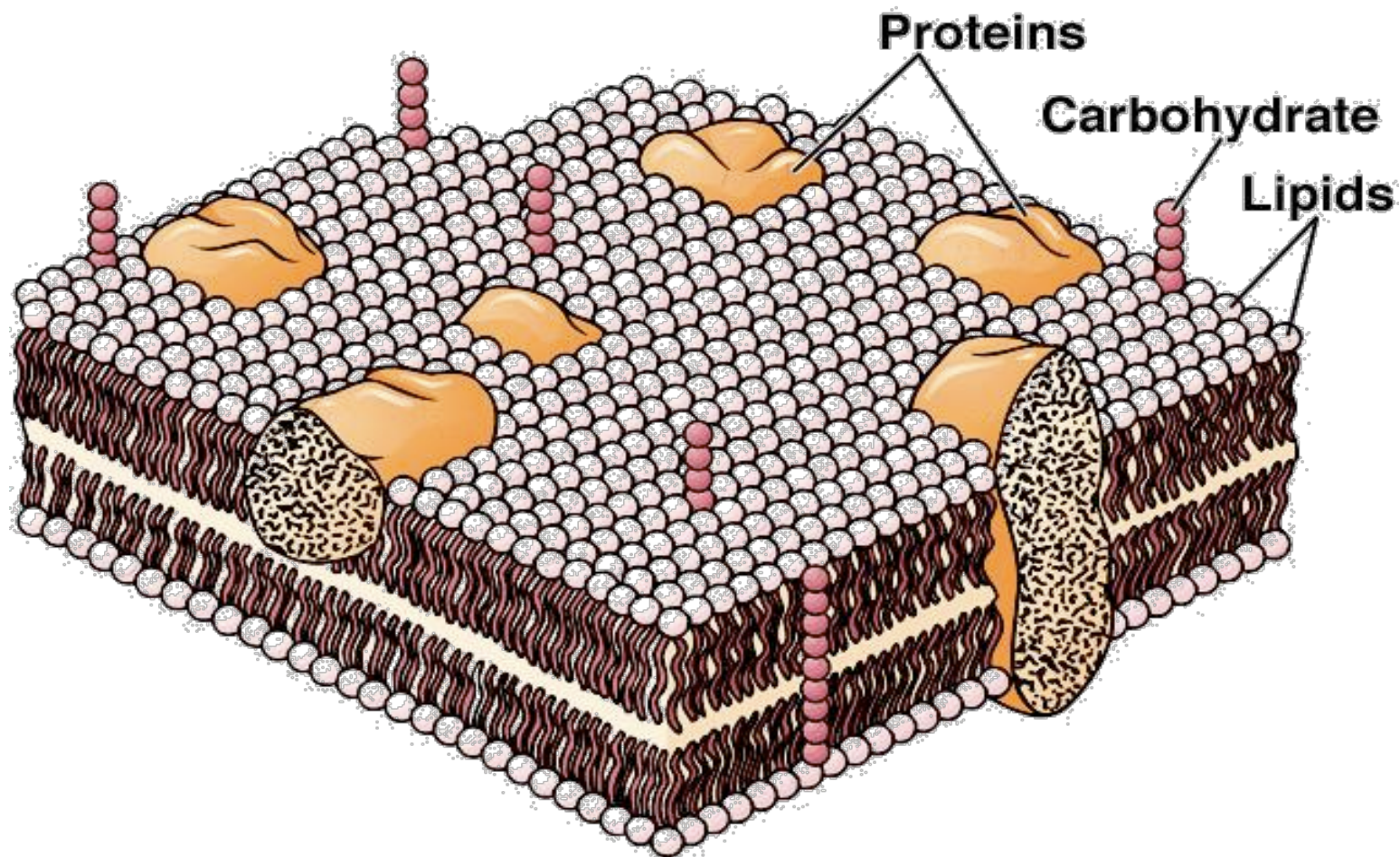
Advantage: at high densities still good acceptance

Summary

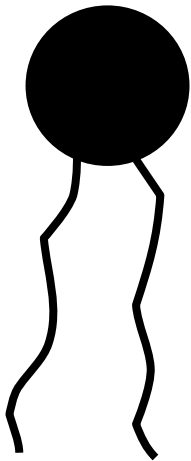
Ensemble	Constant (Imposed)	Fluctuating (Measured)	Function
NVT	N,V,T	P	$\beta F = -\ln Q(N,V,T)$
NPT	N,P,T	V	$\beta G = -\ln Q(N,P,T) = \beta F + \beta PV$
μVT	μ, V, T	N	$\beta \Omega = -\ln Q(\mu, V, T) = -\beta PV$
$\Delta\mu VT$	$\Delta\mu VT$	μ_i, N_i	$\beta Y = -\ln \Xi(\Delta\mu, V, T)$

Exotic ensembles

What to do with a biological membrane?

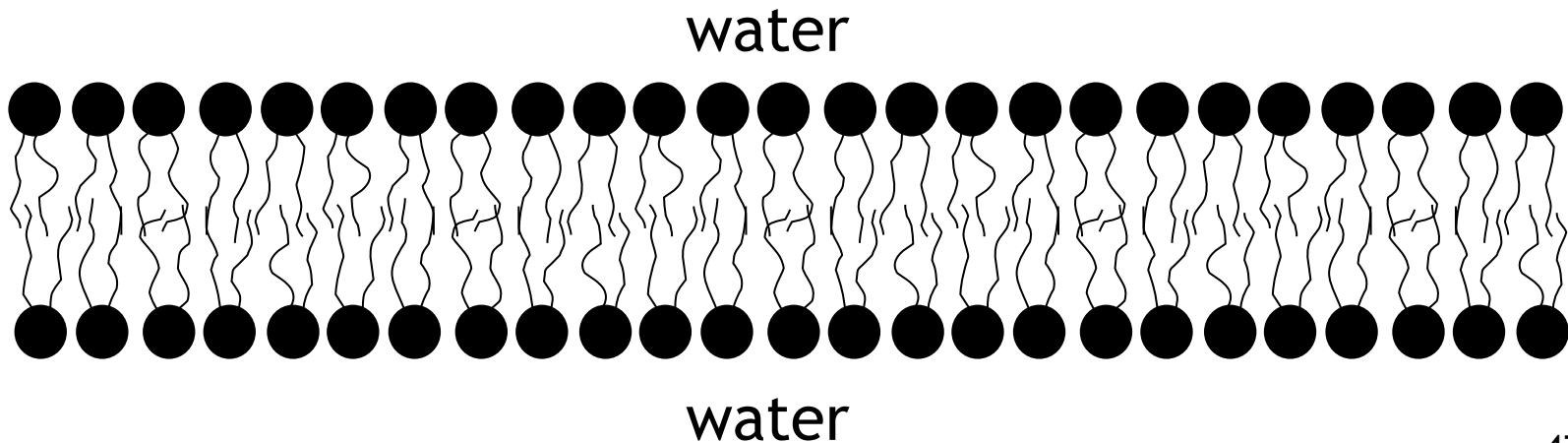


Model membrane: Lipid bilayer

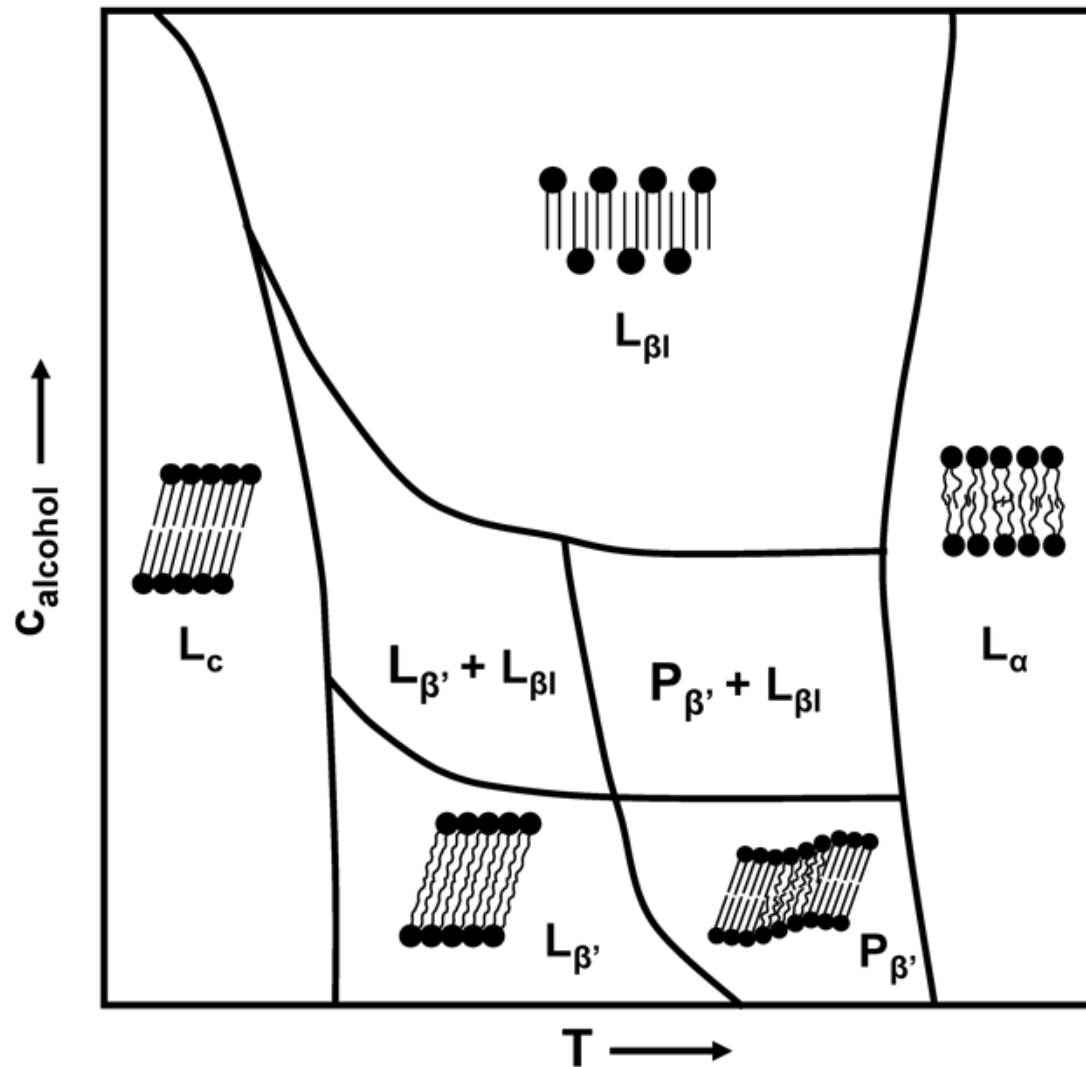


hydrophilic head group

two hydrophobic tails



Phase diagram: lipid-alcohol mixture

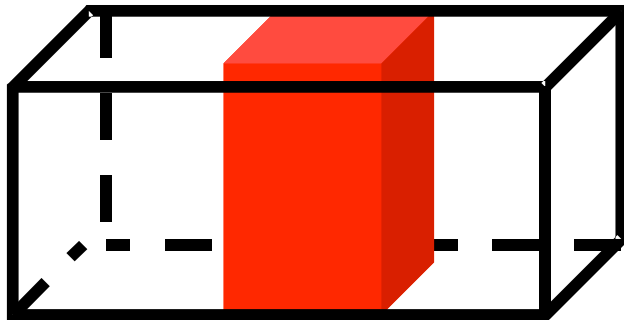


Surface Tension γ controls area per lipid

$\gamma < 0$	compressed bilayer
$\gamma = 0$	tensionless bilayer
$\gamma > 0$	stretched bilayer

Simulations at imposed surface tension

- Simulation to a constant surface tension
 - Simulation box: allow the area of the bilayer to change in such a way that the volume is constant.



Constant surface tension simulation

$$N_{N\gamma T} (A, \mathbf{r}^N) \propto \exp[-\beta (U(\mathbf{r}^N) - \gamma A)]$$

$$A \longrightarrow A'$$

$$L \longrightarrow L'$$

$$AL = A'L' = V$$

$$P_{acc} = \min(1, \exp\{-\beta[U(\mathbf{s}^N; A') - U(\mathbf{s}^N; A) - \gamma(A' - A)]\})$$

Tensionless state: $\gamma = 0$

$$\gamma(\text{Ao}) = -0.3 \pm 0.6$$

$$\gamma(\text{Ao}) = 2.5 \pm 0.3$$

$$\gamma(\text{Ao}) = 2.9 \pm 0.3$$

