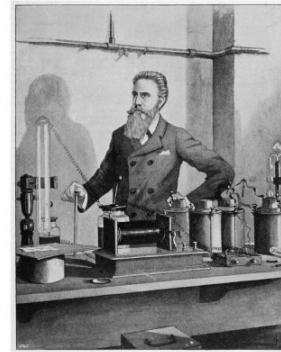


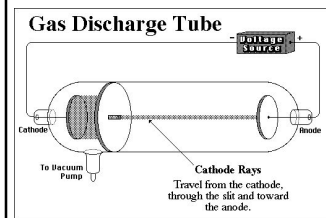
X-ray Spectroscopy

- Lecture 1: Introduction & experimental aspects
- Lecture 2: Atomic Multiplet Theory
Crystal Field Theory
CTM4XAS program
- Lecture 3: Charge Transfer Multiplet Theory
Resonant Inelastic X-ray Scattering
X-ray Spectroscopy on nanomaterials

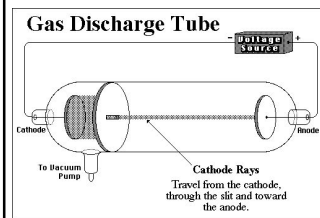
Röntgen's experiment in 1895



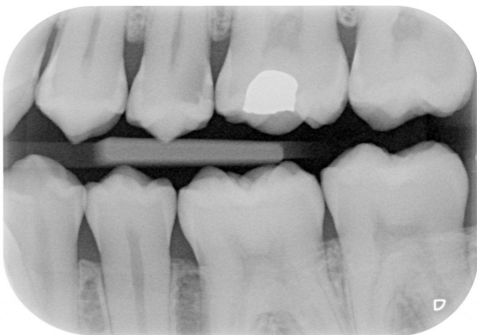
Röntgen's experiment



Röntgen's experiment



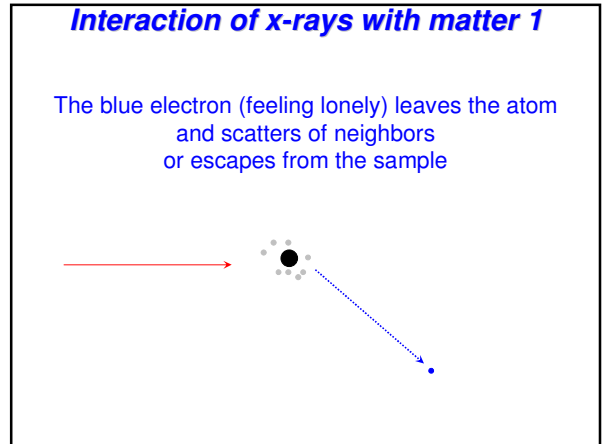
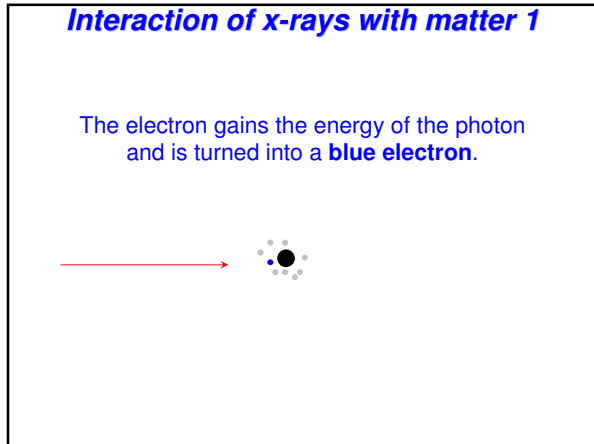
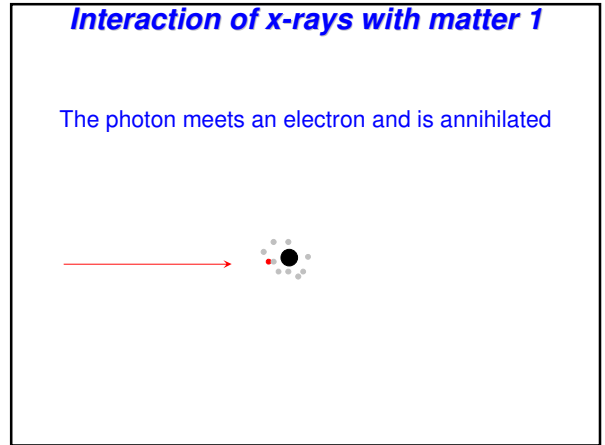
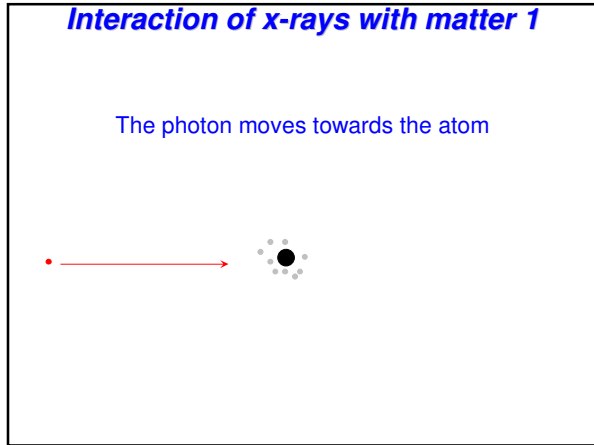
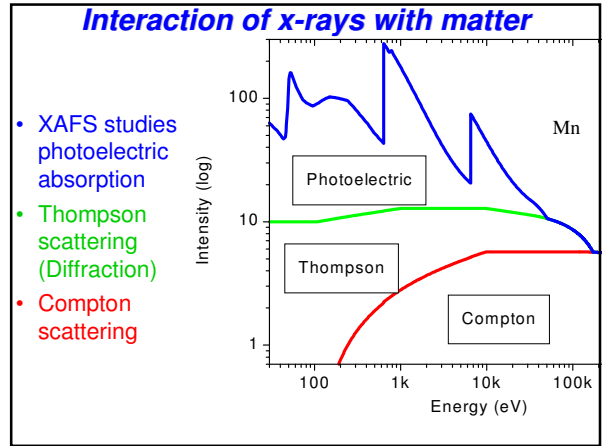
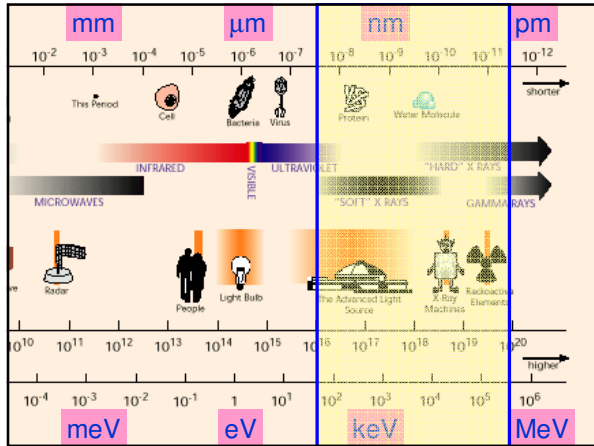
Interaction of x-rays with matter



x-rays

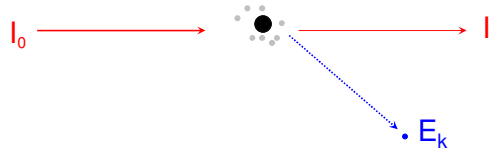
Wavelength: 10^{-10} m

Frequency: $C/\lambda = 3 \cdot 10^{18}$ Hz

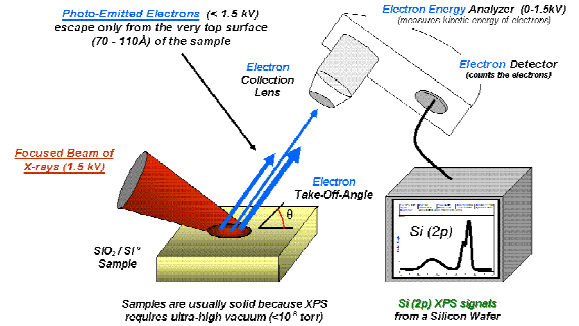


Interaction of x-rays with matter 1

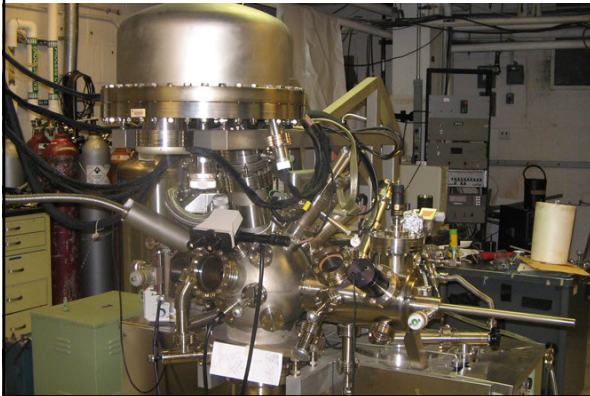
The probability of photon annihilation determines the intensity of the transmitted photon beam



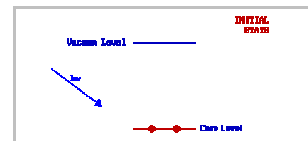
XPS machine



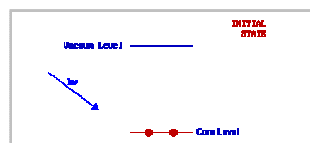
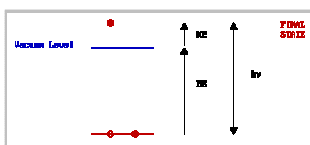
XPS machine



X-ray photoemission



X-ray photoemission



X-ray photoemission: workfunction

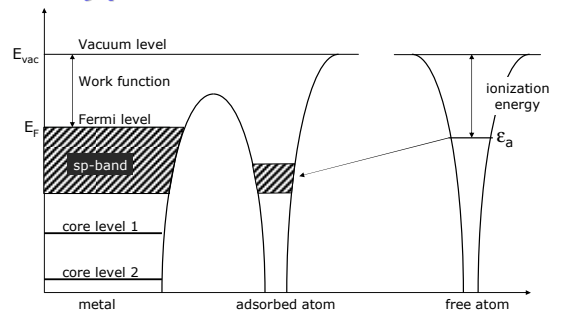
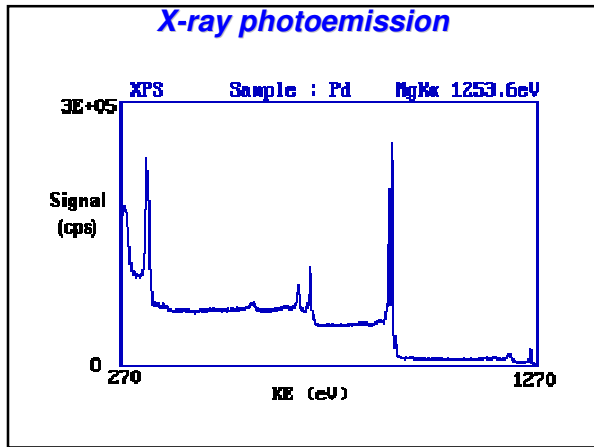
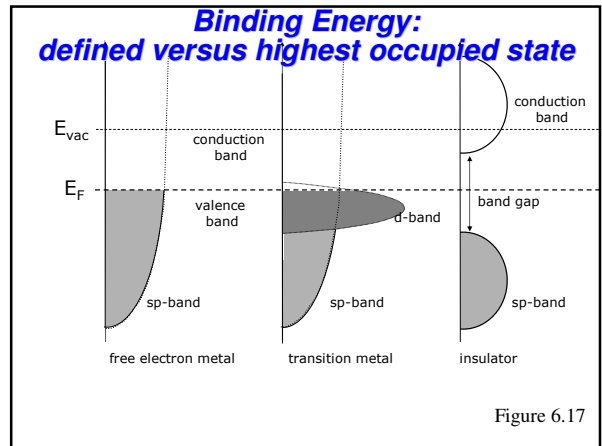
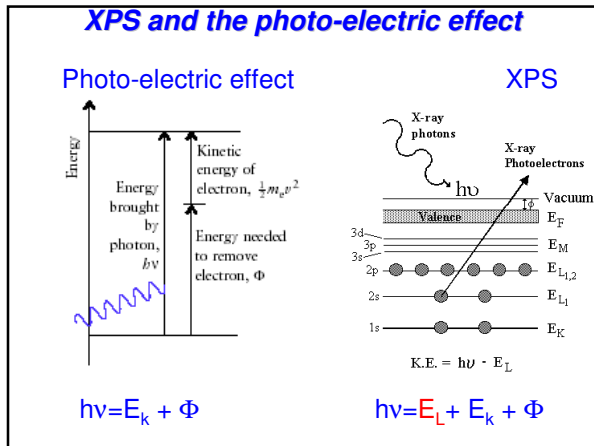


Figure 6.11



Atomic binding energies

THE HIGH FREQUENCY SPECTRA OF THE ELEMENTS (1913)

Moseley (1887-1915)

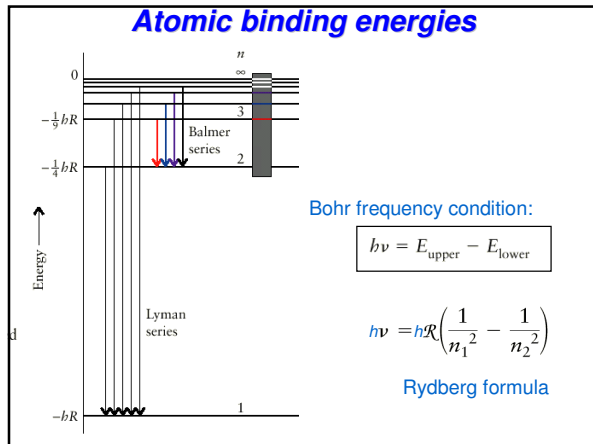
Atomic binding energies

... simple laws have been found which [...] make it possible to predict with confidence the position of the principal lines in the spectrum of any element from aluminum to gold.

Moseley, *Phil. Mag.*, 26, 156(1913); online link at website

Atomic binding energies

$\nu = 2.5 \cdot 10^{15} (Z - \alpha)^2 \text{ Hz}$



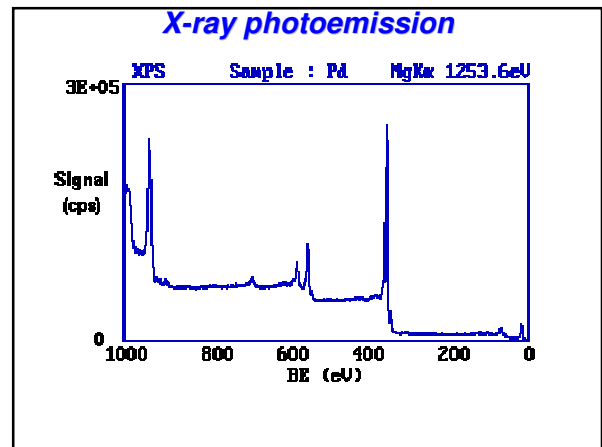
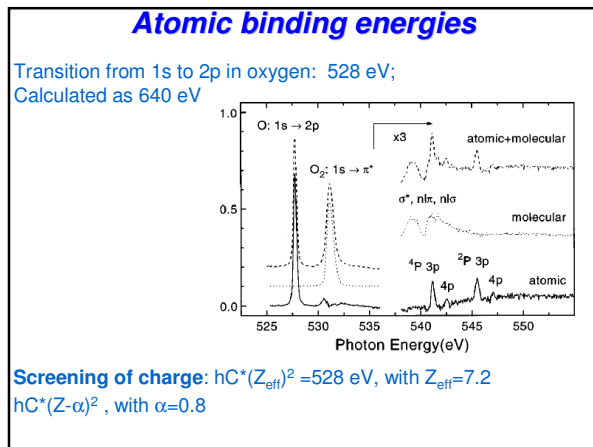
Atomic binding energies

Transition from 1s to 2p in hydrogen:

$$h\nu = hR \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right) = 13.6 * (1/1 - 1/4) \sim 10 \text{ eV}$$

Transition from 1s to 2p in oxygen:

$$h\nu = hR \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right) \cdot Z^2$$

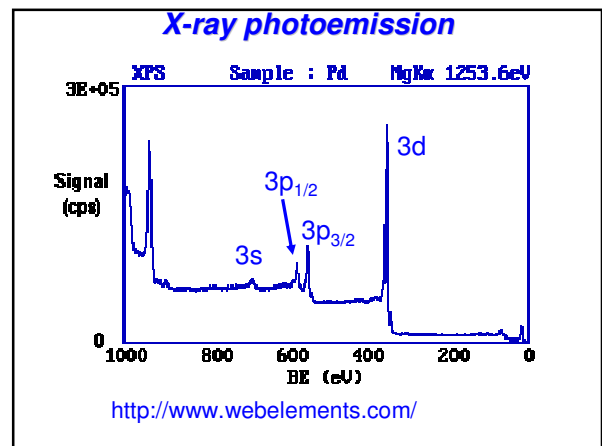
$$= 10 * 8^2 = 10 * 64 = 640 \text{ eV}$$


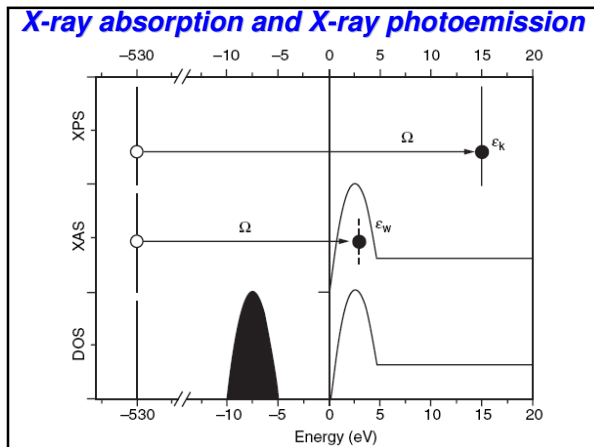
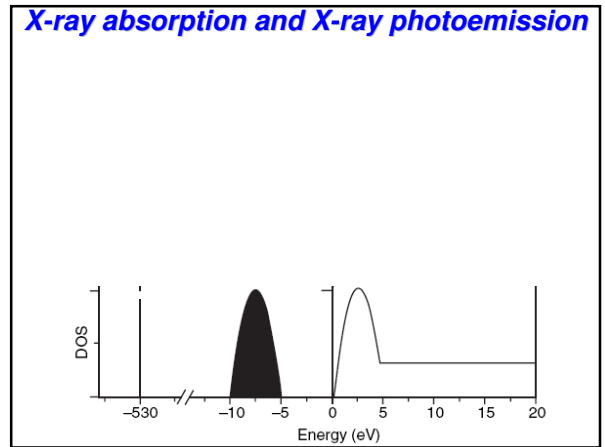
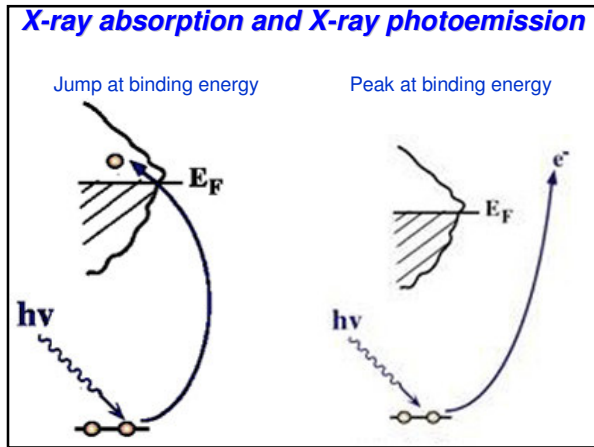
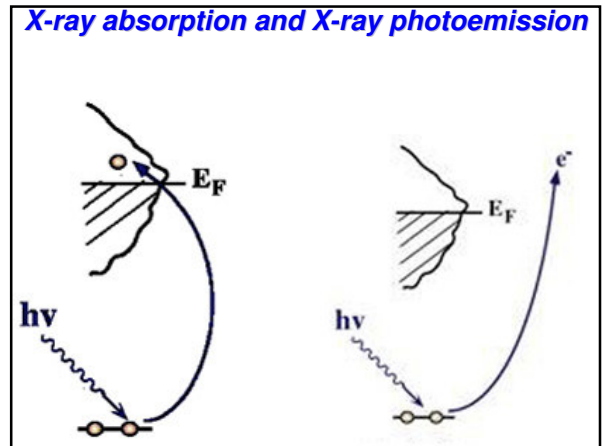
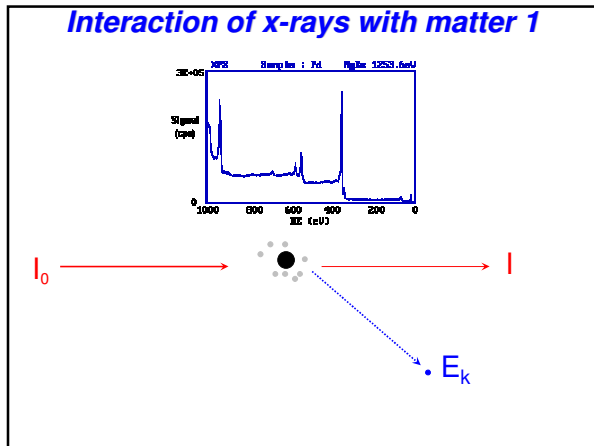
X-ray photoemission edges

Palladium Electron binding energies

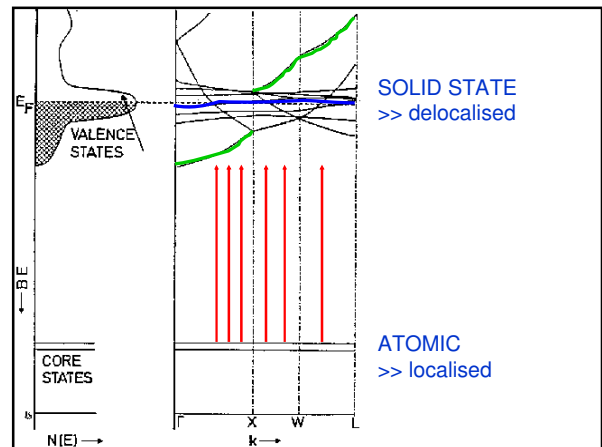
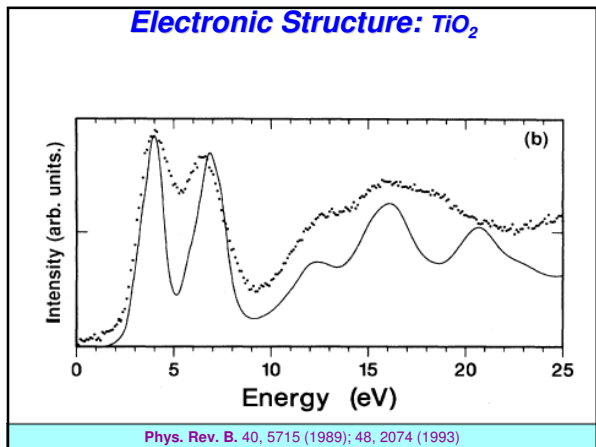
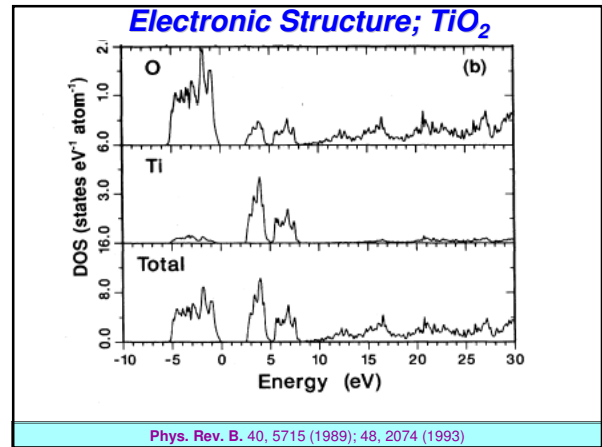
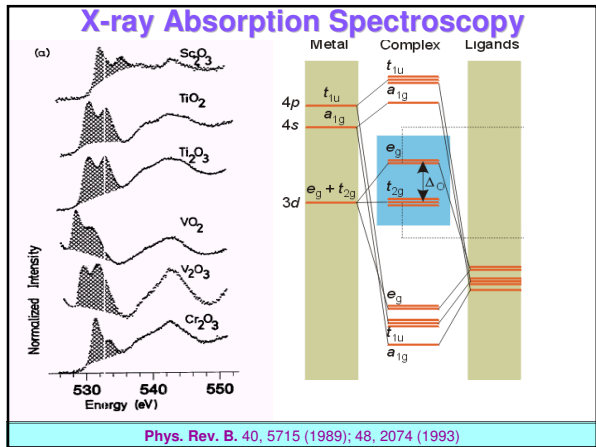
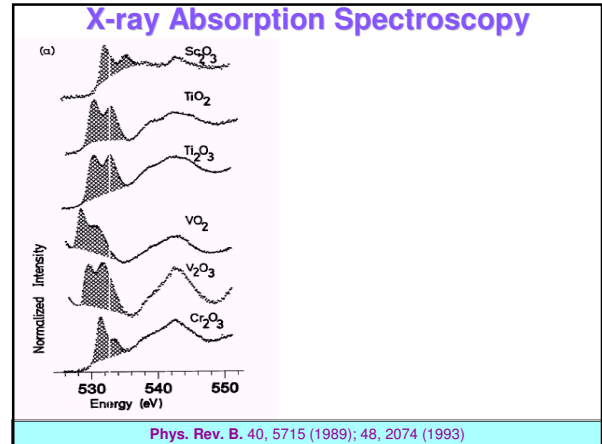
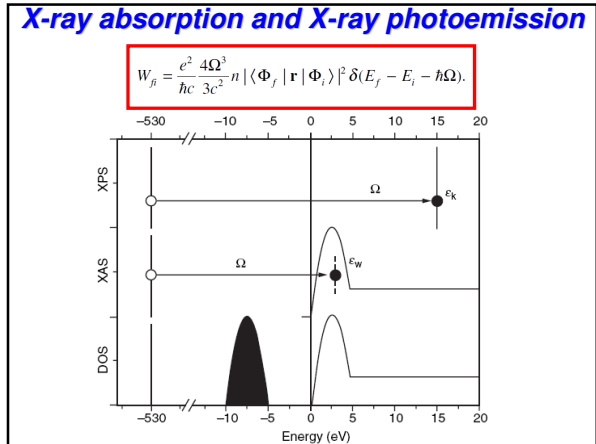
M I	3s	671.6
M II	3p _{1/2}	559.9
M III	3p _{3/2}	532.3
M IV	3d _{3/2}	340.5
M V	3d _{5/2}	335.2
N I	4s	87.1
N II	4p _{1/2}	55.7
N III	4p _{3/2}	50.9

<http://www.webelements.com/>





X-ray absorption experiments



X-ray absorption

Excitation of core electrons to empty states.

Spectrum given by the **Fermi Golden Rule**
(name Golden Rule given by Fermi; rule itself given by Dirac)

$$I_{XAS} \sim \sum_f \left| \langle \Phi_f | \hat{e} \cdot r | \Phi_i \rangle \right|^2 \delta_{E_f - E_i - \hbar\omega}$$

X-ray absorption

Fermi Golden Rule:

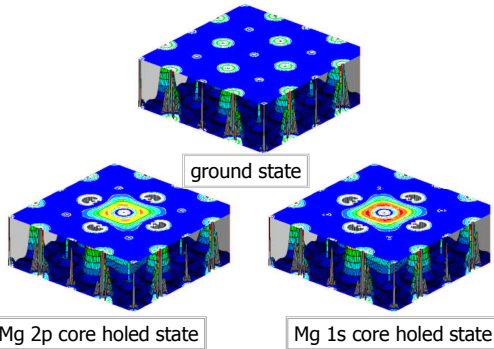
$$I_{XAS} = |\langle \Phi_f | \text{dipole} | \Phi_i \rangle|^2 \delta_{[\Delta E=0]}$$

$$\begin{aligned} \left| \langle \Phi_f | \hat{e}_q \cdot r | \Phi_i \rangle \right|^2 &= \left| \langle \Phi_i | \underline{c} \underline{\mathcal{E}} | \hat{e}_q \cdot r | \Phi_i \rangle \right|^2 \\ &= ?? \left| \langle \mathcal{E} | \hat{e}_q \cdot r | c \rangle \right|^2 \end{aligned}$$

Single electron (excitation) approximation:

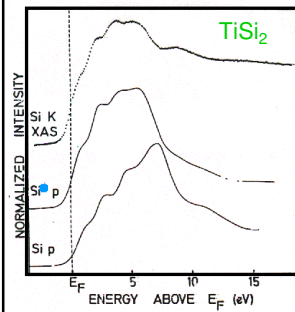
$$I_{XAS} = |\langle \Phi_{\text{empty}} | \text{dipole} | \Phi_{\text{core}} \rangle|^2 \rho$$

X-ray absorption: core hole effect



Tanaka et al. *J. Am. Ceram. Soc.* 88, 2013 (2005).

X-ray absorption: core hole effect



- Final State Rule:**
Spectral shape of XAS looks like final state DOS

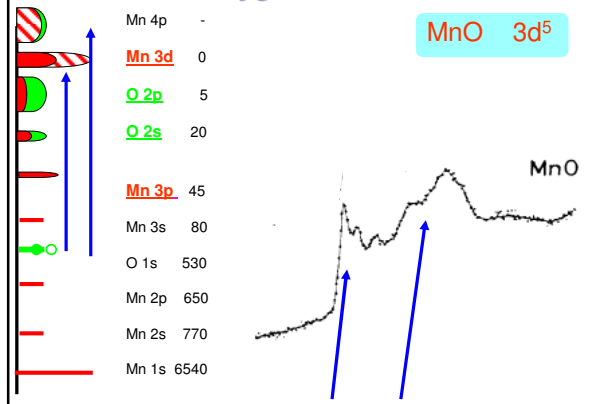
- Initial State Rule:**
Intensity of XAS is given by the initial state

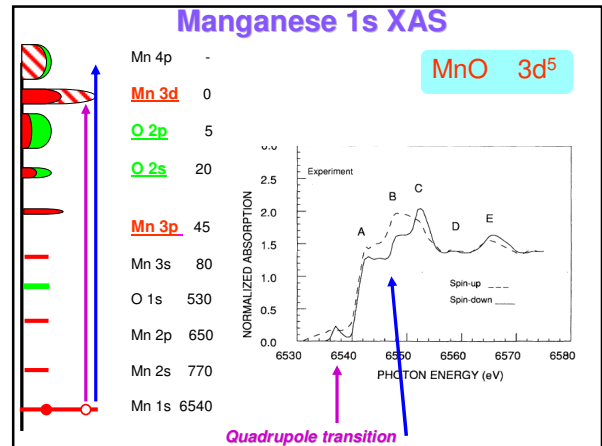
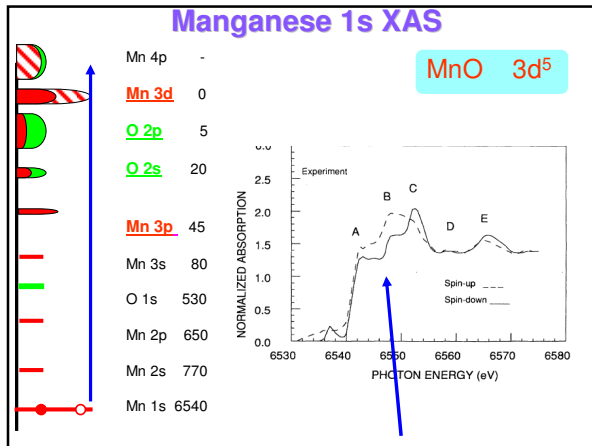
Phys. Rev. B. 41, 11899 (1991)

Ground State



Oxygen 1s XAS

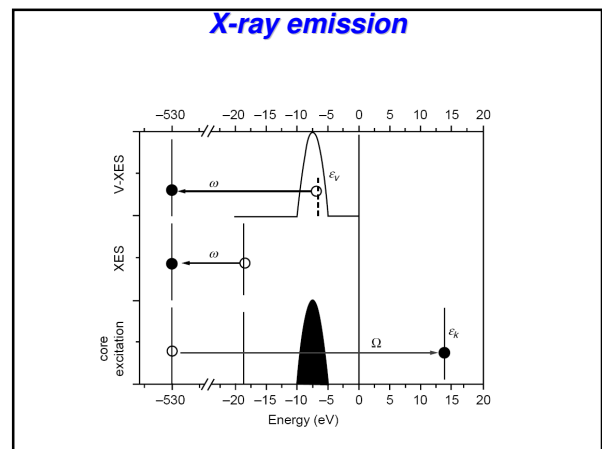
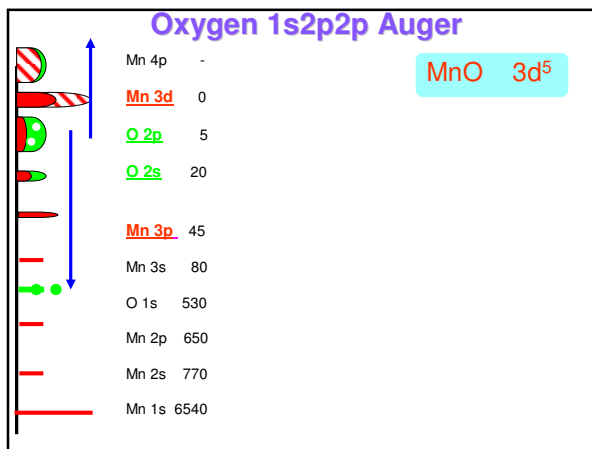
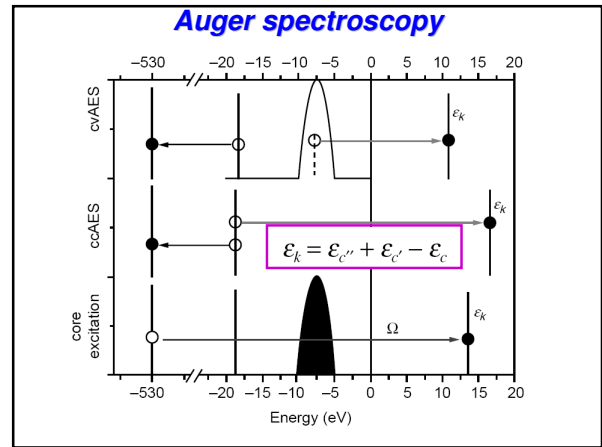


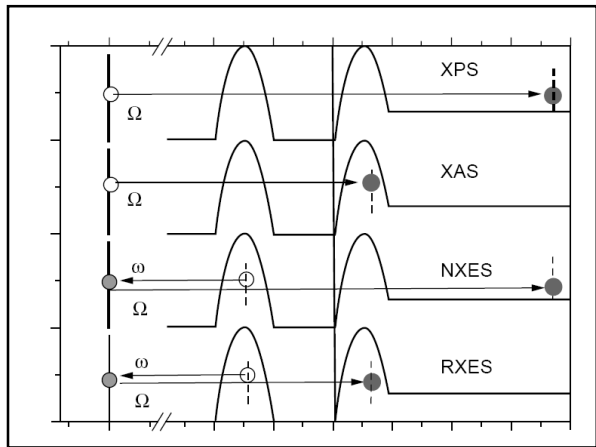
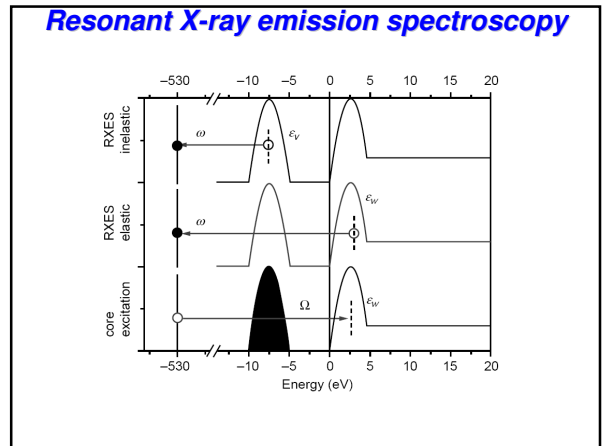
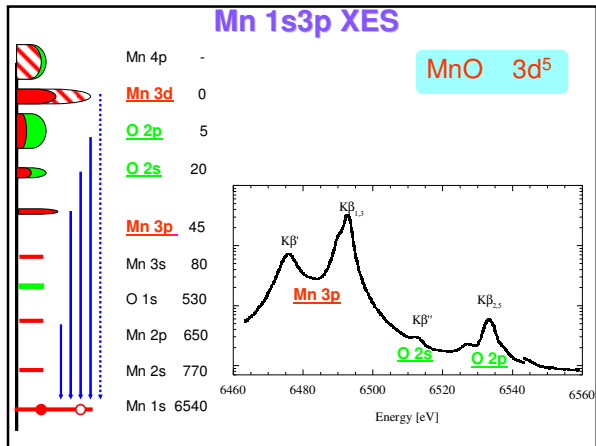


The life of a Core Hole

Heisenberg uncertainty relation
 $\Gamma\tau \cong \hbar$ ($\sim 10^{-16}$ eV s).

Broadening of XAS: ~ 0.1 eV
 Lifetime = 1 femtosecond





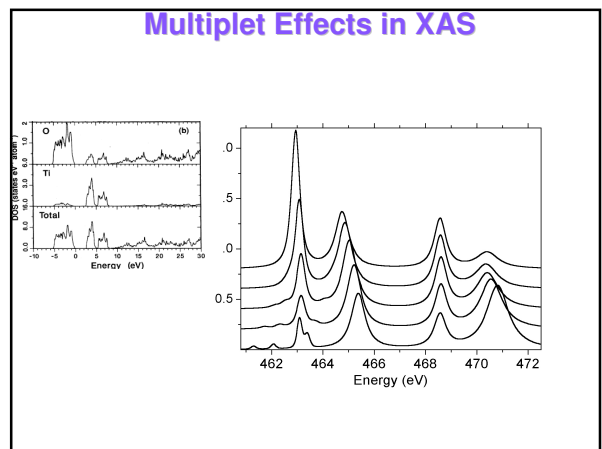
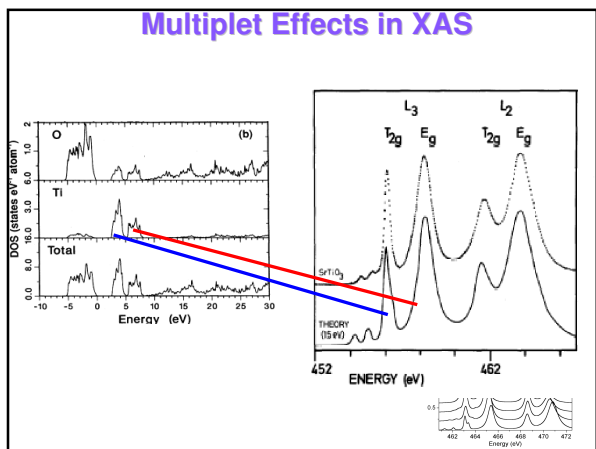
X-ray absorption

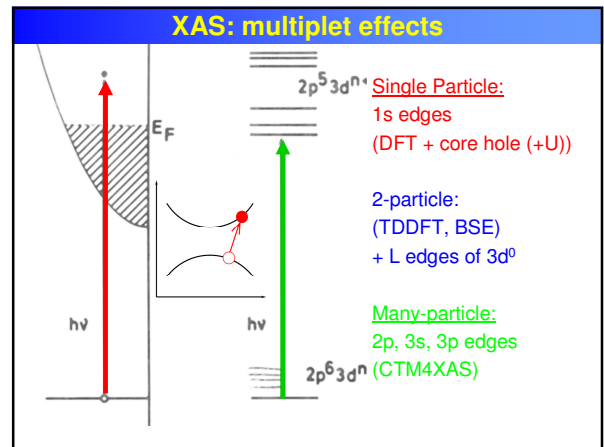
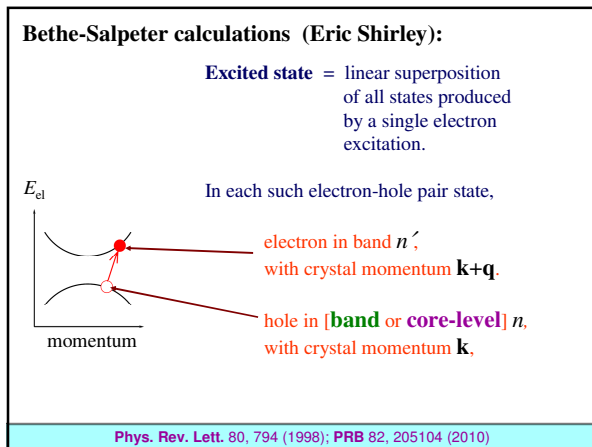
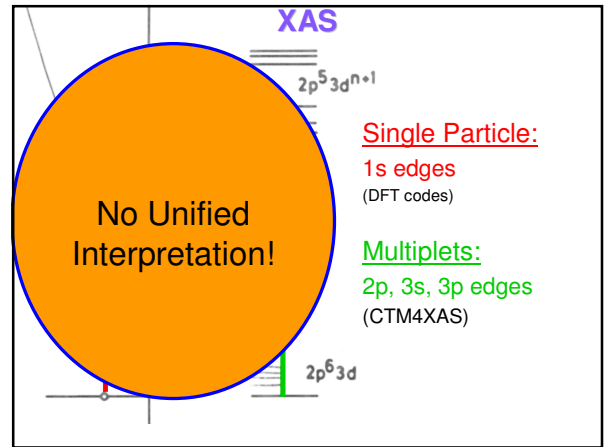
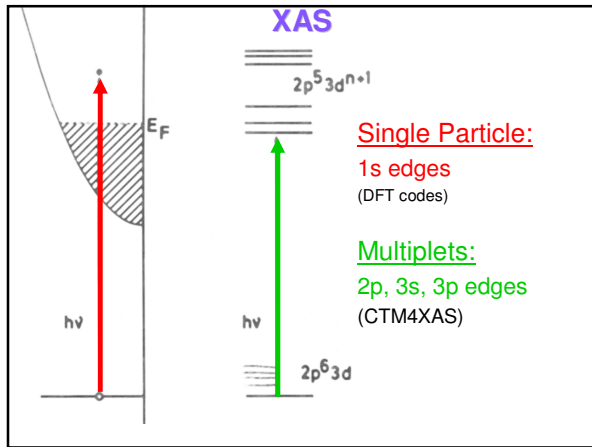
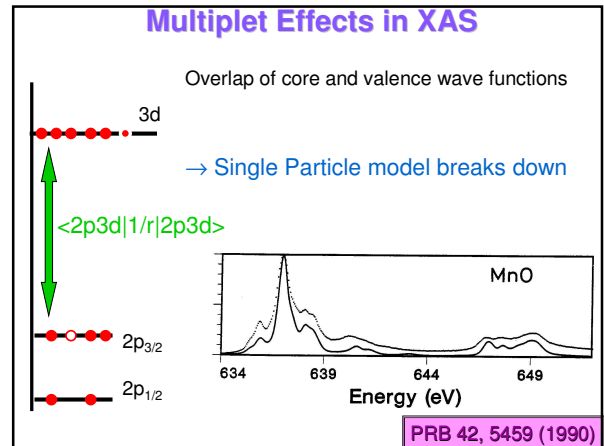
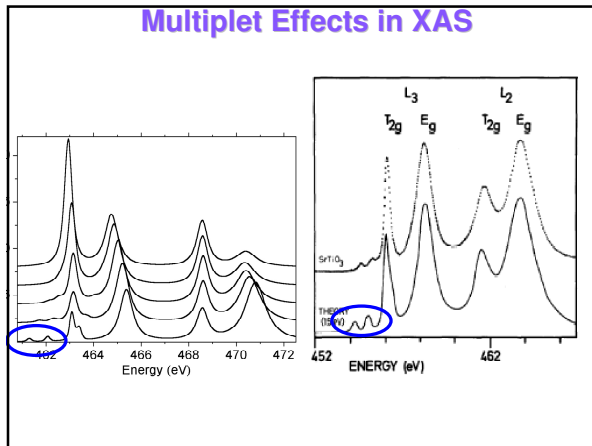
Fermi Golden Rule:
 $I_{XAS} = |\langle \Phi_f | \text{dipole} | \Phi_i \rangle|^2 \delta_{[\Delta E=0]}$

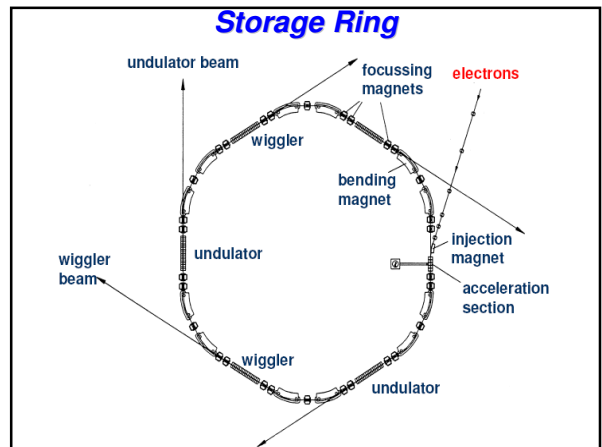
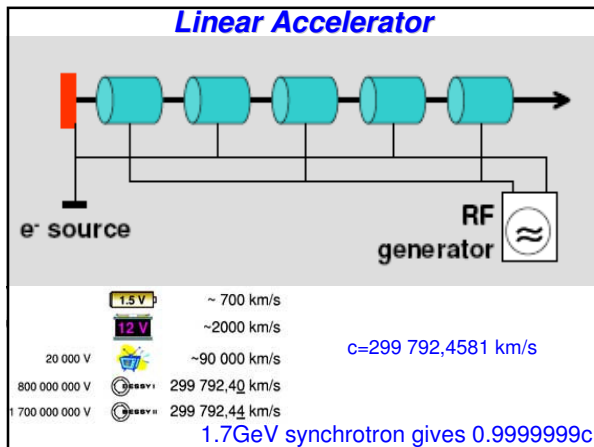
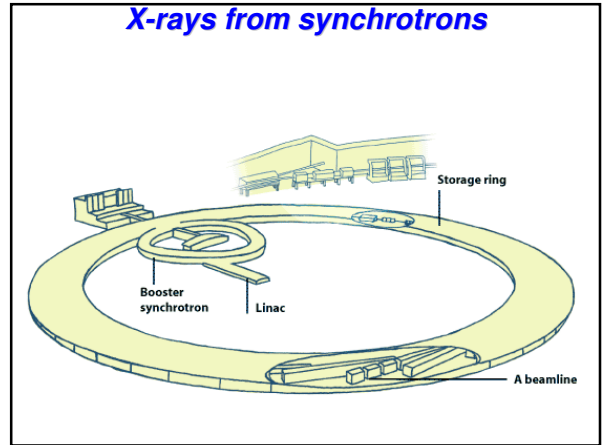
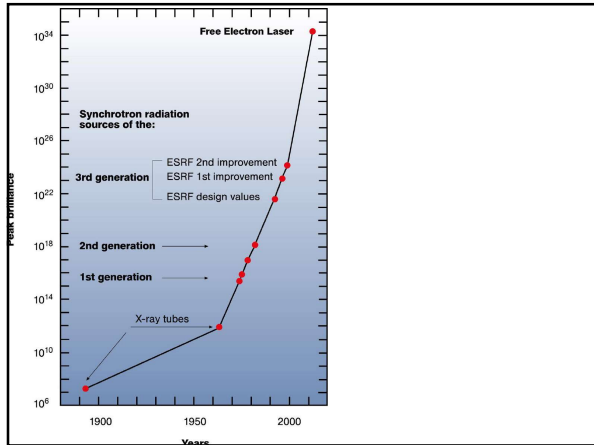
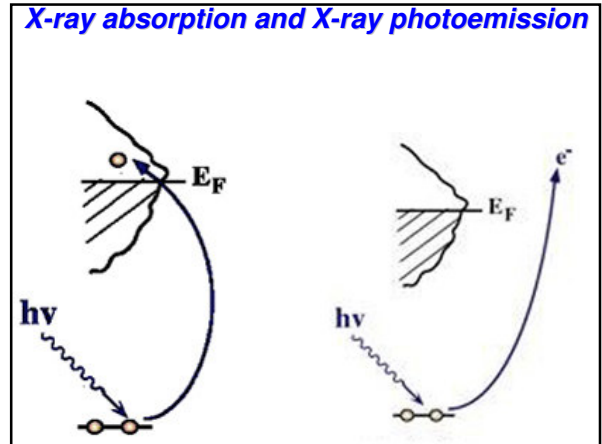
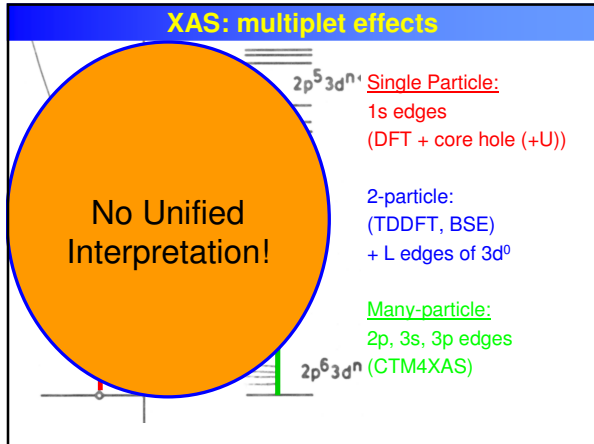
$$\left| \langle \Phi_f | \hat{e}_q \cdot r | \Phi_i \rangle \right|^2 = \left| \langle \Phi_i | \underline{c} \mathcal{E} | \hat{e}_q \cdot r | \Phi_i \rangle \right|^2$$

$$= ?? \left| \langle \mathcal{E} | \hat{e}_q \cdot r | c \rangle \right|^2$$

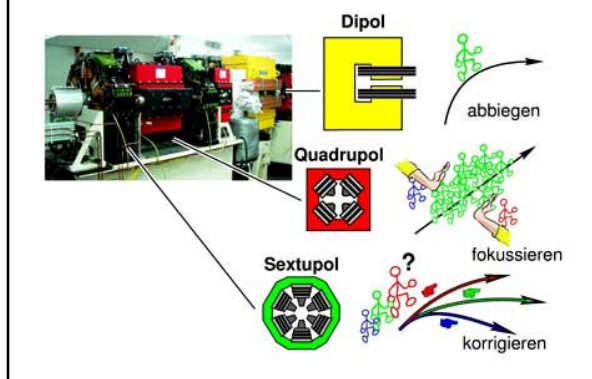
Single electron (excitation) approximation:
 $I_{XAS} = |\langle \Phi_{\text{empty}} | \text{dipole} | \Phi_{\text{core}} \rangle|^2 \rho$



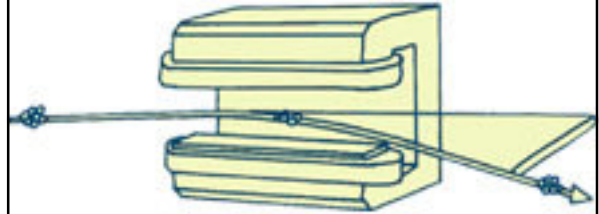




Magnets in the storage ring

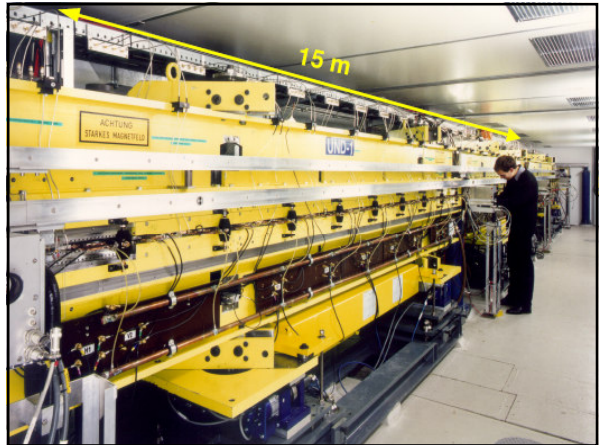
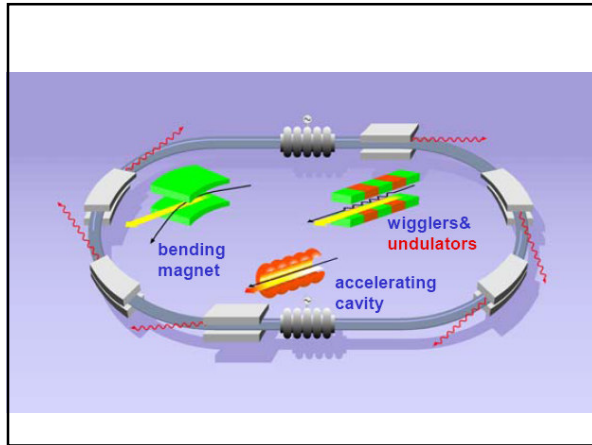


X-rays from synchrotrons

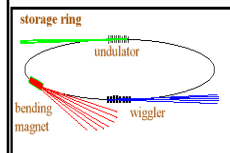


Magnetic Force

- $F = q \cdot v \cdot B$
- Perpendicular to Field and propagation direction

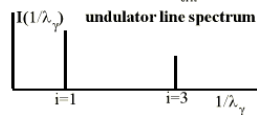
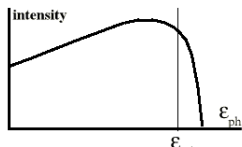


X-rays from synchrotrons

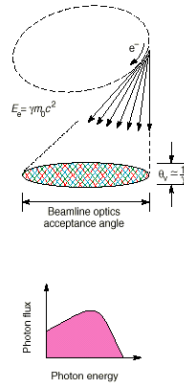


- Bending Magnets
- Insertion Devices:
 - Wiggler
 - Undulator

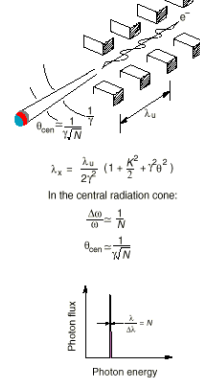
continuous spectrum for wiggler, bending magnet and wave length shifter:

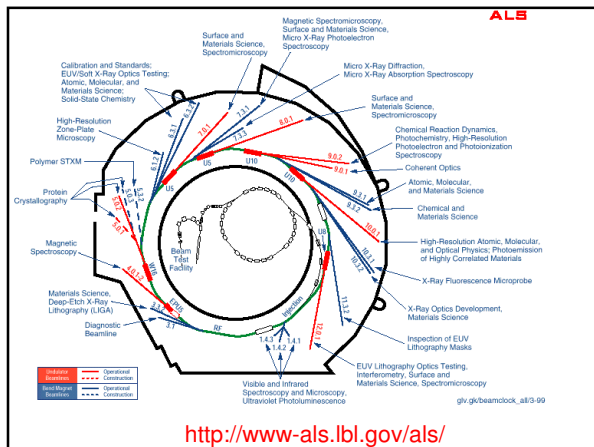
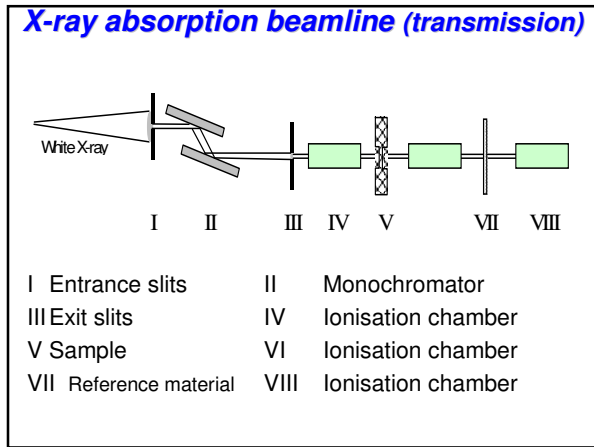
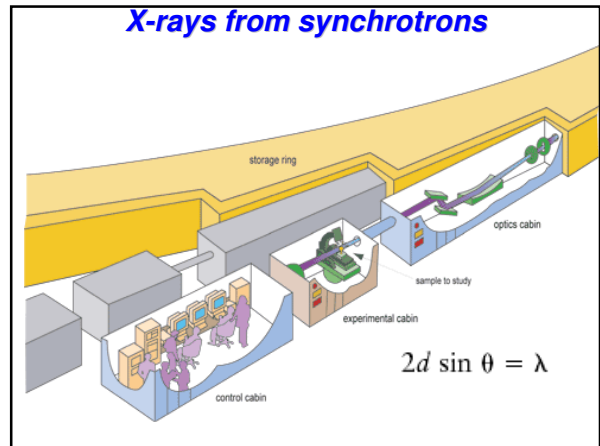
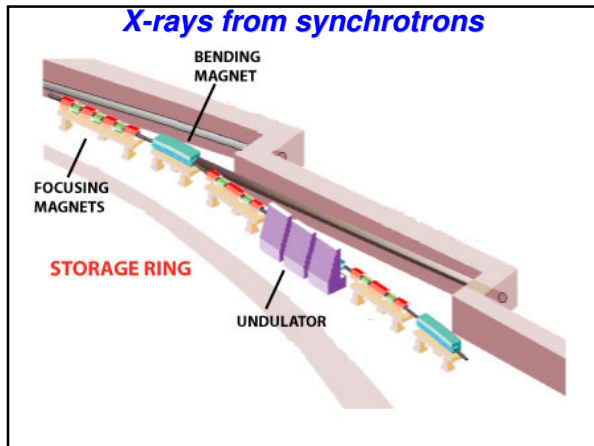


Bend-Magnet Radiation



Undulator Radiation

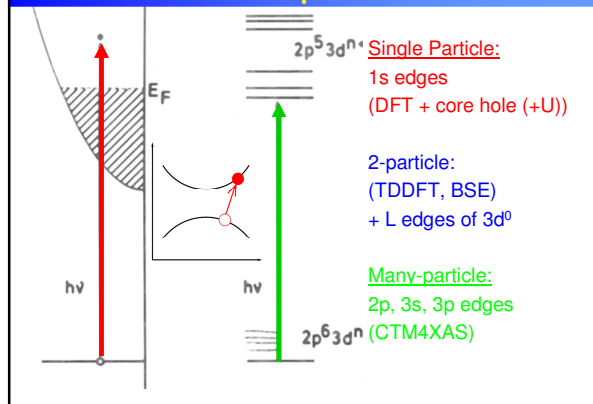




Why a synchrotron?

- **Energy:** tunable source
- **Intensity:** 10^6 - 10^{12} higher than x-ray tube
- **Space:** spot-size 1x1 mm (unfocussed)
down to 20x20 nm (focussed)
- **Time:** pulse 50 ps, sliced down to 50 fs.
- **Polarization:**
Angular dependence & Circular dichroism:
- **NOT coherent** (no laser)

XAS: multiplet effects



Charge Transfer Multiplet program

Used for the analysis of XAS, EELS,
Photoemission, Auger, XES,

ATOMIC PHYSICS
↓
GROUP THEORY
↓
MODEL HAMILTONIANS

The CTM4XAS program

